

## FARADAY COMMUNICATIONS

First Kinetic Measurements on  $\text{SiD}_2$ : Absolute Rate Constant for the Reaction  $\text{SiD}_2 + \text{H}_2$ 

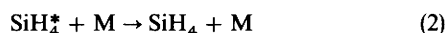
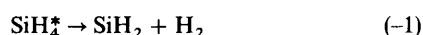
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$\text{SiD}_2$ , generated *via* the laser flash photolysis of phenyl[ $^2\text{H}_3$ ]silane, has been detected *via* a ro-vibrational transition in its  $\tilde{\text{A}}(^1\text{B}_1) \leftarrow \tilde{\text{X}}(^1\text{A}_1)$  absorption band. It has been monitored for the first time in time-resolved studies to give an absolute rate constant of  $(3.8 \pm 0.2) \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$  for the reaction  $\text{SiD}_2 + \text{H}_2$ . Product  $\text{SiH}_2$  from this reaction was also observed.

Silylene,  $\text{SiH}_2$ , is a key intermediate in the photolytic and thermal decompositions of silanes and plays an important role in the process for chemical vapour deposition of thin films of amorphous silicon. In 1985 Inoue and Suzuki<sup>1</sup> reported the first directly measured rate constants for reactions of  $\text{SiH}_2$ . Since then the rate constants for several other reactions have been obtained by time-resolved methods.

Kinetic studies of  $\text{SiD}_2$  have not previously been carried out although they are of particular interest in two respects, *viz.* (i) a comparison with  $\text{SiH}_2$  data will help elucidate reaction mechanisms and (ii) many  $\text{SiH}_2$  reactions exhibit pressure dependences often making it problematical to obtain high-pressure limiting rate constants. Isotopic scrambling effects make it possible to circumvent this problem with  $\text{SiD}_2$  studies. For example the reaction  $\text{SiH}_2 + \text{H}_2$  can be conveniently described by three elementary steps:



Pressure dependence of this reaction arises from the behaviour of the vibrationally excited molecular product,  $\text{SiH}_4^*$ , which can either revert back to reactants or be stabilized by a third body. For the analogous  $\text{SiD}_2$  reaction the vibrationally excited product is  $\text{SiD}_2\text{H}_2^*$  which can redissociate to  $\text{SiHD} + \text{HD}$  or  $\text{SiH}_2 + \text{D}_2$  as well as reverting to the reactants. Since  $\text{SiD}_2$  will only be reformed once in every six times on a statistical basis, following the argument of Jasinski,<sup>2</sup> there should be virtually no pressure effect on its rate of disappearance which will therefore closely approximate the rate of the true insertion reaction.

$\text{SiD}_2$  kinetic studies were carried out using the laser flash-laser absorption technique previously used to study  $\text{SiH}_2$  kinetics.<sup>3,4</sup> Transient concentrations of  $\text{SiD}_2$  were produced by 193 nm laser flash photolysis of  $\text{PhSiD}_3$  (synthesized by  $\text{LiAlD}_4$  deuteration of  $\text{PhSiCl}_3$ , and of isotopic purity >92.6%). The concentration of  $\text{SiD}_2$  was monitored by high-resolution absorption spectroscopy with light from a single-frequency ring-dye laser (linewidth  $3 \times 10^{-5} \text{ cm}^{-1}$ ). Prior to this a high-resolution spectrum of  $\text{SiD}_2$  was recorded in order to find a strong characteristic transition for probing. Previously recorded spectra,<sup>5–7</sup> although not sufficiently resolved to reveal a suitable transition, indicated that the  $^{\text{P}}\text{Q}_1$  bandhead of  $\text{SiD}_2$  occurred at  $17385 \text{ cm}^{-1}$  for the  $\tilde{\text{A}}(^1\text{B}_1, 0,3,0) \leftarrow \tilde{\text{X}}(^1\text{A}_1, 0,0,0)$  transitions. Based on this a spectral search was made over the wavenumber range  $17384$ – $17391 \text{ cm}^{-1}$  at intervals of  $0.03 \text{ cm}^{-1}$ , using a series of one-shot photolysis experiments with gas mixtures consisting

of 10 mTorr  $\text{PhSiD}_3$  in 5 Torr  $\text{SF}_6$ . The resulting spectrum is shown in Fig. 1, although we have not yet been able to assign the observed transitions. Preliminary calculations of the  $\text{SiD}_2$  spectrum confirm the existence of a number of strong ro-vibrational lines in this wavenumber region.<sup>8</sup> A check at  $17259.50 \text{ cm}^{-1}$  (the location of a strong  $\text{SiH}_2$  transition) suggested a small photolytic yield of  $\text{SiH}_2$  from  $\text{PhSiD}_3$  of *ca.* 10% of the intensity for the equivalent photolysis of  $\text{PhSiH}_3$ . In the latter photolysis we found a weak absorption due to  $\text{SiH}_2$  on probing at  $17387.07 \text{ cm}^{-1}$  (the location of the strongest  $\text{SiD}_2$  transition observed) which was *ca.* 16% of that from  $\text{PhSiD}_3$  under identical conditions. However these results indicate a less than 2% contribution to the signal at  $17387.07 \text{ cm}^{-1}$  from  $\text{SiH}_2$  formed from impurity (or by photolytic scrambling) from  $\text{PhSiD}_3$  and ensure that  $\text{SiD}_2$  can be monitored at  $17387.07 \text{ cm}^{-1}$  quantitatively to a good approximation.

For kinetic studies  $\text{SiD}_2$  was monitored at both  $17387.07$  and  $17385.18 \text{ cm}^{-1}$ . No detectable difference was observed between results at the two wavelengths. Reaction mixtures generally contained 4 mTorr of  $\text{PhSiD}_3$ , 0–3.2 Torr of hydrogen and added  $\text{SF}_6$  buffer gas (total pressure 5 Torr). The temperature was 294 K. First-order decays were averaged for four to ten laser shots at 50 to 140 mJ pulse<sup>-1</sup>. All decays were good-quality single exponentials. The results shown in Fig. 2, confirm that second-order kinetics were obeyed. The slope of this plot yielded a rate constant of  $(3.8 \pm 0.2) \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ , where the quoted error represents two standard deviations. Variation of the total pressure by a factor of 50 between 2 and 100 Torr indicated that the observed rate was independent of pressure

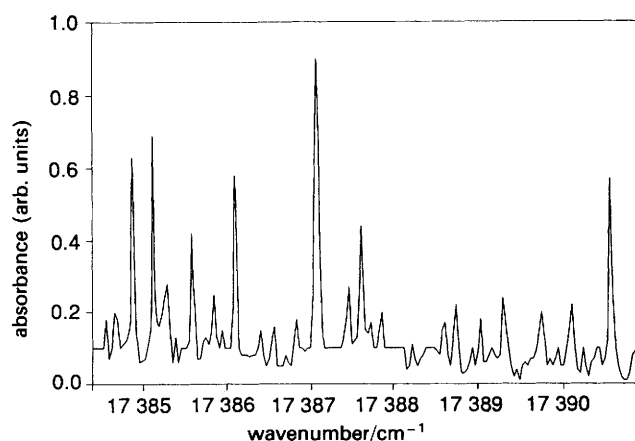


Fig. 1 Observed gas-phase absorption spectrum of  $\text{SiD}_2$

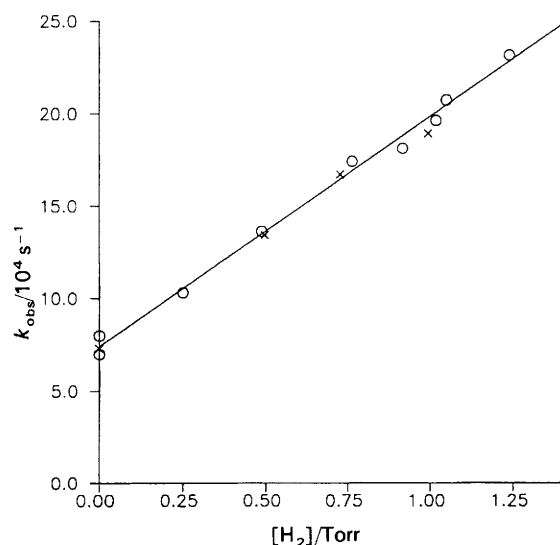


Fig. 2 Second-order plot for reaction of  $\text{SiD}_2$  with  $\text{H}_2$  at 294 K (total pressure = 5 Torr),  $\text{SiD}_2$  detection wavenumber: (○) 17387.07, (×) 17385.18  $\text{cm}^{-1}$ .

(within  $\pm 20\%$ ). A comparison of the rate constant reported here for  $\text{SiD}_2 + \text{H}_2$  with the high-pressure limiting rate constant of  $3.2 \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$  reported for  $\text{SiH}_2 + \text{H}_2$  by Jasinski and Chu<sup>9</sup> and the rate constants of  $(2.6 \pm 0.7) \times 10^{-12}$  and  $(1.9 \pm 0.2) \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$  for  $\text{SiH}_2 + \text{D}_2$  reported by Jasinski<sup>2</sup> and Baggott *et al.*<sup>4</sup>

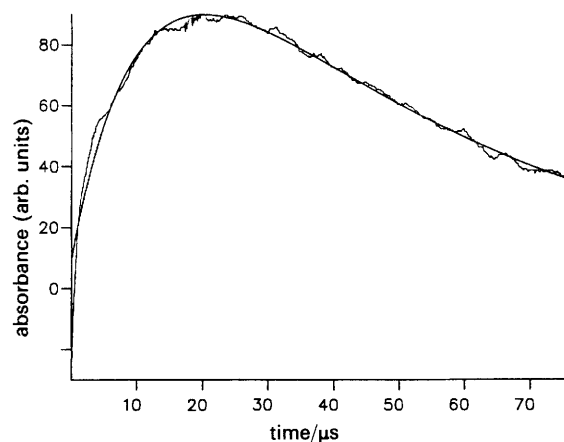


Fig. 3 Time evolution of trace for  $\text{SiH}_2$  from the reaction of  $\text{SiD}_2 + \text{H}_2$ .  $[\text{H}_2] = 1$  Torr. Smooth curve is biexponential best fit.

Table 1 Second-order rate constants found for reaction of  $\text{SiH}_2$  with  $\text{H}_2$

total pressure/Torr	$k/10^{-13} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$	
	this paper <sup>a</sup>	Jasinski and Chu <sup>b</sup>
3	$1.5 \pm 0.3$	2.7
5	$3.4 \pm 0.5$	3.9
7	$4.5 \pm 0.3$	5.0
10	$4.8 \pm 0.5$	6.5

<sup>a</sup> Total pressures in  $\text{H}_2$ . <sup>b</sup> Ref. 9. Total pressures in He.

respectively, show it to be of a reasonable magnitude. They also indicate the magnitude of the true insertion process isotope effects.

In addition to monitoring  $\text{SiD}_2$  concentration decay curves we have carried out additional experiments where the generation of  $\text{SiH}_2$  from isotopic scrambling and its subsequent disappearance by reaction with  $\text{H}_2$  were observed. Reaction mixtures containing 2 mTorr of  $\text{PhSiD}_3$  and made up to total pressures in hydrogen of 1–10 Torr were photolysed and  $\text{SiH}_2$  absorption was monitored at 17 259.50  $\text{cm}^{-1}$ . Averaged absorption growth–decay traces were obtained by summation of five traces using an excimer laser photolysis energy of 100 mJ pulse<sup>-1</sup>. An example of such a trace is shown in Fig. 3. Although the noise level in these traces was not excessive, the short rise-time (slightly perturbed by the laser pulse) made analysis of the growth part of the curves difficult. Despite this the results obtained from experiments with 1, 3 and 5 Torr of  $\text{H}_2$  gave an average growth rate constant of  $(4 \pm 2) \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ . Although the error limit is large this is the same as the value found by monitoring  $\text{SiD}_2$ . By contrast there was no time-resolution difficulty in fitting the decay part of the  $\text{SiH}_2$  traces and so a set of reasonably precise rate constants for reaction of  $\text{SiH}_2$  and  $\text{H}_2$  were obtained. These are shown in Table 1 and demonstrate the expected pressure dependence of this reaction. Comparison with the values of Jasinski and Chu<sup>9</sup> for the same reaction in He buffer gas shows reasonable agreement.

The rate constant obtained here for  $\text{SiD}_2 + \text{H}_2$  may be compared with that<sup>9</sup> for  $\text{SiH}_2 + \text{H}_2$ , viz.  $3.2 \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$  (high-pressure limit). There is thus an inverse isotope effect of ca. 0.84. This can be accounted for by transition-state theory.<sup>10</sup> In future publications we will present: (i) a temperature-dependent study of the reaction of  $\text{SiD}_2$  with  $\text{H}_2$  and (ii) a full experimental and theoretical study of the isotope effects in the  $\text{SiH}_2$ – $\text{SiD}_2 + \text{H}_2$ – $\text{D}_2$  reaction systems.

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