

Synthesis and Structural Characterization of the layered Selenogallate RbGaSe_2

Daniel Friedrich,^[a] Marc Schlosser,^[a] and Arno Pfitzner*^[a]

Dedicated to Professor Wolfgang Schnick on the Occasion of his 60th Birthday

Abstract. The chalcogenogallate RbGaSe_2 was synthesized by thermal decomposition of rubidium azide in the presence of gallium selenide and selenium. RbGaSe_2 crystallizes in the monoclinic space group $C2/c$ (no. 15) with the lattice parameters $a = 10.954(1) \text{ \AA}$, $b = 10.949(1) \text{ \AA}$, $c = 16.064(1) \text{ \AA}$, $\beta = 99.841(4)^\circ$, $V = 1898.2(2) \text{ \AA}^3$, and $Z = 16$ (single-crystal data, 20°C) in the TlGaSe_2 structure type. Its crystal structure features anionic layers ${}_2[\text{Ga}_4\text{Se}_8^{4-}]$ with a van der

Waals distance of $3.30(1) \text{ \AA}$. Thermal analysis revealed a melting point of about 930°C . Using UV/Vis diffuse reflectance spectroscopy, a wide bandgap of 3.16 eV was determined for the colorless semiconductor. The bonding situation in the compound was further compared with the isotopic cesium phases based on the results from Raman spectroscopy and DFT calculations. The slight shifts of the Raman bands suggest a slightly higher covalency of the rubidium compound.

Introduction

Due to their semiconducting properties, group 13 chalcogeno-metallates containing alkali metal cations $M_xT_yQ_z$ (M = alkali metal, T = triel, Q = chalcogen) are interesting materials for technical applications.^[1] These chalcogenotrirelates are known to crystallize in a large variety of diverse crystal structure types.^[2] The smallest structural building unit of these structures are TQ_4 tetrahedra. Condensation of these building blocks yields more complex polymeric one-, two- or three-dimensional anions, which are embedded in a cationic surrounding of the alkali metal cations. The MTQ_2 phases are known for almost all possible element combinations among these chalcogenotrirelates. The dominating structure type among these solids is the TlGaSe_2 structure type,^[3] featuring anionic layers ${}_2[T_4Q_8^{4-}]$, which seems to have a very high tolerance for different combinations of the elements involved.^[4] Another possible anionic substructure of these MTQ_2 phases are anionic chains ${}_1[Q_2^-]$ as found in the Tl_2Se_2 ^[5] and KFeSe_2 structure types.^[6]

In this paper, we report on the synthesis and structural characterization of RbGaSe_2 , one of the very few missing phases MTQ_2 in the ternary systems of alkali metal, triel, and chalcogen. The title compound was further characterized by Raman- and UV/Vis spectroscopy. We examined the solid at high temperatures regarding a possible phase-transition similar to the isotopic cesium compounds $\text{CsGaQ}_2\text{-}mC64$ ($Q = \text{S, Se}$).^[7] This

is especially interesting in order to further investigate the chemical similarities of Rb with Cs and K.

Results and Discussion

Crystal Structure Determination of RbGaSe_2

RbGaSe_2 forms air- and moisture sensitive, colorless, plate-shaped crystals. In order to minimize the significant stacking disorder present in this compound and isotopic phases,^[8] a small rod-like fragment was broken off from a suitable plate-shaped crystal. The crystal structure was solved from single-crystal X-ray diffraction data collected at 20°C . Due to the high fluorescence of the rubidium and selenium atoms when using $\text{Mo-}K_\alpha$ radiation, $\text{Cu-}K_\alpha$ radiation was used for the measurement. The compound crystallizes in the monoclinic space group $C2/c$ (no. 15) with the lattice parameters $a = 10.954(1) \text{ \AA}$, $b = 10.949(1) \text{ \AA}$, $c = 16.064(1) \text{ \AA}$, $\beta = 99.841(4)^\circ$, $V = 1898.2(2) \text{ \AA}^3$, and $Z = 16$ (single-crystal data, 20°C). The crystal structure was solved by charge flipping methods using SUPERFLIP,^[9] implemented in Jana 2006,^[10] and refined to $R_1 = 0.0404$ and $wR_2 = 0.0863$ (all data) using Jana2006.^[10] Similar to the isotopic compound $\text{CsGaSe}_2\text{-}mC64$, RbGaSe_2 single crystals are notoriously twinned. Using TwinRotMat of the PLATON software package,^[11] a twin matrix identical to $\text{CsGaSe}_2\text{-}mC64$ was found.^[7a] Details on the structure solution and refinement are listed in Table 1. The atomic coordinates, displacement parameters as well as selected interatomic distances and angles are provided in the supplementary material (Tables S1–S3, Supporting Information).

RbGaSe_2 crystallizes isotypic to $\text{CsGaSe}_2\text{-}mC64$ and many related compounds MTQ_2 ($M = \text{Na, K, Rb, Cs, Tl}$; $T = \text{Al, Ga, In}$; $Q = \text{S, Se, Te}$) in the TlGaSe_2 structure type.^[3] The anisotropic crystal structure of this solid features anionic layers ${}_2[\text{Ga}_4\text{Se}_8^{4-}]$, composed of corner sharing $\text{Ga}_4\text{Se}_{10}$ super-tetra-

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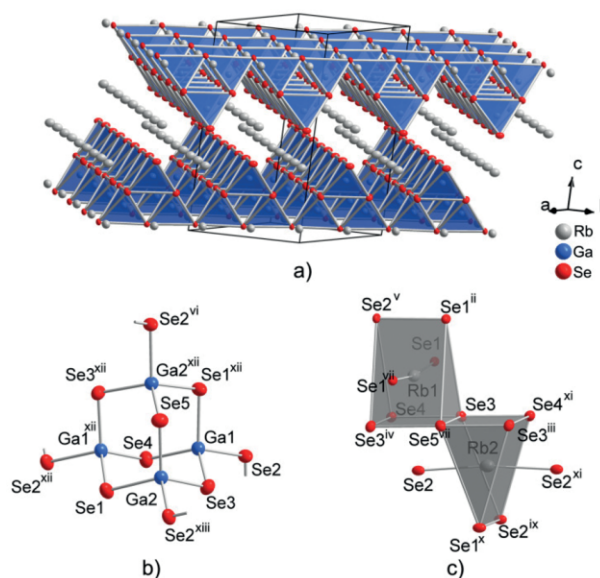
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Table 1. Crystallographic data for RbGaSe₂.

	RbGaSe ₂
Formula weight $M / \text{g} \cdot \text{mol}^{-1}$	313.11
Color and shape	Colorless rod-like fragment
Crystal size /mm ³	$0.24 \times 0.04 \times 0.02$
Space group	$C2/c$ (no. 15)
Crystal system	monoclinic
Lattice parameters from single crystal	
$a / \text{\AA}$	10.954(1)
$b / \text{\AA}$	10.949(1)
$c / \text{\AA}$	16.064(1)
$\beta / ^\circ$	99.841(4)
Unit cell volume $V / \text{\AA}^3$	1898.2(2)
Number of formula units Z	16
Calculated density $\rho / \text{g} \cdot \text{cm}^{-3}$	4.382
Diffractometer	Rigaku Supernova
Temperature $T / ^\circ\text{C}$	20(1)
Wavelength $\lambda / \text{\AA}$	1.54184
Absorption coeff. $\mu(\text{Cu-K}\alpha) / \text{mm}^{-1}$	36.236
θ range / $^\circ$	5.59–76.52
Index range	$-13 \leq h \leq 12$ $-13 \leq k \leq 13$ $-20 \leq l \leq 20$
No. of reflections collected	11049
Independent reflections	2045
R_{int}, R_σ	0.0119, 0.0375
Completeness to $\theta = 76.52^\circ$	99.9 %
Absorption correction	Numerical, CrysAlisPro [12]
Structure solution	Charge flipping, Superflip [9]
Structure refinement	Jana2006 [10]
Twin law	$-1 \ 0 \ 0, 0 \ -1 \ 0, \frac{1}{2} \ 0 \ 1$
Twin fraction	26.1(4) %
Data / restraints / parameters	2045 / 0 / 76
GooF, $F(000)$	1.70, 2176
$R_1, wR_2 [I > 3\sigma(I)]$	0.0298, 0.0737
$R_1, wR_2 [\text{all data}]$	0.0404, 0.0863
Extinction coefficient G_{iso}	0.144(1)
Largest diff. peak and hole / $\text{e} \cdot \text{\AA}^{-3}$	1.26, −1.32

hedra (Figure 1). These layers are stacked along [001] and are separated by Rb atoms. The van der Waals distance between two adjacent layers is 3.30(1) Å. Both independent Ga³⁺ sites are tetrahedrally coordinated by Se^{2−} anions. The interatomic distances in these slightly distorted tetrahedra range from $d(\text{Ga}–\text{Se}) = 2.388(1) \text{ \AA}$ to $d(\text{Ga}–\text{Se}) = 2.428(1) \text{ \AA}$. The angles Se–Ga–Se range from 106.04(3)° to 113.45(3)°. These values differ only slightly from the sum of the ionic radii,^[13] $d(\text{Ga}^{3+}–\text{Se}^{2-}) \approx 2.5 \text{ \AA}$, and are in good agreement with comparable compounds like Cs₂Ga₂Se₅,^[14] and both polymorphs of CsGaSe₂^[7a] ($Q = \text{S}, \text{Se}$). The two independent Rb⁺ sites are 6+2 fold coordinated by Se^{2−}, resulting in bicapped trigonal prisms (Figure 1). The interatomic distances in these prisms range from $d(\text{Rb}–\text{Se}) = 3.428(1) \text{ \AA}$ to $d(\text{Rb}–\text{Se}) = 3.555(1) \text{ \AA}$ while the distances to the capping atoms ranging from $d(\text{Rb}–\text{Se}) = 3.896(1) \text{ \AA}$ to $d(\text{Rb}–\text{Se}) = 3.911(1) \text{ \AA}$ are significantly longer. The RbSe₈ prisms are connected by common edges, faces, and the capping atoms (Se1, Se2). The eightfold interconnected rubidium atoms form an irregular topological network. The distances between the rubidium atoms in one layer are $d(\text{Rb}–\text{Rb}) < 5 \text{ \AA}$, whereas the distances between adjacent layers are in the range of $5.4 \text{ \AA} < d(\text{Rb}–\text{Rb}) < 5.5 \text{ \AA}$.

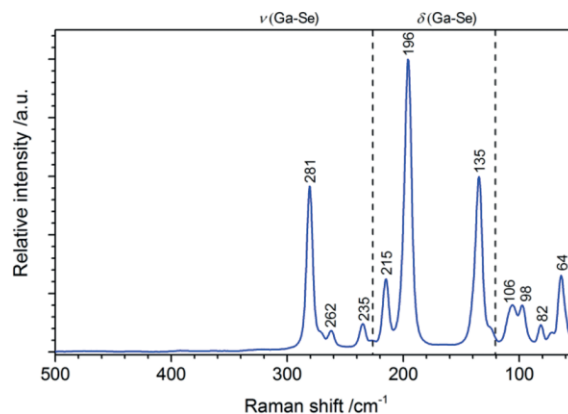
**Figure 1.** The crystal structure of RbGaSe₂, showing: (a) stacking of the anionic selenogallate layers; (b) a supertetrahedral Ga₄Se₁₀ unit; (c) coordination polyhedra of the Rb atoms.

Thermal Analysis

In order to analyze the behavior of RbGaSe₂ at high temperature, the solid was investigated by high-temperature X-ray powder diffraction in the temperature range from 20–900 °C. The measurement revealed no phase-transitions for RbGaSe₂ contrary to the isotopic phases CsGaQ₂-mC64 ($Q = \text{S}, \text{Se}$)^[7] as shown in Figure S1 (Supporting Information). As no melting point of the sample could be observed in the temperature region investigated by diffraction experiments, we performed a differential thermal analysis (DTA) up to 1000 °C. This measurement revealed a melting point of about 930 °C (Figure S2, Supporting Information).

Raman Spectroscopy

The Raman spectrum of RbGaSe₂ is shown in Figure 2. The bands at 281, 262, and 235 cm^{−1} result from Ga–Se stretching modes, whereas the vibrations at 215, 196, and 135 cm^{−1} can be attributed to Ga–Se deformation modes.^[7a,14–15] All vi-

**Figure 2.** Raman spectrum of RbGaSe₂.

brations at lower frequencies are tentatively assigned to lattice vibrations. The Raman spectra of RbGaSe_2 and $\text{CsGaSe}_2\text{-mC64}$ ^[7a] are almost identical. The vibration frequencies in the rubidium compound, however, show a small blue-shift of about $2\text{--}10\text{ cm}^{-1}$ which indicates stronger Ga–Se bonding interactions in RbGaSe_2 compared to $\text{CsGaSe}_2\text{-mC64}$.

Optical Properties

The optical bandgap of RbGaSe_2 was determined by using UV/Vis diffuse reflectance spectroscopy (Figure 3). The absorption spectrum was calculated using a modified Kubelka-Munk function.^[16] Extrapolation of the linear part onto the baseline revealed a wide bandgap of 3.16 eV which is in good agreement with the colorless powdered substance.

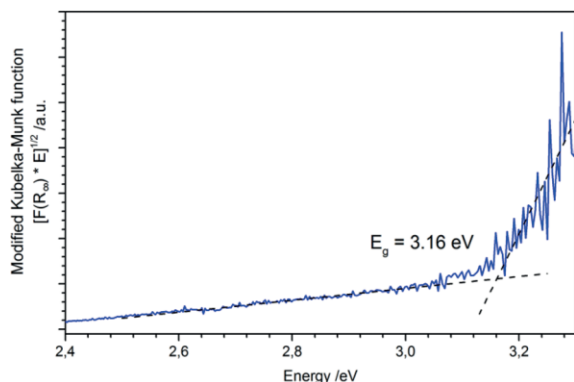


Figure 3. Diffuse reflectance spectrum of RbGaSe_2 . The optical bandgap was determined by extrapolation of the linear part of the modified Kubelka-Munk function onto the baseline, which is indicated by the dashed lines.

We further attempted to calculate the electronic band structure of RbGaSe_2 from relativistic DFT calculations. Similar to $\text{CsGaSe}_2\text{-mC64}$,^[7a] we observed significant discrepancies between the experimental and calculated bandgaps using the GGA^[17] and LSDA^[18] functionals. The most plausible result is depicted in Figure 4. Interestingly, the bandgap values are identical for $\text{CsGaSe}_2\text{-mC64}$ ($E_g^{\text{exp}} = 3.16\text{ eV}$, $E_g^{\text{calc}} \approx 2.0\text{ eV}$) and RbGaSe_2 ($E_g^{\text{exp}} = 3.16\text{ eV}$, $E_g^{\text{calc}} \approx 2.0\text{ eV}$). Furthermore, the electronic structures of both compounds are basically iden-

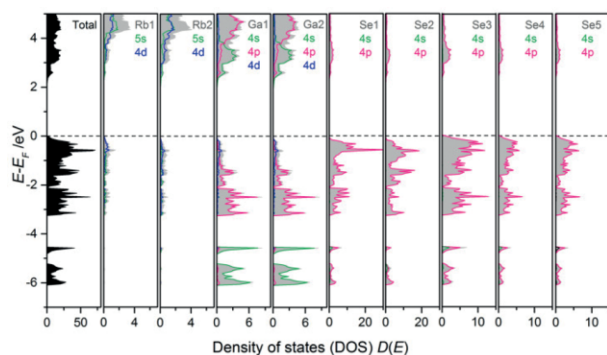


Figure 4. Calculated total (TDOS) and partial density of states (PDOS) for RbGaSe_2 . The highest occupied bands are mainly due to Se-4p states whereas the lowest energies of the conduction band show contributions mostly from Ga-4s states.

tical except for the alkali metal states, which of course do not contribute to the bandgap energy. The highest occupied bands are mainly due to Se-4p states, whereas the lowest energies of the conduction band show contributions mostly from Ga-4s states. This observation further confirms our previous assumption that the electronic properties of these compounds strongly depend on the covalent interactions between gallium and chalcogen atoms, whereas the alkali metals do not significantly contribute to the bandgap.^[7,19–20] The fact that the DFT calculations for the selenides show these discrepancies, which were not observed in case of the sulfides,^[7b] indicates that conventional calculations significantly underestimate the Ga–Se interactions in the selenogallate layers.

Conclusions

RbGaSe_2 , a new phase in the ternary systems of alkali metal, gallium, and selenium was synthesized and structurally characterized. Contrary to the isotypic cesium phases, RbGaSe_2 does not undergo a polymorphic phase-transition at elevated temperatures below its melting point. DFT calculations combined with Raman spectroscopy revealed that even though the electronic structure of RbGaSe_2 and $\text{CsGaSe}_2\text{-mC64}$ are almost identical, slightly stronger covalent Ga–Se interactions can be observed in the rubidium compound.

Experimental Section

Synthesis of the Starting Materials: GaSe was prepared by chemical vapor transport reaction from gallium (Chempur 99.99 %) and selenium (Chempur 99.999 %) using iodine (Sigma-Aldrich 99.8 %) as transporting agent.^[21] Pure RbN_3 was obtained by passing hydrazoic acid (prepared by acidifying an aqueous solution of NaN_3 (Sigma-Aldrich 99.0 %) into an aqueous solution of Rb_2CO_3 (Sigma-Aldrich 99 %)).^[22]

Caution: Condensed HN_3 is highly explosive. Tools made from transition metals must be avoided.

Synthesis of RbGaSe_2 : The title compound was synthesized by controlled thermal decomposition (0.2 g batch size, heating rate $0.5\text{ K}\cdot\text{min}^{-1}$) of RbN_3 combined with a stoichiometric mixture of GaSe and selenium in a quartz glass ampoule under dynamic vacuum conditions. Upon heating, the azide decomposes to the pure alkali metal which immediately reacts with the other starting materials, resulting in an inhomogeneous raw product. A flame sealed ampoule containing this raw product was annealed at $850\text{ }^\circ\text{C}$ to yield pure RbGaSe_2 . The advantage of this azide route is the avoidance of handling larger amounts of pure and highly reactive alkali metals, which are difficult to dose in proper portions. However, since the azide decomposes to the elements, the reaction can be regarded as equivalent to a batch containing only the elements.

Single-Crystal X-ray Diffraction: Data collections on suitable single-crystals were performed with a Rigaku Supernova at $20\text{ }^\circ\text{C}$ using monochromatized $\text{Cu-K}\alpha$ radiation ($\lambda = 1.54184\text{ \AA}$). The diffraction data were corrected for Lorentz- and polarization effects and corrected by a numerical absorption correction using the CrysAlisPro software package.^[12] The resulting data set had a completeness of 99.9 % within $75^\circ\theta$. The crystal structure was solved by charge flipping methods using SUPERFLIP^[9] (implemented in Jana2006^[10]) and refined on F^2

with Jana2006^[10] using full-matrix least-squares methods. The experimental parameters and details on the structure solution and refinement are summarized in Table 1.

Further details of the crystal structures investigations may be obtained from the Fachinformationszentrum Karlsruhe, 76344 Eggenstein-Leopoldshafen, Germany (Fax: +49-7247-808-666; E-Mail: crysdata@fiz-karlsruhe.de, <http://www.fiz-karlsruhe.de/request> for deposited data.html) on quoting the depository numbers CSD-433600.

X-ray Powder Diffraction: The X-ray powder diffraction experiments were performed with a STOE STADI P diffractometer equipped with a Dectris Mythen 1 K detector and a high-temperature capillary furnace. Monochromatized Mo- $K_{\alpha 1}$ radiation ($\lambda = 0.70930 \text{ \AA}$) was used in all experiments. The air sensitive samples were flame sealed in a 0.3 mm quartz glass capillary and placed in another 0.5 mm quartz glass capillary because of the furnace architecture. The WinX^{POW} software package from STOE & Cie was used for data collection and processing.^[23] The furnace temperature was controlled by a Eurotherm 24.16 controller ($\Delta T = \pm 1 \text{ }^{\circ}\text{C}$). A comparison of the diffraction patterns of phase-pure RbGaSe₂, measured with Cu- $K_{\alpha 1}$ radiation ($\lambda = 1.540598 \text{ \AA}$), with a calculated pattern based on the single crystal data is shown in Figure S3 (Supporting Information).

Differential Thermal Analysis: Differential thermal analysis of the compound was performed with a SETARAM TG-DTA 92 16.18 in an evacuated quartz glass ampoule using Al₂O₃ as reference material. The measurements were performed in the range of 25–1000 °C with a heating- and cooling rate of 10 K·min⁻¹ (2 cycles).

Vibrational Spectroscopy: The Raman spectrum was recorded with a DXRTM SmartRaman Spectrometer from ThermoScientific (excitation wavelength $\lambda = 532 \text{ nm}$) in the range of 50–1000 cm⁻¹ with a resolution of 0.5 cm⁻¹.

UV/Vis Spectroscopy: Diffuse reflectance measurements were performed with a Bruins Omega 20 UV/Vis spectrometer using BaSO₄ as a reference (100% reflectance). The absorption spectrum was calculated using a modified Kubelka-Munk function.^[16]

DFT Calculations: The first-principles calculations were carried out within the framework of density functional theory (DFT) with exchange-correlation functionals in the generalized gradient approximation (GGA) according to Perdew-Burke-Ernzerhof (PBE),^[17] and the local-spin-density approximation (LSDA).^[18] The full potential local orbital code FPLO14^[24] with a k -grid mesh of $12 \times 12 \times 12$ was applied for total energy and band structure calculations based on the experimentally obtained structures. All of these calculations converged with a maximum divergence of $<10^{-7}$ Hartree.

Supporting Information (see footnote on the first page of this article): Tables S1–S3 display atomic coordinates, anisotropic displacement parameters, and selected interatomic distances of RbGaSe₂. Figures S1–S3 show the evolution of the X-ray powder diffractogram in the range 20–900 °C, a DTA curve from 25 to 1000 °C, and an experimental X-ray powder diffractogram in comparison to a calculated one.

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Keywords: Crystal structure; Density functional calculations; Layered material; Selenogallate; Semiconductors

References

- [1] a) J.-J. Zondy, F. Bielsa, A. Douillet, L. Hilico, O. Acef, V. Petrov, A. Yelissev, L. Isaenko, P. Krinitsin, *Opt. Lett.* **2007**, *32*, 1722–1724; b) V. Petrov, A. Yelissev, L. Isaenko, S. Lobanov, A. Titov, J.-J. Zondy, *Appl. Phys. B: Laser Opt.* **2004**, *78*, 543–546.
- [2] B. Krebs, *Angew. Chem.* **1983**, *95*, 113–134.
- [3] D. Müller, F. E. Poltmann, H. Hahn, *Z. Naturforsch.* **1974**, *29b*, 117–118.
- [4] a) M. D. Ward, E. A. Pozzi, R. P. Van Duyne, J. A. Ibers, *J. Solid State Chem.* **2014**, *212*, 191–196; b) F. Q. Huang, B. Deng, D. E. Ellis, J. A. Ibers, *J. Solid State Chem.* **2005**, *178*, 2128–2132; c) S. J. Ewing, D. I. Woodward, A. V. Powell, P. Vaqueiro, *J. Solid State Chem.* **2013**, *204*, 159–165; d) B. Eisenmann, A. Hofmann, *Z. Kristallogr.* **1991**, *197*, 171–172; J. Kim, T. Hughbanks, *J. Solid State Chem.* **2000**, *149*, 242–251; e) P. Lemoine, D. Carre, M. Guittard, *Acta Crystallogr., Sect. C* **1984**, *40*, 910–912; f) K. Feng, D. Mei, L. Bai, Z. Lin, J. Yao, Y. Wu, *Solid State Sci.* **2012**, *14*, 1152–1156; B. Eisenmann, A. Hofmann, *Z. Kristallogr.* **1991**, *195*, 318; g) P. Wang, X.-Y. Huang, Y.-L. Liu, J. Li, H.-Y. Guo, *Huaxue Xuebao* **2000**, *58*, 1005–1008; h) A. Kumari, K. Vidyasagar, *Acta Crystallogr., Sect. E* **2005**, *61*, i193–i195; i) E. J. Wu, M. A. Pell, T. M. Fuelberth, J. A. Ibers, *Z. Kristallogr.* **1997**, *212*, 91; j) H.-Y. Zeng, F.-K. Zheng, R.-P. Chen, Z.-C. Dong, G.-C. Guo, J.-S. Huang, *J. Alloys Compd.* **2007**, *432*, 69–73; k) H. Lin, J.-N. Shen, L. Chen, L.-M. Wu, *Inorg. Chem.* **2013**, *52*, 10726–10728.
- [5] J. A. A. Ketelaar, W. H. t'Hart, M. Moerel, D. Polder, *Z. Kristallogr.* **1939**, *101*, 396–405.
- [6] J. W. Boon, C. H. Mac Gillavry, *Recueil Trav. Chim. Pays-Bas* **1942**, *61*, 910–920.
- [7] a) D. Friedrich, M. Schlosser, A. Pfitzner, *Cryst. Growth Des.* **2016**, *16*, 3983–3992; b) D. Friedrich, M. Schlosser, R. Weihrich, A. Pfitzner, *Inorg. Chem. Frontiers* **2017**, *4*, 393–400.
- [8] L. Kienle, V. Duppel, A. Simon, M. Schlosser, O. Jarchow, *J. Solid State Chem.* **2004**, *177*, 6–16.
- [9] L. Palatinus, G. Chapuis, *J. Appl. Crystallogr.* **2007**, *40*, 786–790.
- [10] V. Petříček, M. Dušek, L. Palatinus, *Z. Kristallogr., Crystalline Mater.* **2014**, *229*, 345.
- [11] A. Spek, *J. Appl. Crystallogr.* **2003**, *36*, 7–13.
- [12] Rigaku Oxford Diffraction, CrysAlisPro, Version 1.171.38.43, **2015**.
- [13] R. D. Shannon, *Acta Crystallogr., Sect. A* **1976**, *32*, 751–767.
- [14] D. Friedrich, M. Schlosser, A. Pfitzner, *Z. Anorg. Allg. Chem.* **2014**, *640*, 826–829.
- [15] J. Do, M. G. Kanatzidis, *Z. Anorg. Allg. Chem.* **2003**, *629*, 621–624.
- [16] a) H. Kisch, *Angew. Chem. Int. Ed.* **2013**, *52*, 812–847; b) P. Kubelka, F. Munk, *Z. Technol. Phys.* **1931**, 593.
- [17] J. P. Perdew, K. Burke, M. Ernzerhof, *Phys. Rev. Lett.* **1996**, *77*, 3865–3868.
- [18] J. W. Negele, *Phys. Rev. C* **1970**, *1*, 1260–1321.
- [19] D. Friedrich, F. Pielhofer, M. Schlosser, R. Weihrich, A. Pfitzner, *Chem. Eur. J.* **2015**, *21*, 1811–1817.
- [20] D. Friedrich, M. Schlosser, A. Pfitzner, *Cryst. Growth Des.* **2017**, *17*(9), 4887–4892.
- [21] T. Ishii, N. Kambe, *J. Cryst. Growth* **1986**, *76*, 489–493.
- [22] R. Suhrmann, K. Clusius, *Z. Anorg. Allg. Chem.* **1926**, *152*, 52–58.
- [23] Stoe & Cie GmbH, WinX^{POW}, Version 3.10, Darmstadt, Germany, **2017**.
- [24] K. Koepnik, H. Eschrig, *Phys. Rev. B* **1999**, *59*, 1743–1757.

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