# The vibrational spectra of some dialkylamido derivatives of phosphorus and arsenic 

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#### Abstract

Infrared and Raman spectra have been obtained and assigned for $\mathrm{X}\left(\mathrm{NR}_{2}\right)_{3}$. where $\mathrm{X}=\mathrm{P}$ or As. $\mathbf{R}=\mathrm{Me}$ or Et . The internal alkyl group modes could be analysed using a local symmetry of $C_{s}$ for an isolated $\mathrm{NR}_{2}$ unit. Skeletal modes were most consistent with an overall molecular symmetry of $C_{3}$.


## INTRODUCTION

Compounds containing coordinated dialkylamide groups have not been studied very frequently by vibrational spectroscopists. A review by Bradley [1] summarised data available on the vibrational spectra of such compounds of metals up to 1971. Dimethyl-amidodichloro- and -difluorophosphines have been examined $[2,3]$, the data on the latter being much more complete [3]. A trans-geometry, of $C_{s}$ symmetry, sufficed to explain the observed spectrum of $\left(\mathrm{Me}_{2} \mathrm{~N}\right)$ $\mathrm{PF}_{2}$, and no evidence was found for the presence of rotational isomers. The arsenic-chloro analogue ( $\left.\mathrm{Me}_{2} \mathrm{~N}\right) \mathrm{AsCl}_{2}$, howẹver, exists as a mixture of two isomers, trans ( $C_{s}$ symmetry) and gauche ( $C_{1}$ symmetry), even at 80 K [4].

The only published work on $\mathrm{P}\left(\mathrm{NMe}_{2}\right)_{3}$ has been that of Goubeau et al. [5], who have also examined XP$\left(\mathrm{NMe}_{2}\right)_{2} \mathrm{Cl}, ~ \mathrm{XP}\left(\mathrm{NMe}_{2}\right) \mathrm{Cl}_{2}$ and $\mathrm{XP}\left(\mathrm{NMe}_{2}\right)_{3-n} \mathrm{Me}_{n}$ ( $\mathrm{X}=\mathrm{O}$ or $\mathrm{S} ; n=0,1,2$ or 3 ) [6, 7]. For $\mathrm{P}\left(\mathrm{NMe}_{2}\right)_{3}$, it was suggested that the highest molecular symmetry compatible with the observations was $C_{s}$, although it is impossible to assess the validity of their conclusions, as complete spectral data were not listed and no Raman polarizations were reported.

A gas-phase electron diffraction study of $\mathrm{P}\left(\mathrm{NMe}_{2}\right)_{3}$. together with $\mathrm{P}\left[\mathrm{N}\left(\mathrm{CH}_{2}\right)_{2}\right]_{3}$. on the other hand [8] indicates that the tris(dimethylamido)phosphine does possess three-fold symmetry.

In order to resolve this ambiguity, we have recorded the i.r. and Raman spectra of $\mathrm{P}\left(\mathrm{NR}_{2}\right)_{3}$ and $\mathrm{As}\left(\mathrm{NR}_{2}\right)_{3}$, where $\mathrm{R}=\mathrm{Me}$ or Et . and assigned as many bands as possible to normal vibrational modes.

## EXPERIMENTAL

All preparations were carried out under an atmosphere of pure dry argon. Dimethylamine and diethylamine were dried by distillation from sodium wire, and the ether solvent was dried using $\mathrm{LiAlH}_{4}$. The general method for preparing the four dialkylamide-compounds was based on those of BURG and Slota [9] and Moedritzer [10]. A six-fold molar excess of the appropriate amine was added to an ethereal solution of the metal trichloride. with stirring. at -10 C . After warming to room temperature the precipitated amine hydrochloride was filtered off under vacuum. The ether was removed using a rotary evaporator, and the tris(dialkyl-
amide) isolated by distillation in vacuo. Satisfactory $\mathbf{C}, \mathrm{H}$ and N analyses were obtained in all cases.

Infrared spectra were obtained using a Perkin-Elmer 521 spectrometer ( $4000-250 \mathrm{~cm}^{-1}$ ). Solid-. liquid- and vapourphase samples were used (for dimethylamido compounds; liquid-phase only for those containing diethylamido groups). The spectra were calibrated using known wavenumbers of $\mathrm{CH}_{4} . \mathrm{HBr} . \mathrm{CO}, \mathrm{NH}_{3}$ and $\mathrm{H}_{2} \mathrm{O}$. All of the observed wavenumbers are accurate to $\pm 2 \mathrm{~cm}^{-1}\left( \pm 5 \mathrm{~cm}^{-1}\right.$ for very weak or broad features.)

A Cary 81 spectrometer. modified by Anaspec Ltd., was used to record the Raman spectra. the excitation source being a Spectra-Physics 164 Argon ion laser (power output ca. 1 W at both 488.0 and 514.5 nm ). Liquid samples were distilled directly into glass capillaries (approx. 1 mm i.d.). and polarization measurements carried out by examining the spectrum with the incident light successively parallel and perpendicular to the axis of a polaroid analyser. The depolarization values so obtained were proportional to the true values.

## RESULTS

The i.r. and Raman spectra of $\mathrm{P}\left(\mathrm{NMe}_{2}\right)_{3}, \mathrm{P}\left(\mathrm{NEt}_{2}\right)_{3}$, $\mathrm{As}\left(\mathrm{NMe}_{2}\right)_{3}$ and $\mathrm{As}\left(\mathrm{NEt}_{2}\right)_{3}$ are listed in Table 1-4. The proposed assignments are summarised below.

## DISCUSSION

It will be convenient to divide this section as follows: (a) internal modes of the $\mathrm{NMe}_{2}$ group, (b) internal modes of the $\mathrm{NEt}_{2}$ group (c) skeletal modes of $\mathrm{P}\left(\mathrm{NC}_{2}\right)_{3}$ and (d) skeletal modes of $\mathrm{As}\left(\mathrm{NC}_{2}\right)_{3}$. In sections (a) and (b). data from both the P and $\mathrm{A} s$ compound will be included.

## (a) Internal $\mathrm{NMe}_{2}$ modes

As assignment of the fundamental modes of dimethylamine is available [11. 12]. which may be used as a basis for this assignment. The electron diffraction data of Vilkov et al. [ 8$]$. suggest that the $\mathrm{NMe}_{2}$ unit in $\mathrm{P}\left(\mathrm{NMe}_{2}\right)_{3}$ has a planar skeleton (local symmetry $C_{2 v}$ ). It is known [13], however, that pyramidal $\mathrm{NX}_{3}$ units have a very low barrier to inversion at the N atom (in the absence of restraining factors such as ring formation, and therefore it is likely that this apparent planarity is due to rapid inversion. We will assume $\mathrm{XNMe}_{2}$ ( $\mathrm{X}=\mathrm{P}$ or As) to be pyramidal. of local symmetry $C_{s}$. The numbers and symmetry types of the

Table 1. The vibrational spectrum of tris(dimethylamido)phosphine (all numbers in wavenumbers $/ \mathrm{cm}^{-1}$ )

| Solid | Infrared liquid | Vapour | Raman liquid | Proposed assignment |
| :---: | :---: | :---: | :---: | :---: |
| 2997 m | $3000 \mathrm{sh}^{*}$ | $3000 \mathrm{sh}^{\text {r }}$ |  |  |
| 2965 w |  | 2974 ms | 2974, sh', dp | $\mathrm{CH}_{3}$ stretch ( $\mathrm{A}^{\prime \prime}$ ) |
| 2930 s | 2940 vs | 2940 vs | 2931 ms dp | $\mathrm{CH}_{3}$ stretch ( $\mathrm{A}^{*}$ ) |
| 2870 vs | 2890 s | 2888 s | 2889 ms dp | $\mathrm{CH}_{3}$ stretch ( $\mathrm{A}^{\prime \prime}$ ) |
| 2785 s | 2798 ms | 2798 s | 2797 s pol | $\mathrm{CH}_{3}$ stretch ( $A^{\prime}$ ) |
| 2485 m | $\begin{aligned} & 2470 \mathrm{~m} \\ & 2368 \mathrm{~ms} \end{aligned}$ |  |  |  |
|  |  | 2340 w |  |  |
|  | 2150 w |  |  |  |
|  |  | 2070 w |  |  |
| 1648 m | 1658 w | 1658 w | 1630 w |  |
|  | 1481 ms | $1486 \mathrm{sh}^{\prime}$ | 1481 s. dp | $\mathrm{CH}_{3}$ deformation ( $A^{*}$ ) |
| $1465 \mathrm{sh}^{\text {r }}$ | 1463 ms | 1465 ms | 1465 ms | $\mathrm{CH}_{3}$ deformation (A') |
| 1454 s | 1458 ms | 1444 sh $^{*}$ |  | $\mathrm{CH}_{3}$ deformation ( $A^{\prime}$ ) |
|  |  |  | $1441 \mathrm{~s} . \mathrm{dp}$ | $\mathrm{CH}_{3}$ deformation ( $A^{*}$ ) |
| 1419 w1406 w | $1417 \mathrm{sh}^{*}$ |  | $1417 \mathrm{~m} . \mathrm{dp}$ | $\mathrm{CH}_{3}$ deformation ( $A^{\prime \prime}$ ) |
|  |  | 1404 w |  |  |
| 1364 w | 1376 vw |  |  |  |
| 1356 w |  |  |  |  |
| 1322 m | 1318 s |  |  | $\mathrm{CH}_{3}$ deformation ( $A^{\prime}$ ) |
| $1295 \mathrm{sh}^{\prime}$ |  |  |  |  |
| 1278 s |  | 1277 ms | 1273 m | $1280 \mathrm{mw}, \mathrm{dp}$ | $\mathrm{CH}_{3}$ rock $\left(A^{\prime}\right)$ |
| 1246 m | 1248 w |  |  |  |
| 1190 ) ${ }^{\text {vs }}$ | 1191 s | 1204 ms | 1195 vw, dp | $\mathrm{CH}_{3}$ rock ( $A^{\prime \prime}$ ) |
|  | $1178 \mathrm{sh}^{\mathbf{r}}$ |  |  |  |
| 1164 s |  | $\left.\begin{array}{l} 1166 \\ 1159 \\ 1148 \end{array}\right\}$ |  |  |
| 1150 3 | $\begin{aligned} & 1150 \mathrm{sh}^{r} \\ & 1123 \mathrm{vw} \end{aligned}$ |  | 1147 w, poi | $\mathrm{CH}_{3} \operatorname{rock}\left(A^{\prime}\right)$ |
|  |  | 1148 |  |  |
| 1105 m | $1104 \mathrm{sh}^{\text {r }}$ |  |  |  |
|  | 1086 m |  |  | $\mathrm{CH}_{3}$ rock ( $\mathrm{A}^{\prime}$ ) |
| 1068 s | 1069 m |  |  | $\mathrm{NC}_{2}$ stretch (E) |
| $1060 \mathrm{sh}^{\prime}$ |  | 1062 mw |  |  |
| $1034 \mathrm{mw}$ |  | 1040 w |  |  |
|  | $1016 \mathrm{sh}^{\text {r }}$ | 1024 w |  |  |
| 988 vs | 984 vs | 978 sh" | $981 \mathrm{~ms}, \mathrm{dp}$ | $\mathrm{NC}_{2}$ stretch (E) |
| 958 vvs | 962 vs | 960 s | $963 \mathrm{sh}^{\text {r }}$, pol | $\mathrm{NC}_{2}$ stretch (A) |
| 946 s | 947 sh ${ }^{\text {r }}$ 904 shr $^{\text {r }}$ | $945 \mathrm{sh}^{\prime}$ |  | $\mathrm{NC}_{2}$ stretch ( $A$ ) |
| $\begin{aligned} & 895 \mathrm{~m} \\ & 807 \mathrm{w} \end{aligned}$ |  |  |  |  |
|  |  |  |  |  |  |  |  |
|  |  |  |  |  |
|  | 738 w | $\begin{aligned} & 745 \mathrm{~m} \\ & 720 \mathrm{~m} \end{aligned}$ |  |  |
| 706 \$ | 704 mw | 705 sh | $697 \mathrm{sh}^{\text {r }}$ | $\mathrm{PN}_{3}$ stretch (E) |
| 675 s | 680 s | 675 s | 676 vs. pol | $\mathrm{PN}_{3}$ stretch (A) |
| 654 s | 659 m | $652 \mathrm{sh}^{\text {r }}$ | $647 \mathrm{sh}^{\text {r }}$ |  |
| 570 w | 550 w |  |  |  |
| 506 w | 510 w |  |  |  |
| 491 m | 494 mw |  | 499 w | PNC deformation ( $E$ ) |
| 417 s | 420 m |  | 416 w | PNC deformation (E) |
|  | 405 mw |  |  | $\begin{aligned} & \mathrm{NC}_{2} \text { deformation (A } \\ & \text { or } E \text { ) } \end{aligned}$ |
| 389 s | 392 mw |  | 390 w | $\mathrm{NC}_{2}$ deformation (A or $E$ ) |
|  | 384 w |  |  |  |
|  | 360 vw |  |  |  |
| 340 mw | 345 mm |  |  |  |
| 333 mw | 332 ww |  | 337 s. pol | PNC deformation ( $A$ ) |
|  | 310 m |  |  |  |
|  |  |  | 295 s , pol | $\mathrm{CH}_{3}$ torsion ( $\mathrm{A}^{\prime}$ ) |
|  |  |  | 193 w. dp | $\mathrm{NC}_{2}$ torsion (E) |
|  |  |  | 127 mw , dp | $\mathrm{PN}_{3}$ deformation ( $E$ ) |

Table 2. The vibrational spectrum of tris(dimethylamido) arsine (all numbers in wavenumbers $/ \mathrm{cm}^{-1}$ )

| Solid | Infrared |  | Raman liquid | Proposed assignment |
| :---: | :---: | :---: | :---: | :---: |
|  | Liquid | Vapour |  |  |
| $\begin{aligned} & 2965 \mathrm{~s} \\ & 2955 \mathrm{~s} \end{aligned}$ | $2993 \mathrm{sh}^{-}$ | $2998 \mathrm{sh}^{\text {r }}$ |  |  |
|  | 2963 sh ${ }^{\text {r }}$ | $2964 \mathrm{sh}^{\text {r }}$ | 2965 ms, dp | $\mathrm{CH}_{3}$ stretch ( $\mathrm{A}^{\prime \prime}$ ) |
|  |  |  |  |  |
| 2860 ms | 2932 ms | 2932 ms | 2927 ms, dp | $\mathrm{CH}_{3}$ stretch ( $\mathrm{A}^{*}$ ) |
|  | 2886 ms | 2888 ms |  | $\mathrm{CH}_{3}$ stretch ( $\mathrm{A}^{*}$ ) |
|  | 2838 m | $2840 \mathrm{sh}^{\text {r }}$ | 2839 ms, pol | $\mathrm{CH}_{3}$ stretch ( $A^{\prime}$ ) |
| 2785 s | 2798 ms | 2798 ms | 2794 ms, pol | $\mathrm{CH}_{3}$ stretch ( $\mathrm{A}^{\prime}$ ) |
| 2450 s | $\begin{aligned} & 2440 \mathrm{w} \\ & 1830 \mathrm{w} \end{aligned}$ | 2450 w |  |  |
| 1610 m |  |  |  |  |



Table 3. The vibrational spectrum of tris(diethylamido)phosphine (all figures in wavenumbers $/ \mathrm{cm}^{-1}$ )

| Infrared <br> Liquid | Raman <br> Liquid | Proposed Assignment |
| :--- | :--- | :--- |

$522 \mathrm{w} \quad 525 \mathrm{vvw}, \mathrm{dp}$
508 vvw, dp
$495 \mathrm{w} \quad 492 \mathrm{vw}, \mathrm{pol}$
$460 \mathrm{w} \quad 462 \mathrm{vw}$, pol
420 w
$392 \mathrm{w} \quad 387 \mathrm{~m}, \mathrm{pol}$
$311 \mathrm{~s}, \mathrm{pol}$
PNC deformation ( $A$ )
?PNC deformation
$\mathrm{NC}_{2}$ deformation ( $A$ )
$\mathrm{CH}_{3}$ torsion ( $A^{\prime}$ )
N.B. i.p. and o.o.p. are abbreviations for in-phase, out-ofphase respectively.

Table 4. The vibrational spectrum of tris(diethylamido)arsine (all figures in wavenumbers $/ \mathrm{cm}^{-1}$ )

| Infrared liquid | Raman liquid | Proposed assignment |
| :---: | :---: | :---: |
| 2973 s | 2975 m | $\mathrm{CH}_{3}$ stretch ( $A^{\prime \prime}$ ) |
| 2968 s | 2965 m, dp | $\mathrm{CH}_{3}$ stretch ( $A^{\prime \prime}$ ) |
| 2933 s | 2931 s. pol | $\mathrm{CH}_{2}$ stretch (i.p.) ( $A^{\prime}$ ) |
|  | 2901 shr, dp | $\mathrm{CH}_{2}$ stretch (o.o.p) ( $A^{* *}$ ) |
|  | 2885 shr, dp | $\mathrm{CH}_{3}$ stretch ( $A^{\prime \prime}$ ) |
| 2873 s | 2871 s , pol | $\mathrm{CH}_{2}$ stretch (i.p.) ( $A^{\prime}$ ) |
|  | 2863 s | $\mathrm{CH}_{2}$ stretch (0.0.p) ( $A^{\prime \prime}$ ) |
| 2843 s |  | $\mathrm{CH}_{3}$ stretch ( $A^{\prime}$ ) |
| 2783 m | 2781 w | $\mathrm{CH}_{3}$ stretch ( $A^{\prime}$ ) |
| 2753 shr |  | ? $\mathrm{CH}_{3}$ stretch ( $A^{\prime}$ ) |
| 2728 shr |  |  |
| 2714 w | 2719 vw |  |
| 2483 w |  |  |
| 2391 w |  |  |
| 1595 w |  |  |
| 1485 shr | 1480 shr | $\mathrm{CH}_{3}$ deformation ( $A^{\prime \prime}$ ) |
| 1462 m | 1458 shr | $\mathrm{CH}_{3}$ deformation ( $A^{\prime}$ ) |
| 1452 m | 1452 s , dp | $\mathrm{CH}_{3}$ deformation ( $A^{\prime \prime}$ ) |
| 1394 shr | 1398 w | $\mathrm{CH}_{3}$ deformation ( $A^{\prime}$ ) |
| 1373 m | 1371 w, dp | $\mathrm{CH}_{2}$ o.o.p. scissors ( $A^{*}$ ) |
| 1362 shr | 1364 w | $\mathrm{CH}_{2}$ i.p. scissors ( $\mathrm{A}^{\prime}$ ) |
| 1344 w | 1343 dp | $\mathrm{CH}_{3}$ deformation ( $A^{\prime \prime}$ ) |
|  | 1323 m | $\mathrm{CH}_{3}$ deformation ( $A^{\prime \prime}$ ) |
| 1292 m | $1292 \mathrm{~m}, \mathrm{dp}$ | $\mathrm{CH}_{3}$ rock ( $A^{\prime \prime}$ ) |
| 1190 m | 1194 w. pol | $\mathrm{CH}_{2}$ i.p. wag ( $A^{\prime}$ ) |
| 1160 m | 1161 w | $\mathrm{CH}_{2}$ 0.0.p. wag ( $A^{*}$ ) |
| 1102 w | 1099 w, dp | $\mathrm{CH}_{3}$ rock (? $A^{\prime}$ ) |
|  | 1077 s, pol | $\mathrm{CH}_{3}$ rock ( $A^{\prime}$ ) |
| 1052 m |  | $\mathrm{CH}_{2}$ i.p. twist ( $A^{\prime}$ ) |
|  | 1055 w, dp | $\mathrm{CH}_{2}$ o.o.p. twist ( $A^{\prime \prime}$ ) |
|  | 1048 w | $\mathrm{CH}_{3}$ rock ( $A^{\prime \prime}$ ) |
| 1017 shr | $1020 \mathrm{~m}, \mathrm{pol}$ | $\mathrm{C}-\mathrm{C}$ stretch ( $A^{\prime}$ ) |
| $\begin{array}{r} 1005 \mathrm{~m} \\ 917 \mathrm{w} \end{array}$ | 1007 w | $\mathrm{C}-\mathrm{C}$ stretch ( $A^{\prime \prime}$ ) |
|  | 917 m, dp | $\mathrm{NC}_{2}$ stretch ( $E$ ) |
|  | 898 m , dp | $\mathrm{NC}_{2}$ stretch ( $E$ ) |
| 877 m | 875 w, dp | $\mathrm{NC}_{2}$ stretch (? $A$ ) |
| 808 shr | 809 w | $\mathrm{CH}_{2}$ o.o.p. rock ( $A^{\prime \prime}$ ) |
| 790 m777 shr $\mathrm{789} \mathrm{w} \quad \mathrm{CH}_{2}$ 1.p. rock ( $A^{\prime}$ ) |  |  |
| 723 w |  |  |
| 700 w | 704 vw |  |
| 592 w | $\begin{aligned} & 597 \mathrm{~s} . \mathrm{pol} \\ & 501 \mathrm{w} \end{aligned}$ | $\operatorname{AsN}_{3} \operatorname{stretch}(A+E)$ |
| 492 w |  |  |
| 477 w | 479 w | AsNC deformation ( $E$ ) |
|  | 433 w | ? AsNC deformation ( $E$ ) |
|  | 383 w | $\mathrm{NC}_{2}$ deformation ( $E$ ) |
|  | 374 w | $\mathrm{NC}_{2}$ deformation ( $E$ ) |
|  | 353 w | AsNC deformation ( $A$ ) |
|  | 333 w, dp | ${ }^{\text {? }} \mathrm{CH}_{3}$ torsion ( $A^{\prime \prime}$ ) |
|  | 297 s. pol | $\mathrm{CH}_{3}$ torsion ( $A^{\prime}$ ) |

N.B. i.p. o.o.p. are abbreviations for in-phase and out-of phase respectively.
$\mathrm{NMe}_{2}$ normal modes are set out in Table 5. If significant coupling occurs between the $\mathrm{NMe}_{2}$ units, then further bands would arise, and the problem would have to be treated under the overall symmetry of the molecule ( $C_{3}$ or $C_{3 v}$ ).

Finch et al. carried out an approximate normal coordinate analysis of an $\mathrm{XNMe}_{2}$ group with X taken to be of variable mass (1-100 a.m.u.) [12] This showed that for larger $X$ there is likely to be extensive mixing of modes, especially involving the methyl rocks and the $\mathrm{N}-\mathrm{X}$ stretches.

Table 5. Vibrational modes of an isolated - $\mathrm{NMe}_{2}$ ligand of $C_{s}$ symmetry

| $\mathrm{CH}_{3}$ stretches | $3 A^{\prime}+3 A^{\prime \prime}$ |
| :--- | :--- |
| $\mathrm{CH}_{3}$ deformations | $3 A^{\prime}+3 A^{\prime \prime}$ |
| $\mathrm{CH}_{3}$ rocks | $2 A^{\prime \prime}+2 A^{\prime \prime}$ |
| $\mathrm{CH}_{3}$ torsion | $A^{\prime}+A^{\prime \prime}$ |
| $\mathrm{NC}_{2}$ deformation | $A^{\prime}$ |
| $\left(-\mathrm{NC}_{2}\right.$ torsion | $\left.A^{\prime \prime}\right)$ |

In describing the assignment of internal $\mathrm{NMe}_{2}$ modes, figures quoted will refer to $\mathrm{P}\left(\mathrm{NMe}_{2}\right)_{3}$. The equivalent features in $\mathrm{As}\left(\mathrm{NMe}_{2}\right)_{3}$ almost always occur at very similar wavenumbers.

The highest-wavenumber fundamentals will, of course, be the $\mathrm{C}-\mathrm{H}$ stretches. Table 6 shows that for an isolated $-\mathrm{N}\left(\mathrm{CH}_{3}\right)_{2}$ group we expect 6-three of which should be polarized in the Raman spectrum. The anti-

Table 6. Vibrational modes for an $\mathrm{NEt}_{2}$ unit of $C_{s}$ symmetry

| $\mathrm{CH}_{3}$ stretches | $3 A^{\prime}+3 A^{\prime \prime}$ |
| :--- | :---: |
| $\mathrm{CH}_{2}$ stretches | $2 A^{\prime}+2 A^{\prime \prime}$ |
| $\mathrm{CH}_{3}$ deformations | $3 A^{\prime}+3 A^{\prime \prime}$ |
| $\mathrm{CH}_{2}$ scisors | $A^{\prime}+A^{\prime \prime}$ |
| $\mathrm{CH}_{2}$ wags | $A^{\prime}+A^{\prime \prime}$ |
| $\mathrm{CH}_{3}$ rocks | $2 A^{\prime}+2 A^{\prime \prime}$ |
| $\mathrm{CH}_{2}$ twists | $A^{\prime}+A^{\prime \prime}$ |
| $\mathrm{CH}_{2}$ rocks | $A^{\prime}+A^{\prime \prime}$ |
| $\mathrm{CH}_{3}$ torsions | $A^{\prime}+A^{\prime \prime}$ |
| $\mathrm{C}-\mathrm{C}$ stretches | $A^{\prime}+A^{\prime \prime}$ |
| $\mathrm{NC}_{2}$ stretches | $A^{\prime}+A^{\prime \prime}$ |
| $\mathrm{NC}_{2}$ deformation | $A^{\prime}$ |
| NCC deformations | $2 A^{\prime}+2 A^{\prime \prime}$ |

symmetric $\mathrm{CH}_{3}$ stretches generally occur at higher wavenumber than the symmetric, and can be assigned here to three medium-strong Raman bands, all depolarized at. 2974, 2931 and $2889 \mathrm{~cm}^{-1}$; all have i.r. counterparts. Of the three predicted symmetric modes. however. only two give observable features-strong and polarized Raman bands at 2843 and $2797 \mathrm{~cm}^{-1}$ (with corresponding i.r. features). These are in the region expected for this type of mode in $\mathrm{N}\left(\mathrm{CH}_{3}\right)_{2}$ compounds [14], but no band due to the final symmetric mode could be detected.
The methyl deformations should also give six bands,
three antisymmetric, predicted to fall within the range $1470-1410 \mathrm{~cm}^{-1}$, and three symmetric, of much more variable wavenumber (dependent upon the electronegativity of the attached atom) [15]. Depolarized Raman bands at 1481,1441 and $1417 \mathrm{~cm}^{-1}$ are definitely due to the three antisymmetric deformations, but no polarized Raman bands occur anywhere near these wavenumbers. Strong i.r. absorptions are, however, noted at 1463,1458 and $1318 \mathrm{~cm}^{-1}$, and these are assigned as the deformations of $A^{\prime}$ symmetry. This complete absence of symmetric modes from the Raman spectrum is most unusual.

The modes described so far are probably quite free from "mixing," but this is no longer true when discussing the $\mathrm{CH}_{3}$ rocking and $\mathrm{NC}_{2}$ stretching modes. Descriptions of modes appearing in the $900-1300 \mathrm{~cm}^{-1}$ region are likely to be approximate, as shown by the calculations of Finch et al. [12], and assignments will be tentative. The following assignments are all quite reasonable, however. Four methyl rocking modes are expected, and can be assigned to four strong i.r. absorptions, at $1277,1190,1150$ and $1086 \mathrm{~cm}^{-1}$. A pglarized Raman band at $1147 \mathrm{~cm}^{-1}$ shows that this corresponds to an $A^{\prime}$ mode, while depolarized bands at 1280 and $1195 \mathrm{~cm}^{-1}$ indicate that these are of $A^{\prime \prime}$ symmetry. The i.r. band at $1086 \mathrm{~cm}^{-1}$ has no Raman counterpart, but is necessarily the remaining $A^{\prime}$ mode.

For the modes which involve only the $\mathrm{CH}_{3}$ groups, no evidence has been found for coupling between the $\mathrm{NMe}_{2}$ groups. For the $\mathrm{NC}_{2}$ stretches, however, it is possible that such coupling may be significant. As Table 5 shows, an isolated $\mathrm{NMe}_{2}$ group would give only two such modes ( $A^{\prime}+A^{\prime \prime}$ ), but Table 7 (which will be discussed in detail for the skeletal modes of X $\left.\left(\mathrm{NMe}_{2}\right)_{3}\right)$ reveals that for the "whole molecule" model there would be 4 distinct $\mathrm{NC}_{2}$ stretches, all i.r. and Raman active for $C_{3}$ symmetry $(2 A+2 E)$, but only three active modes for $C_{3 v}$ symmetry $\left(A_{1}+2 E ; A_{2}\right.$ totally inactive.)

Table 7. Skeletal vibrations of $\mathrm{X}\left(\mathrm{NC}_{2}\right)_{3}$ units

| Mode type | Effective symmetry |  |
| :--- | :---: | :---: |
| $C_{3}$ | $C_{3 v}$ |  |
| $\mathrm{X}-\mathrm{N}$ stretch | $A+E$ | $A_{1}+E$ |
| $\mathrm{NC} C_{2}$ stretch | $2 A+2 E$ | $A_{1}+A_{2}+2 E$ |
| $\mathrm{X}-\mathrm{N}-\mathrm{C}$ deformation | $2 A+2 E$ | $A_{1}+A_{2}+2 E$ |
| $\mathrm{NC}_{2}$ deformation | $A+E$ | $A_{1}+E$ |
| $\mathrm{NC}_{2}$ torsion | $A+E$ | $A_{2}+E$ |

There is some uncertainty as to the position expected of the antisymmetric $\mathrm{NC}_{2}$ stretch. Thus, DURIG and CASPER [4] favour a value of $c a .1250 \mathrm{~cm}^{-1}$, while other workers [3.6] favour a lower wavenumber. Since in $\mathrm{HNMe}_{2}$ the $\mathrm{NC}_{2}$ stretches are at $930 \mathrm{~cm}^{-1}\left(A^{\prime}\right)$ and $1024 \mathrm{~cm}^{-1}\left(A^{\prime \prime}\right)$, we follow the latter alternative. In the Raman spectrum of $\mathrm{P}\left(\mathrm{NMe}_{2}\right)_{3}$ we observe bands at $963 \mathrm{~cm}^{-1}$ (pol) and $981 \mathrm{~cm}^{-1}$ (depol.), which are clearly
due to symmetric and antisymmetric modes respectively. There are very strong i.r. absorptions corresponding to both of these and, in addition, absorptions at 1069 and $947 \mathrm{~cm}^{-1}$ which are of medium intensity in the liquid-phase, and strong in solid-phase spectra. These can also be assigned as $\mathrm{NC}_{2}$ stretching fundamentals and we can therefore conclude (i) that there is sufficient vibrational coupling between the $\mathrm{NMe}_{2}$ units to break down the "local symmetry" approximation for $\mathrm{NC}_{2}$ stretches, and (ii) that the effective molecular symmetry must be $C_{3}$ rather than $C_{30}$, since in the latter case only three stretches would be seen.

The $\mathrm{X}-\mathrm{N}-\mathrm{C}$ and $\mathrm{NC}_{2}$ deformations are most conveniently discussed as molecular skeletal modes, leaving only the methyl torsions to be assigned in this section. In $\mathrm{HNMe}_{2}$ these are at $290 \mathrm{~cm}^{-1}\left(A^{\prime}\right)$ and $250 \mathrm{~cm}^{-1}\left(A^{\prime \prime}\right)$ [11] in $\mathrm{P}\left(\mathrm{NMe}_{2}\right)_{3}$ a polarized Raman band is present at $295 \mathrm{~cm}^{-1}$, while in $\operatorname{As}\left(\mathrm{NMe}_{2}\right)_{3}$ a very similar feature is seen at $273 \mathrm{~cm}^{-1}$; these are assigned to the symmetric torsion, and no evidence was found for the antisymmetric mode.

## (b) Internal $\mathrm{NEt}_{2}$ modes

In this discussion data from $\mathrm{As}\left(\mathrm{NEt}_{2}\right)_{3}$ will generally be used for illustration, as the spectrum of $\mathrm{P}\left(\mathrm{NEt}_{2}\right)_{3}$ was less well resolved. In addition, the greater complexity of the data will lead to uncertainties in the assignments. and only very brief discussions will be given. The numbers and symmetry types of vibrations for an $\mathrm{NEt}_{2}$ group of $C_{s}$ symmetry are summarised in Table 6. Ten CH stretches should be present, and there are 10 observed wavenumbers, but assignment to specific modes is not easy. The assignments in Table 4 are, however, consistent with accepted characteristic wavenumbers in this region.

Methylene scissors deformation modes, when the $\mathrm{CH}_{2}$ is adjacent to an amine residue, are generally within the range $1475-1445 \mathrm{~cm}^{-1}$ [16]. We have two candidates, at 1485 and $1462 \mathrm{~cm}^{-1}$; the polarizations of neither could be detected in the Raman spectrum, so they are assigned arbitrarily as $A^{\prime \prime}, A^{\prime}$ respectively. Six bands can be assigned as $\mathrm{CH}_{3}$ deformations, 1452$1323 \mathrm{~cm}^{-1}$ (Table 4).

All of the remaining $\mathrm{CH}_{2}$ deformations (twist, wag, rock) and $\mathrm{CH}_{3}$ rocking modes can be assigned similarly to features in the normaily expected regions, and are summarised in Table 4.
$\mathrm{C}-\mathrm{C}$ stretching modes in a number of diethylamido derivatives of boron are assigned ${ }^{17}$ as $1008 \mathrm{~cm}^{-1}$ (symmetric) and $1080 \mathrm{~cm}^{-1}$ (antisymmetric). In As$\left(\mathrm{NEt}_{2}\right)_{3}$, a medium-intensity, polarized Raman band is seen at $1020 \mathrm{~cm}^{-1}$, and is assigned as the $A^{\prime}$ mode (i.r. at $1017 \mathrm{~cm}^{-1}$ ). The only band near this which could be due to the $A^{\prime \prime}$ mode is at $1007 \mathrm{~cm}^{-1}$ (of undetermined polarization).

As for the $\mathrm{X}\left(\mathrm{NMe}_{2}\right)_{3}$ compounds, more $\mathrm{NC}_{2}$ stretches are seen than can be accounted for by a single, uncoupled $\mathrm{NR}_{2}$ group. In $\mathrm{P}\left(\mathrm{NEt}_{2}\right)_{3}$, four bands are seen. as expected for a $C_{3}$ model, at $973.887 \mathrm{~cm}^{-1}(A$ symmetry). $942.917 \mathrm{~cm}^{-1}(E$ symmetry $) . \operatorname{As}\left(\mathrm{NEt}_{2}\right)_{3}$
gives only three such bands, $876(A) .917,898 \mathrm{~cm}^{-1}$ $(E)$-but the same symmetry is likely to apply here also. These wavenumbers are lower than for the $\mathrm{NMe}_{2}$ compounds. This may be due to mass effects, or to coupling with $\nu \mathrm{C}-\mathrm{C}$ or $\mathrm{CH}_{3}$ rocks and $\mathrm{CH}_{2}$ twists.

As for the $\mathrm{NMe}_{2}$ groups, the $\mathrm{X}-\mathrm{N}-\mathrm{C}$ and $\mathrm{NC}_{2}$ deformations are discussed as skeletal modes. This leaves only the $\mathrm{CH}_{3}$ torsion-for which a Raman band (polarized) at $297 \mathrm{~cm}^{-1}$ is assigned as the $A^{\prime}$ mode, with a depolarized feature at $333 \mathrm{~cm}^{-1}$ possibly being the $A^{\prime \prime}$ mode (As compound); the former feature is at $311 \mathrm{~cm}^{-1}$ in $\mathrm{P}\left(\mathrm{NEt}_{2}\right)_{3}$.
(c) Skeletal modes of $\mathrm{P}\left(\mathrm{NR}_{2}\right)_{3}$

Detailed discussion will be given for $R=M e$, as assignments for the ethyl analogue are very similar. Predicted modes are summarised in Table 7. The $\mathrm{NC}_{2}$ stretches have already been described, and the nexthighest fundamentals will be the $\mathrm{PN}_{3}$ antisymmetric stretches. Goubeau et al. showed [6] that for OP$\left(\mathrm{NMe}_{2}\right)_{3}$ and $\mathrm{SP}\left(\mathrm{NMe}_{2}\right)_{3}$ these were at $752.742 \mathrm{~cm}^{-1}$ respectively, with $v_{s} \mathrm{PN}_{3}$ at $636,722 \mathrm{~cm}^{-1}$ respectively. Exocyclic PN stretches in $\mathrm{P}_{3} \mathrm{~N}_{3} \mathrm{~F}_{6-n}\left(\mathrm{NMe}_{2}\right)_{n}$ gave bands at $680 \mathrm{~cm}^{-1}$ (symmetric) and $748 \mathrm{~cm}^{-1}$ (antisymmetric) [18]. $v_{s} \mathrm{PN}_{3}$ in $\mathrm{P}\left(\mathrm{NMe}_{2}\right)_{3}$ is very easily assigned to a very strong, polarized Raman band at $671 \mathrm{~cm}^{-1}$, while $v_{a s}$ is apparently only slightly higher, at $700 \mathrm{~cm}^{-1}$.

Both $\mathrm{PN}_{3}$ modes are shifted to lower wavenumber in $\mathrm{P}\left(\mathrm{NEt}_{2}\right)_{3}$, probably due to simple mass effects ( $v_{s}$ $655 \mathrm{~cm}^{-1}, v_{a s} 667 \mathrm{~cm}^{-1}$ ).

PNC and $\mathrm{NC}_{2}$ deformation modes will certainly be mixed extensively, and so assignments will be approximate. In $\mathrm{HNMe}_{2}, \delta \mathrm{NC}_{2}$ is at $397 \mathrm{~cm}^{-1}$ [12], and at $393 \mathrm{~cm}^{-1}$ in $\mathrm{MeNPF}_{2}$ [3]. Consequently, a tentative assignment of two $\mathrm{NC}_{2}$ deformations in $\mathrm{P}\left(\mathrm{NMe}_{2}\right)_{3}$ is made to bands at 390 (i.r.), $392 \mathrm{~cm}^{-1}$ (Raman, unknown polarization) and $405 \mathrm{~cm}^{-1}$ (i.r. only) $\mathrm{cm}^{-1}$. It is impossible to differentiate between that of $A$ and that of $E$ symmetry. This leaves candidates for $\delta$ PNC as follows: (symmetric) $337 \mathrm{~cm}^{-1}$ (strong, polarized Raman band); (antisymmetric) $494 \mathrm{~cm}^{-1}$ (i.r.)/497 $\mathrm{cm}^{-1}$ (Raman). and $420 \mathrm{~cm}^{-1}$ (i.r.)/416 $\mathrm{cm}^{-1}$ (Raman). No evidence was found for the second symmetric deformation.

Two depolarized Raman bands are seen, at 197 and $127 \mathrm{~cm}^{-1}$, due to the $\mathrm{NC}_{2}$ torsion and $\delta_{a s} \mathrm{PN}_{3}$ respectively. The $\delta_{s} \mathrm{PN}_{3}$ would be expected below $100 \mathrm{~cm}^{-1}$, and was not detected, while the symmetric torsion mode is likely to only weakly allowed, since it is derived from an $A_{2}$ mode (forbidden) under $C_{3 v}$ symmetry.
(d) Skeletal modes of $\mathrm{As}\left(\mathrm{NR}_{2}\right)_{3}$

As in (c) above. the case with $\mathrm{R}=\mathrm{Me}$ will be considered in greater detail. The $\mathrm{NC}_{2}$ stretches, as for $\mathrm{P}\left(\mathrm{NMe}_{2}\right)_{3}$ are all above $900 \mathrm{~cm}^{-1}$, the $A$ modes giving a broad. polarized feature with maxima at 951 and $937 \mathrm{~cm}^{-1}$, the $E$ modes being at 1063 and (probably) $1024 \mathrm{~cm}^{-1}$.

The next features to be considered at those due to $v_{s}$
and $v_{a s}$ AsN $_{3}$. Kober has reported [19] that $v \mathrm{AsN}$ in several $\mathrm{Me}_{2} \mathrm{AsNR} 2_{2}$ compounds is always close to $580 \mathrm{~cm}^{-1}$, while DURIG showed that in $\mathrm{Me}_{2} \mathrm{NAsCl}_{2}$ it is at $585 \mathrm{~cm}^{-1}$ (trans) or $569 \mathrm{~cm}^{-1}$ (gauche) [4]. Thus, in $\mathrm{As}\left(\mathrm{NMe}_{2}\right)_{3}$ it is clear that $v_{5} \mathrm{AsN}_{3}$ corresponds to the strong, polarized Raman band at $574 \mathrm{~cm}^{-1}$ (with an i.r. counterpart). In $P\left(\mathrm{NR}_{2}\right)_{3}$, as seen above $v_{s}$ and $v_{a s} \mathrm{PN}_{3}$ are separated by $c a .20 \mathrm{~cm}^{-1}$. It is a general rule that on increasing the mass of the central atom $X$, the separation between $v_{s}$ and $v_{a s}$ for $\mathrm{XY}_{n}$ decreases. In fact no other band is seen between 580 and $600 \mathrm{~cm}^{-1}$ for $\mathrm{As}\left(\mathrm{NMe}_{2}\right)_{3}$ or $\mathrm{As}\left(\mathrm{NEt}_{2}\right)_{3}$, and thus $v_{s}$ and $v_{a s}$ appear to be accidentally degenerate. Support for this view comes from the observation of only one $\nu \mathrm{SN}$ band in $\mathrm{S}\left(\mathrm{NMe}_{2}\right)_{2}$ [20], and only one $v \mathrm{SbN}$ in $\mathrm{Sb}\left(\mathrm{NMe}_{2}\right)_{3}$ (the only assignment given for this species) [21].

As in the phosphorus analogues, some mixing is expected between $\delta \mathrm{AsNC}$ and $\delta \mathrm{NC}_{2}$ modes, so descriptions will be approximate. An i.r. band at $480 \mathrm{~cm}^{-1}$ in $\mathrm{As}\left(\mathrm{NMe}_{2}\right)_{3}$ (no Raman counterpart) is assigned as $\delta_{a s} A s N C$, while the equivalent symmetric mode is at $308 \mathrm{~cm}^{-1}$ (strong, polarized Raman band), and a second antisymmetric mode at $339 \mathrm{~cm}^{-1}$ (depolarized, Raman only). The symmetric $\mathrm{NC}_{2}$ deformation is at $397 \mathrm{~cm}^{-1}$ (strong, polarized Raman band), with an antisymmetric deformation at $378 \mathrm{~cm}^{-1}$.

An $\mathrm{NC}_{2}$ torsional mode gives a depolarized Raman band at $153 \mathrm{~cm}^{-1}$, with $\delta_{a 5} A s N_{3}$ at $109 \mathrm{~cm}^{-1}$. The symmetric $\mathrm{AsN}_{3}$ deformation ( $<100 \mathrm{~cm}^{-1}$ ) was not detected.

## CONCLUSION

We have been able to assign satisfactorily the internal modes of the $\mathrm{NR}_{2}$ units in $\mathrm{X}\left(\mathrm{NR}_{2}\right)_{3}$, where $\mathrm{X}=\mathrm{P}$ or As; $\mathbf{R}=\mathrm{Me}$ or Et. Except for $\mathrm{NC}_{2}$ modes, a model involving non-interacting $\mathrm{NR}_{2}$ units, of $C_{s}$ symmetry, was sufficient. For skeletal modes (including $\mathrm{NC}_{2}$ ) the possible symmetries were $C_{3}$ and $C_{3 v}$. For $v \mathrm{NC}_{2}$ there were definitely too many fundamental bands for the latter, and so $C_{3}$ appears to be the effective molecular symmetry.

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