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LOW-TEMPERATURE PHASE TRANSITIONS AND MAGNETIC SUSCEPTIBILITIES  
OF LOW-DIMENSIONAL URANIUM FLUORIDES

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Measurements have been made on the magnetic susceptibilities of the quasi-one-dimensional magnetic  $\alpha$ -UF<sub>5</sub> and the cluster magnetic U<sub>2</sub>F<sub>9</sub> (a four-nuclear type) over the range 4.2-300 K. The Curie-Weiss law applies for  $\alpha$ -UF<sub>5</sub> above 200 K, with  $\mu_{\text{ef}} = 2.37\mu_{\text{B}}$ ,  $\theta = -148$  K; that law also applies to U<sub>2</sub>F<sub>9</sub> above 150 K, with  $\mu_{\text{ef}} = 2.9\mu_{\text{B}}$  and  $\theta = 118$  K. The antiferromagnetic ordering is discussed for the one-dimensional chains in  $\alpha$ -UF<sub>5</sub> together with the phase transition observed in U<sub>2</sub>F<sub>9</sub> at 6 K. It is assumed that the U<sup>4+</sup> and U<sup>5+</sup> valency states are stabilized below the phase-transition point, and only then can U<sub>2</sub>F<sub>9</sub> be considered as the compound UF<sub>4</sub>·UF<sub>5</sub>.

Spin cluster compounds and quasi-one-dimensional and quasi-two-dimensional ones are familiar for the d elements. Out of compounds of the F elements, only the uranium fluorides  $\alpha$ -UF<sub>5</sub> and U<sub>2</sub>F<sub>9</sub> have structures corresponding to quasi-one-dimensional and cluster magnetics. The magnetic parameters of these have not been examined. In general, the magnetochemical parameters of uranium compounds have been inadequately examined. In particular, among the fluorides, we have data only for UF<sub>4</sub> [1] and  $\beta$ -UF<sub>5</sub> [2]. Here we present results on the magnetic susceptibilities of the major representatives of the higher paramagnetic uranium fluorides  $\alpha$ -UF<sub>5</sub> and U<sub>2</sub>F<sub>9</sub> over the range 4.2-300 K.

The  $\alpha$ -UF<sub>5</sub> was made by direct fluorination of uranium tetrafluoride with elemental fluorine at 420 K, while U<sub>2</sub>F<sub>9</sub> was made by reacting UF<sub>4</sub> with uranium hexafluoride at 520 K. The initial UF<sub>4</sub> was made by reducing uranium hexafluoride with silicon at 590 K. The fluorine was purified from oxygen by rectification [3]. The  $\alpha$ -UF<sub>5</sub> and U<sub>2</sub>F<sub>9</sub> were identified by x-ray analysis.

The measurements were made by Faraday's methods over the range 4.2-300 K. A specimen of mass 60 mg was sealed into a silica tube. The empty tube was measured separately subsequently. The value of  $\chi$  for this was incorporated in calculating the susceptibility. The diamagnetic correction was derived from an additive scheme via the data of [4]. The measurements were made in fields of 0-10.9 kOe. There was a weak  $\chi(H)$  dependence evidently caused by traces of ferromagnetic materials, in particular iron or nickel. Figure 1 shows the observed temperature dependence of the susceptibility for  $\alpha$ -UF<sub>5</sub>;  $\chi$  behaves nonlinearly below 200 K, while above it the Curie-Weiss law applies:

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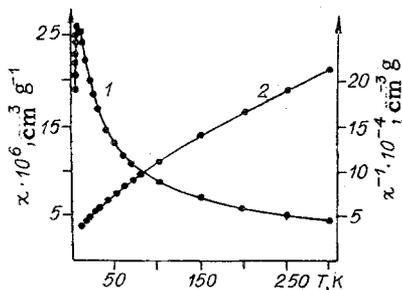


Fig. 1

Fig. 1. Dependence of the specific susceptibility  $\chi$  (curve 1) and  $\chi^{-1}$  (curve 2) on temperature for  $\alpha\text{-UF}_5$ .

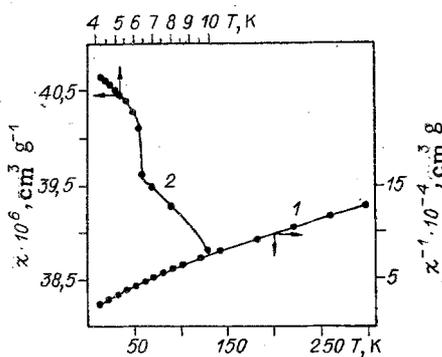


Fig. 2

Fig. 2. Temperature dependence of the specific magnetic susceptibility for  $\text{U}_2\text{F}_9$ : 1)  $\chi^{-1}(T) = 10\text{-}300$  K; 2)  $\chi(T) = 4.2\text{-}10$  K.

$$\chi = \frac{C}{T - \theta}$$

with parameters  $\mu_{\text{ef}} = \sqrt{8C} = 2.37\mu\text{B}$  and  $\theta = -148$  K.

The susceptibility does not follow the Curie-Weiss law at low temperatures, and there is a peak in  $\chi$  at 6.3 K, which may be related to antiferromagnetic ordering in the one-dimensional chains. The contribution from the orbital moment is not small for  $\text{U}^{5+}$  (terms  ${}^2\text{F}_{5/2}$  and  ${}^2\text{F}_{7/2}$ ), which can lead to substantial anisotropy in the exchange interaction. The exact solution for a one-dimensional Ising chain is [5]

$$\chi = \frac{N\bar{\mu}^2}{kT} \theta^{2J/kT},$$

$$\bar{\mu} = g\beta,$$

where  $g$  is the factor and  $\beta$  is the Bohr magneton.

The observed  $\chi$  are compared with the theoretical values for 4.2-20 K to derive the optimum parameters  $J = -2.84$  cm $^{-1}$  and  $g = 0.718$ . This  $g$  factor differs from the pure spin value, which indicates that a pure Ising chain model is not applicable for  $\alpha\text{-UF}_5$ . Susceptibility calculations for the linear Heisenberg chain led Bonner and Fisher [6] to estimate the exchange interaction parameter ( $S = 1/2$ ) for  $\text{U}^{5+}$  ions. The theory gives the following relations for an antiferromagnetic chain:

$$|J|\chi_{\text{max}}/Ng^2\beta^2 = 0.07346; \quad kT(\chi_{\text{max}})/|J| = 1.282.$$

On using the observed  $\chi = 8685 \cdot 10^{-6}$  cm $^3 \cdot \text{mole}^{-1}$  and  $T(\chi_{\text{max}}) = 6.3$  K, we get  $|J| = 4.189$  cm $^{-1}$  and  $g = 1.376$ .

Therefore, neither the Heisenberg one-dimensional chains nor the Ising ones give a satisfactory description of the observed behavior of the susceptibility of  $\alpha\text{-UF}_5$ , which may indicate a substantial contribution from chain interaction. The theoretical  $\chi(T)$  curve [6] for the range 4.2-20 K describes the observed result satisfactorily; the observed and theoretical values of  $\chi$  deviate substantially at  $T > 20$  K. The nonlinear behavior of  $\chi^{-1}(T)$  below 100 K is due evidently to the joint action of the octahedral field and the spin-orbit interaction, which displaces the  ${}^2\text{F}_{5/2}$  and  ${}^2\text{F}_{7/2}$  multiplets, and this is also responsible for  $\theta$  being larger than the temperature of antiferromagnetic ordering. Only combined incorporation of the exchange interactions and the splittings for  $\text{U}^{5+}$  can explain the results. The effects from excited states lead to an additional temperature-independent (van Vleck) contribution to the susceptibility. We have incorporated the nonlinearity of  $\chi^{-1}(T)$  by introducing this term via the formula

$$\chi = \frac{C}{T - \theta} + \chi_{\text{VV}}.$$

The best agreement with experiment is attained with  $C = 0.454$ ,  $\chi_{\text{VV}} = 321 \cdot 10^{-6}$  cm $^3 \cdot \text{mole}^{-1}$ ,  $\mu_{\text{ef}} = 1.914$ , and  $\theta = -70$  K.

Figure 2 gives the observed susceptibility of  $U_2F_9$ ; the structure contains four-nuclear  $U_4F_{18}$  clusters, in which the two  $U^{4+}$  atoms and the two  $U^{5+}$  ones are localized at the structurally equivalent vertices of a regular tetrahedron. Strictly speaking, this compound has fractional valency at ordinary temperatures, while the black color is due to the rapid  $f^1 \rightleftharpoons f^2$  electron migration. The six bridge fluorine atoms lie near the edges of the tetrahedron, while the other 12 fluorine atoms supplement the coordination polyhedra for each of the uranium atoms to octahedra. There is close packing of the  $U_4F_{18}$  particles. In such a four-nucleus cluster consisting of ions of variable valency,  $U^{4+}$  and  $U^{5+}$ , which have configurations  $f^1$  and  $f^2$ , one expects to find spin ordering and unpaired-electron localization at low temperatures [6, 7, 9]. No signs of magnetic ordering have been observed throughout the temperature range used (curve 1 in Fig. 2), which may mean either that the exchange constant is very small or that spin localization effects are too strong [8]. There is a step change in the susceptibility at 6 K (curve 2 in Fig. 2). One assumes that the change is due to a structural phase transition on account of localization and ordering for the  $f^1$  and  $f^2$  electrons. The localization reduces the initial high symmetry of the crystal and stabilizes the  $U^{4+}$  and  $U^{5+}$  in two structurally nonequivalent positions.

The general argument can be given on the trend in  $\chi(T)$  in this transition (neglecting the effects of structure change in the environment of the uranium ion). The unpaired-electron localization on the  $U^{4+}$  and  $U^{5+}$  in the ordered phase should increase  $\chi$  by comparison with the high-temperature phase, in which the electrons migrate from ion to ion. The migrating electrons should be characterized by the Pauli component of the susceptibility  $\chi_{\text{Pauli}}$ , which is less by two orders of magnitude than the susceptibility for localized electrons. In fact, the observed  $\chi(T)$  curve shows a step with reducing  $\chi$  on transition to the high-temperature phase. The size of the step enables one to estimate the proportion of the delocalized electron state. The contribution to  $\chi$  from these electrons is described by Pauli's formula [4]

$$\chi_{\text{Pauli}} = \frac{3N\beta^2}{\epsilon_F},$$

where  $\epsilon_F$  is the Fermi energy;  $N$ , Avogadro's number; and  $\beta$ , the Bohr magneton. Substitution shows that  $\chi_{\text{Pauli}}$  is less by about two orders of magnitude than the contributions to  $\chi$  from electrons localized on the paramagnetic ions. We neglect  $\chi_{\text{Pauli}}$  to get the ratio of the susceptibilities on the two sides of the step:

$$\frac{\chi_l}{\chi_h} = \frac{\mu_1^2 + \mu_2^2}{\mu_1^2(1+x) + \mu_2^2(1-x)}$$

where  $\mu_1$  is the magnetic moment of  $U^{5+}$  derived here;  $\mu_2$  is the magnetic moment of  $U^{4+}$  [10]; and  $x$  is the proportion of delocalized electrons. The observed ratio is  $\chi_l/\chi_h = 1.02$ . The with  $\mu_1 = 2.37$  and  $\mu_2 = 3.58$ , we get  $x = 0.05$ , i.e., the unpaired electrons are 5% delocalized in the high-temperature phase.

The Curie-Weiss law applies for  $\chi(T)$  above 153 K, with  $\mu_{\text{ef}} = 2.9\mu_B$  and  $\theta = -118$  K.

We can thus say that  $U_2F_9$  has fractional valency only above 6.5 K; at lower temperatures, it can be considered as  $UF_4 \cdot UF_5$  and as having a structure close to that of the initial compound  $U_2F_9$ . One also assumes that the phase transition causes the electron localization on the  $U^{4+}$  and  $U^{5+}$  to be ordered, so the symmetry is reduced. As the point-group symmetry of the  $U_2F_9$  molecule is  $C_2$  with an ordered disposition for the  $U^{4+}$  and  $U^{5+}$  ions, the structure of  $U_2F_9$  at low temperatures can only be of tetragonal symmetry with fourfold axes or a screw  $4_2$  axis or a mirror-rotation  $\bar{4}$  one. As the density should increase on charge ordering, the  $4_2$  symmetry is probably preferable, in which there is alternation of  $U^{4+}$  and  $U^{5+}$  along the  $C_4$  axis. With that packing lacking a center of symmetry, the dipole moments of the individual molecules are parallel, which can lead to some macroscopic electrical parameters (piezoelectric or ferroelectric ones).

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<sup>35</sup>Cl NQR SPECTRA OF COMPLEXES OF TETRACHLOROSTANNANE  
WITH SUBSTITUTED BENZOYL CHLORIDES

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<sup>35</sup>Cl NQR spectra of mixtures of SnCl<sub>4</sub> with 2-, 3-, and 4-XC<sub>6</sub>H<sub>4</sub>COCl were obtained. The electronic and steric structure of the complexes obtained was established. Their formation and structure depend on the nature of substituents X. Complexes with a trigonal-bipyramidal structure are formed with the participation of the carbonyl oxygen atom of the ligand as an electron-donor center, and complexes with trans-octahedral structure, with the participation of the sulfur atom of the X = CH<sub>3</sub>S substituent.

In the reaction of SnCl<sub>4</sub> with para-substituted benzoyl chlorides, complexes with a 1:1 composition are formed, in which the tin atom is pentacoordinated [1]. We obtained a crystalline complex with this structure in the reaction of ortho-methoxybenzoyl chloride with tetrachlorostannane [2]. In the course of time, it converts into a cis-octahedral complex [2]. In the continuation of a <sup>35</sup>Cl NQR study of complexes of tetrachlorostannane with organic ligands, we obtained <sup>35</sup>Cl NQR spectra of mixtures of SnCl<sub>4</sub> with 2-, 3-, or 4-substituted benzoyl chlorides (Tables 1 and 2).

The NQR spectra of equimolar mixtures of 2-XC<sub>6</sub>H<sub>4</sub>COCl (X = F, ClCO) and 3-XC<sub>6</sub>H<sub>4</sub>COCl (X = O<sub>2</sub>N, ClCO) consist of the spectra of initial components (Table 1). The same NQR spectra were obtained for these values of X (except for X = F) for mixtures of SnCl<sub>4</sub> with 4-XC<sub>6</sub>H<sub>4</sub>COCl [1]. This indicates the absence of complexation.

A <sup>35</sup>Cl NQR spectrum of an equimolar mixture of SnCl<sub>4</sub> and 2-ClC<sub>6</sub>H<sub>4</sub>COCl consists of four lines with different width and intensity. The broad and intense lines in the spectrum are due to superposition of several closely positioned lines. Among these lines, the lowest frequency lies in the same region as most of the spectrum of individual SnCl<sub>4</sub>. The high-frequency broad and intense line in the spectrum belongs to the chlorine atom in the aromatic ring. The frequency of its maximum (36.88 MHz) practically coincides with that of the line of the same chlorine atom in the individual 2-ClC<sub>6</sub>H<sub>4</sub>COCl. In the NQR spectrum of a mixture of SnCl<sub>4</sub> with 2-ClC<sub>6</sub>H<sub>4</sub>COCl only one weak single line (29.666 MHz) is observed for the chlorine atom of the COCl grouping. Its frequency practically coincides with that of one of the lines for this chlorine atom in the spectrum of individual 2-ClC<sub>6</sub>H<sub>4</sub>COCl (see Table 1). The NQR spectrum obtained of a mixture of SnCl<sub>4</sub> and 2-ClC<sub>6</sub>H<sub>4</sub>COCl shows that no complexation took place.

In a quadruplet NQR spectrum of an equimolar mixture of SnCl<sub>4</sub> with 2-CH<sub>3</sub>C<sub>6</sub>H<sub>4</sub>COCl, the highest frequency line belongs to the chlorine atom in the COCl grouping, and the three low-

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