

MACROMOLECULAR CHEMISTRY  
AND POLYMERIC MATERIALS

## Specific Features of Interaction of Polyacrylic and Polymethacrylic Acids with Nanoparticles of Amorphous Selenium

V. V. Kopeikin, S. V. Valueva, A. I. Kipper, A. P. Filippov, G. V. Shishkina, L. N. Borovikova, N. V. Rumyantseva, Ya. I. Nazarkina, and V. K. Lavrent'ev

*Institute of Macromolecular Compounds, Russian Academy of Sciences, St. Petersburg, Russia*

Received November 25, 2003; in final form, February 2004

**Abstract**—A study was made of formation of selenium nanoparticles in water from reduction of sodium selenite with ascorbic acid in the presence of polyacrylic and polymethacrylic acids.

It is well known that nanoparticles of zero-valence metals and nonmetals, with their specific structural features and high surface energy, exhibit abnormal physicochemical properties because of the scale effect of matter [1]. Of special interest are nanoparticles of selenium, a chemical element with unique photoelectric, semiconducting, catalytic, and biological properties and X-ray sensitivity.

Ligand-free nanoparticles are metastable and tend to rapidly aggregate in aqueous solutions. Therefore, they are usually stabilized by water-soluble polymers whose interaction with nanoparticles yields nanostructures of various morphologies [2]. However, the nature of this interaction in formation of nanoparticles in aqueous solutions of polymers still remains to be elucidated.

This work was aimed at elucidating the nature of the above-mentioned interactions and their role in polymer stabilization of selenium nanoparticles (nano-Se<sup>0</sup>). To this end, we studied for the first time formation of nano-Se<sup>0</sup> particles in the presence of synthetic polymers differing in the hydrophylic–lipophilic balance. Also, we carried out optical studies of the morphological characteristics of the nanostructures being formed.

### EXPERIMENTAL

We studied nano-Se<sup>0</sup> nanoparticles prepared by reducing selenious acid (SA) with ascorbic acid, with and without polymer additions.

As polymeric stabilizers of nano-Se<sup>0</sup> served synthetic polymers, polyacrylic acid (PAA) and poly-

methacrylic acid (PMAA), with molecular weights  $M_w$  of  $27 \times 10^6$  and  $25 \times 10^3$ , respectively. This choice was dictated by close structures and Kuhn segment sizes of these polymers, of which only PMAA has side hydrophobic methyl groups [3].

The ionic selenium species  $\text{SeO}_3^{2-}$  was reduced at the selenium concentration in solution of 0.01 wt % and the polymer-to-selenium mass ratio of 9 : 1 (i.e., under exhaustive saturation of the adsorption capacity of nanoparticles, as revealed for other selenium–polymer nanostructures [4]).

Using the method of flow birefringence [5], we estimated the molecular dispersity of the solutions of the adsorbates being formed, based on analysis of the gradient dependence of the birefringence  $\Delta n$ . Specifically, we studied the birefringence versus the gradient of the rotor rotation rate  $g$  and the polymer concentration  $c$  at a fixed selenium concentration of 0.01 wt %. We used titanium dynamometer with a 4-cm-high internal rotor and the rotor–stator slit width of 0.03 cm. The solutions for birefringence measurements were kept at a constant temperature of 21°C to prevent variations of their viscosities and optical distortions due to the temperature gradient. The measuring facility was calibrated against 2-phenylethanol having a significant birefringence ( $\Delta n/g = 17 \times 10^{-12}$ ), as well as against the polystyrene–bromoform system. The intrinsic birefringence  $[n] = \lim_{g \rightarrow 0, c \rightarrow 0} (\Delta n/gc\eta_0)$  ( $\eta_0$  is the viscosity of the solvent) was measured accurately to within 10%. The measurements were carried out at  $g < g_k$ , where  $g_k$  is the rate gradient corresponding to the onset of flow turbulence.

The experimentally determined intrinsic birefringence  $[n]$  in a general case ( $dn/dc \neq 0$ , where  $dn/dc$  is the increment of the refractive index of the solution, estimated for PAA–Se<sup>0</sup>–H<sub>2</sub>O and PMAA–Se<sup>0</sup>–H<sub>2</sub>O systems at 0.245 and 0.165, respectively) has three components:  $[n] = [n]_e + [n]_{fs} + [n]_f$ , where  $[n]_e$  is the intrinsic anisotropy;  $[n]_{fs}$ , microform effect; and  $[n]_f$ , macroform effect [5]. The full segment anisotropy  $[n]_{fs} + [n]_e$  is determined by the equilibrium rigidity  $A$  of the polymer chain and by the structure of the polymer unit, and  $[n]_f$  is related to the shape asymmetry parameter  $p$  of the macromolecule (particle) as

$$[n]_f = [(n_s^2 + 2)/3]^2 [M_w (dn/dc)^2 f(p)] / (30\pi RT n_s) \\ = \text{const } M_w (dn/dc)^2 f(p), \quad (1)$$

where  $M_w$  is the molecular weight of the adsorbate macromolecule;  $n_s$ , refractive index of the solvent;  $T$ , absolute temperature;  $R$ , universal gas constant; and  $f(p)$ , tabulated function of the ratio of the axes of a rigid (impermeable by the solvent) ellipsoid approximating the particle [5].

Relationship (1) is valid both for isolated macromolecule and nanostructure; all what is needed of a particle is to be impermeable by the solvent. Thus, for high-molecular-weight flexible-chain and moderately rigid-chain polymers ( $A < 5$  nm), having minor optical anisotropy and low-molecular-weight units, the intrinsic birefringence can be fitted by the approximation  $[n] \approx [n]_f$ , which allows the parameter  $p$  to be estimated directly from the experimentally measured birefringence [5].

Using the elastic (static) light scattering method [6], we estimated  $M_w$  and the root-mean-square inertia radii  $R_g$  of the polymers and nanostructures, as well as their affinity for solvent, water (from the second virial coefficient  $A_2$ ). Based on the ratio of the parameters  $M_w$  for poly acids and nanostructures they form, we determined the number of the macromolecules adsorbed on the nano-Se<sup>0</sup> surface. The reduced intensity of light scattering by solutions  $R_0$  was measured on a Fica photogoniometer. The wavelength of the incident, vertically polarized, light was 546.1 nm. The measurements were carried out at light scattering angles  $\theta$  within 30°–150°. The solutions and solvents were purified by centrifugation at  $15 \times 10^3$  rpm for 1–1.5 h. The refractive index increment  $dn/dc$  was determined from the refractive indices measured on an IRF-23 instrument.

The experimental light scattering data were proc-

essed by the Zimm method, by double extrapolation ( $c \rightarrow 0$ ,  $\theta \rightarrow 0$ ) of the plots of  $Kc/R_\theta$  vs.  $\sin^2(\theta/2) + kc$  ( $K$  is the calibration constant, and  $k$ , numerical constant).

We determined the average hydrodynamic particle sizes  $R_H$  by quasielastic (dynamic) light scattering method [7]. Based on the experimental  $R_g$ -to- $R_H$  ratio for the nanostructures, we estimated the parameter  $\rho^*$  characterizing the conformation of the nanostructure [8–11].

The optical component of the facility for measuring the dynamic light scattering comprised an ALV-SP (Germany) goniometer (a Spectra-Physics He–Ne laser with a wavelength  $\lambda = 632.8$  nm, power ca. 20 mV, as light source). The correlation function of the scattered light intensity was obtained on a 288-channel Photo Cor-FC correlator (Anteks Closed Joint-Stock Company, Russia). Data processing was carried out by standard methods of cumulants and Tikhonov regularization, using appropriate software.

According to the static light scattering data, the parameter  $M_w$  remains unchanged ( $M_w = 27 \times 10^7$ ) in going from the PAA–H<sub>2</sub>O to PAA–Se<sup>0</sup>–H<sub>2</sub>O system, i.e., in the presence of PAA, no nanostructure is formed. For the PMAA–Se<sup>0</sup>–H<sub>2</sub>O system,  $M_w$  was estimated at  $1.5 \times 10^8$ , i.e., it increased by a factor of 4300 relative to free PMAA macromolecules. This unambiguously suggests the deciding importance of hydrophobic interactions between macromolecules and nano-Se<sup>0</sup> in formation of nanostructures.

The second virial coefficient  $A_2$  for the PMAA–Se<sup>0</sup>–H<sub>2</sub>O system is equal to zero, which suggests an ideal thermodynamic quality of the solution. For the PAA–Se<sup>0</sup>–H<sub>2</sub>O system,  $A_2$  is negative, namely,  $-0.8 \times 10^{-4}$  cm<sup>3</sup> mol g<sup>-2</sup>, which suggests a bad thermodynamic quality of the solution. Thus, adsorption of macroions on nano-Se<sup>0</sup> improves the thermodynamic quality of the solution.

The parameter  $R_H$  for nano-Se<sup>0</sup> particles obtained without polymeric stabilizer was estimated at 170 nm. The particle size of the PMAA–Se<sup>0</sup> adsorbate, estimated by dynamic and static light scattering methods, is much smaller (root-mean-square inertia radius  $R_g = 90$  nm and mean hydrodynamic radius  $R_H = 90$  nm). The parameter  $\rho^* = R_g/R_H$  is equal to unity, which suggests virtually ideal spherical shape of the particles of the adsorbate being formed [8–11].

To determine the radius of the selenium particles occurring in the core of the nanostructure, we lyophilically dried an aqueous solution of PMAA–Se<sup>0</sup>,

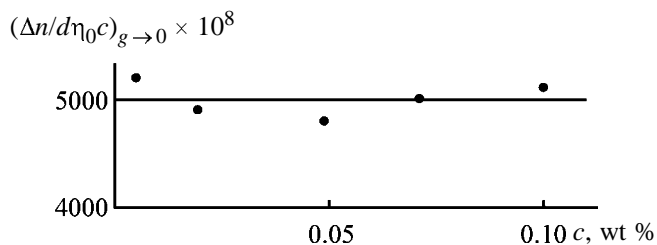
and the resulting dry pelletized product was examined by small-angle X-ray scattering. We estimated the radius of the spherical  $\text{Se}^0$  particles at 15 nm, which is by an order of magnitude smaller than that of selenium particles obtained without polymeric stabilizer. Thus, the nano- $\text{Se}^0$  particles in aqueous solutions have a ca. 75-nm-thick polymer shell.

Optical studies of the PAA and PMAA macromolecules showed that, in the nonionized state, these are Gaussian globules with the asymmetry parameter  $p \sim 2$  [12].

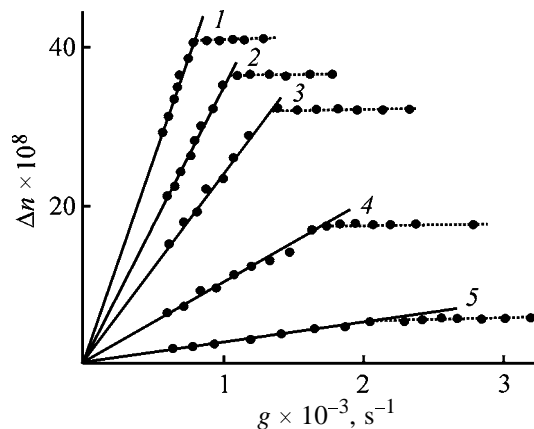
Calculation of the parameter  $p$  in the  $[n] \sim [n]_f$  approximation for the PAA- $\text{Se}^0$ - $\text{H}_2\text{O}$  and PMAA- $\text{Se}^0$ - $\text{H}_2\text{O}$  systems yielded, respectively,  $p \sim 1$  (spherical conformation) and  $p \sim 2$  (which is inconsistent with the value of the parameter  $p^*$ , corresponding to the spherical conformation of the nanostructure) [5]. Also, the fact that the reduced optical anisotropy for the PMAA- $\text{Se}^0$ - $\text{H}_2\text{O}$  system is independent of the concentration (Fig. 1) suggests that the parameter  $p$  tends to unity, i.e., that the nanostructure is spherical [5]. The fact that hybrid nanostructures formed by macromolecules and nanoparticles have a spherical shape, even in the case of such a rigid macromolecule as DNA [13], suggests that the  $[n] \sim [n]_f$  approximation is invalid for the PMAA- $\text{Se}^0$ - $\text{H}_2\text{O}$  system. Evidently, in this case a significant contribution to the optical anisotropy comes from an effect similar to the microshape effect for an isolated macromolecule (Lyubina [12] showed that the intrinsic optical anisotropy even for PMAA with its high degree of ionization is negligible). This is supported by the fact that the  $\Delta n = f(g)$  curves (Fig. 2) for the PMAA- $\text{Se}^0$ - $\text{H}_2\text{O}$  system reach saturation. This suggests an increase in the rigidity of the PMAA macromolecules owing to hydrophobic interactions. Thus, hydrophobic interactions of nonpolar fragments of the macromolecules with one another and with nano- $\text{Se}^0$  constitute the major driving force of adsorption.

In the case of PAA, no nanostructure is formed, and the conformation of the macroions changes, in contrast to nonionic hydrophilic polymers such as polyacrylamide [4]. This is presumably due to dipole-dipole interactions of PAA with selenium nanoparticles.

We showed that selenium nanoparticles are adsorption matrices suitable for accommodating in high local concentrations anionic polyelectrolytes containing hydrophobic methyl groups in the side chains. Our results can be applied to molecular designing of new materials with desired properties.



**Fig. 1.**  $(\Delta n/d\eta_0c)_{g \rightarrow 0} \sim c$  dependence for the PMAA- $\text{Se}^0$ - $\text{H}_2\text{O}$  system.  $(\Delta n)$  Dimensionless birefringence of the polymer;  $(g)$  flow rate gradient,  $\text{s}^{-1}$ ;  $(\eta_0)$  viscosity of the solvent,  $\text{cm s g}^{-1}$ ; and  $(c)$  concentration of the polymer.



**Fig. 2.** Variation of the birefringence  $\Delta n$  with the flow rate gradient  $g$  for the PMAA- $\text{Se}^0$ - $\text{H}_2\text{O}$  system.  $c_{\text{PMAA}}$ , wt %: (1) 0.1, (2) 0.071, (3) 0.049, (4) 0.0195, and (5) 0.0048.

## CONCLUSIONS

- (1) Reduction of sodium selenite with ascorbic acid in an aqueous solution of polymethacrylic acid yielded spherical nano- $\text{Se}^0$  particles with a size of ca. 15 nm.
- (2) Molecular-optical studies showed that the polymethacrylic acid macromolecules adsorbed on nano- $\text{Se}^0$  form an ultra-high-molecular-weight ( $\text{MW} = 150 \times 10^6$ ) spherical nanostructure with a size of ca. 90 nm, comprising over 4000 densely packed macromolecules.
- (3) The PMAA- $\text{Se}^0$ - $\text{H}_2\text{O}$  system exhibits an induced conformational transition of polyacrylic acid from a Gaussian globule to a sphere.
- (4) The hydrophobic interactions of the macromolecules with nano- $\text{Se}^0$  play the deciding role in their adsorption on nanoparticles.

## REFERENCES

1. Pomogailo, A.D., *Nanochastitsy metallov v polimerakh* (Nanoparticles of Metals in Polymers), Moscow: Nauka, 1999.

2. Mayer, A.B.R., *Polym. Adv. Technol.*, 2001, vol. 12, no. 1, pp. 96–106.
3. Nekrasova, T.N., Study of the Structure of Polymethacrylic Acid Molecules in Solution, *Cand. Sci. Dissertation*, Leningrad, 1970.
4. Kopeikin, V.V., Valueva, S.B., Kipper, A.I., *et al.*, *Vysokomol. Soedin., Ser. A*, 2003, vol. 45, no. 4, pp. 615–622.
5. Tsvetkov, V.N., Eskin, V.E., and Frenkel, S.Ya., *Struktura makromolekul v rastvorakh* (Structure of Macromolecules in Solutions), Moscow: Nauka, 1964.
6. Eskin, V.E., *Rasseyanie sveta rastvorami polimerov i svoistva makromolekul* (Light Scattering by Polymer Solutions and Properties of Macromolecules), Leningrad: Nauka, 1986.
7. Brown W., *Dynamic Light Scattering: the Method and Some Applications*, Oxford: Clarendon, 1993.
8. Meewes, M., Ricka, J., De Silva, M., *et al.*, *Macromolecules*, 1991, vol. 24, no. 21, pp. 5811–5816.
9. Hishio, I., Shao Thang Sun, Swislow, G., and Tanaka, T., *Nature*, 1979, vol. 281, no. 5728, pp. 208–209.
10. Konishi, T., Yoshizaki, T., and Yamakawa, H., *Macromolecules*, 1991, vol. 24, no. 20, pp. 5614–5622.
11. *Laser Light Scattering in Biochemistry*, Harding, S.E., Satelle, D.B., and Bloomfield, V.A., Eds., Cambridge: Royal Chem. Soc., Inform. Services, 1992, pp. 3–21.
12. Lyubina, S.Ya., Influence of the Intramolecular Interactions and Secondary Structure on the Dynamooptical Properties of Macromolecules in Solutions, *Cand. Sci. Dissertation*, Leningrad, 1965.
13. Storhoff, J.J., Lazaorides, A.A., Mucic, R.C., *et al.*, *J. Am. Chem. Soc.*, 2000, vol. 122, no. 19, pp. 4640–4650.