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Electron Transfer around a Molecular Corner

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Abstract: The distance dependence of electron transfer (ET) is commonly investigated in linear rigid rod-like compounds, but studies of molecular wires with integrated corners imposing 90° angles are very rare. By using spirobifluorene as a key bridging element and by substituting it at different positions, two isomeric series of donor-bridge-acceptor compounds with either nearly linear or angled geometries were obtained. Photoinduced ET in both series is dominated by rapid through-bond hole hopping across oligofluorenes over distances of up to 70 Å. Despite considerable conformational flexibility, direct through-space and through-solvent ET is negligible even in the angled series. The independence of the ET rate constant on the total number of fluorene units in the angled series is attributed to a rate-limiting tunneling step through the spirobifluorene corner. This finding is relevant for multi-dimensional ET systems and grids in which individual molecular wires are interlinked at 90° angles.

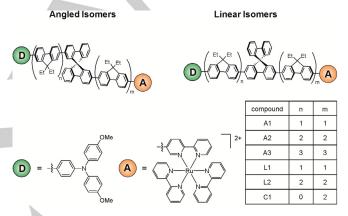
Photoinduced electron transfer has been widely explored in artificial and natural systems. Whilst numerous studies focused on the influence of driving-force and reorganization energy on electron transfer rates,^[1] many other investigations concentrated on their distance dependence.^[2] A key goal was often the generation of long-lived charge-separated states and the mimicry of elementary processes in natural photosynthesis.^[3] Linear rigid rod-like donor-acceptor compounds and straight molecular wires were frequently preferred,[4] because this permits the investigation of through-bond electron transfer over well-defined donor-acceptor distances and a separation of charges over maximal distance. By contrast, rigid molecular bridges or wires containing 90° angles are very uncommon.^[5] Prior investigators employed octahedral metal complexes as photosensitizers with donor and acceptor units attached to different ligands that were coordinated at ca. 90° angles relative to one another,^[6] but this is conceptually much different from the situation in which a corner is integrated into the molecular wire.

We hypothesized that spirobifluorene as a key bridging unit between a donor and an acceptor could provide access to linear and angled isomers of rigid rod-like donor-acceptor dyads (Scheme 1). When complemented with additional fluorene bridging units on both sides of the central spirobifluorene element, donor and acceptor branches of variable lengths should be accessible (Scheme 1). Both branches would then have a rigid rod-like structure, but depending on the substitution pattern at the spirobifluorene unit either linear or angled isomer series should be accessible. Electron delocalization across spiro-centers has been investigated in mixed-valence

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compounds,^[7] but we are unaware of prior studies that explored long-range electron transfer in spirobifluorene-based systems with both linear and angled geometries. This comparison should provide direct insight into the relative efficiencies of different charge transfer pathways, and this is relevant in the greater context of a future molecular electronics technology.^[8] Both linear and π -stacked oligo-fluorenes were previously identified as molecular wires that mediate long-range electron and energy transfer very efficiently.^[9] Spirobifluorene plays a key role in hole transport materials for solar cells,^[10] and it has also been explored in the context of luminescent materials.^[11]



Scheme 1. Molecular dyads investigated in this work. C1 has the same angled structure as A1, A2 and A3, but with different relative bridge lengths between donor and acceptor branches. The linking position of the donor and the acceptor is the same in all 6 compounds.

A triarylamine (TAA) unit and a $[Ru(bpy)_3]^{2^+}$ (bpy = 2,2'bipyridine) complex (Ru) were chosen as a donor and an acceptor, respectively, because this combination has favorable optical spectroscopic properties and there is enough drivingforce for long-range electron transfer after photoexcitation of the metal complex, at least in so-called flash-quench experiments (see below). For the linear series of dyads, we managed to synthesize compounds with either 3 or 5 fluorene units between TAA and Ru (L1, L2). For the angled series, dyads with 3, 5, and 7 fluorene units between TAA and Ru could be synthesized and explored (A1 – A3). Synthetic details and characterization data are in the Supporting Information (SI pages S2-S30).

Not too surprisingly, *in vacuo* Merck Molecular Force Field (MMFF) energy minimization calculations quickly dispersed the notion of the angled series having a strictly rigid rod-like right-angled structure. Instead, two low energy conformers can be readily identified: an open 'winged' conformer, and a closed conformer in which the TAA and $[Ru(bpy)_3]^{2+}$ units approach one another, with N_{TAA}-Ru distances of 7.6, 8.9 and 9.0 Å for A1, A2 and A3, respectively (SI page S40-S41). This second conformer is predicted to be lower in energy by 8 kJ mol⁻¹ for A1, 26 kJ mol⁻¹ for A2, and 27 kJ mol⁻¹ for A3, compared to the open conformer (SI page S41). MMFF is a low, computationally inexpensive level

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of theory and these stabilization values are therefore likely exaggerated by the calculations.^[12] Furthermore, as the calculations are performed *in vacuo*, solvent stabilization effects are not taken into consideration. The key message from these calculations is that in the angled isomer series conformers with relatively short through-space but long through-bond distances can form. Electron transfer in such systems can either proceed through the long bonding system, or across the short space between donor and acceptor.^[13] In the linear series, by contrast, only the through-bond pathway is available. To further separate the role of through-bond and through-space electron transfer, an additional angled control compound, C1, was prepared, in which a close contact between the donor and acceptor cannot form (Scheme 1).

The TAA and [Ru(bpy)₃]²⁺ redox potentials are very similar in the 6 dyads (SI pages S32-S34), leading to essentially identical driving-forces for intramolecular electron transfer in all cases. For photoinduced electron transfer from TAA to ³MLCT-excited $[Ru(bpy)_3]^{2^+}$ the reaction free energy is relatively small $({\Delta G_{\text{ET}}}^0 \approx$ -0.2 eV), making the observation of long-range charge transfer challenging because the rates can get slow compared to the ³MLCT lifetime. However, in presence of excess methylviologen (MV^{2+}) , the ³MLCT-excited complex is oxidatively guenched to yield $[Ru(bpy)_3]^{3+}$ within less than 10 ns (SI page S35),^[14] and the Ru(III) species is then reduced by TAA with substantially greater driving-force ($\Delta G_{\text{ET}}^0 \approx$ -0.7 eV). Bimolecular thermal reverse electron transfer from MV⁺ to TAA⁺ subsequently takes places on a longer timescale (~ 40 µs, SI page S35).^[14] This socalled flash-quench technique was most useful to explore longrange electron transfer in our dyads.

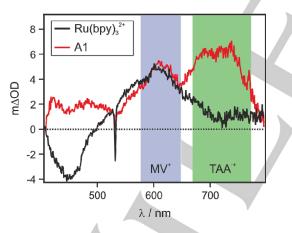


Figure 1. Transient absorption spectra obtained after excitation at 532 nm with laser pulses of ca. 10 ns duration: 15 μ M dyad A1 with 80 mM MV²⁺ (red trace), and 15 μ M [Ru(bpy)₃]²⁺ with 80 mM MV²⁺ (black trace) in de-aerated CH₃CN at 22 °C.

All dyads exhibit intense fluorene-localized π - π^* absorption bands between 370 and 385 nm, but at 532 nm photoexcitation occurs selectively into MLCT absorptions of the photosensitizer units (SI page S36). The transient absorption spectrum of A1 recorded in presence of MV²⁺ (red trace Figure 1) is typical for all 6 dyads (SI page S37). In all cases, the dyad concentration was 15 μ M, whilst MV²⁺ was present at 80 mM in de-aerated

CH₃CN at room temperature. The duration of the laser pulses was ca. 10 ns, and transient absorption spectra were recorded by time-integration over a period of 200 ns immediately after excitation. The key observation is the transient absorption band at 730 nm, which is diagnostic for TAA⁺ (SI page S38),^[15] in addition to a band at 605 nm caused by MV^{+. [16]} Evidently, selective excitation of the Ru sensitizer at 532 nm leads to the formation of both TAA⁺ and MV⁺, and this can be explained by rapid oxidative ³MLCT excited-state quenching by MV²⁺ followed by intramolecular electron transfer from TAA to the [Ru(bpy)₃]³⁺ unit, as noted above. When performing the same experiment with the commercial [Ru(bpy)₃]²⁺ reference complex, one still observes the MV⁺ band at 605 nm (black trace in Figure 1), but instead of the TAA⁺ absorption at 730 nm there is now an MLCT bleach at 450 nm signaling the disappearance of $[Ru(bpy)_3]^{2+}$ and the formation of $[Ru(bpy)_3]^{3+}$, in line with expectation.

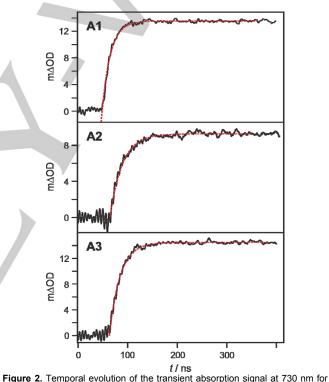


Figure 2. Temporal evolution of the transient absorption signal at 730 nm for compounds A1 – A3 following excitation at 532 nm with laser pulses of ca. 10 ns duration. The dyad and methylviologen concentrations were 15 μ M and 80 mM, respectively, in all cases. The solvent was de-aerated CH₃CN at 22 °C. The dashed red lines are single-exponential fits.

Given these spectral characteristics, the rate constants for intramolecular electron transfer (k_{ET}) from TAA to [Ru(bpy)₃]³⁺ in the 6 dyads can be determined by monitoring the temporal evolution of the transient absorption signal at 730 nm. In the linear dyads L1 and L2, and in the control angled dyad C1, the risetime of the signal at 730 nm is instrumentally limited, and we can only estimate a lower limit of 10^8 s^{-1} for k_{ET} (SI, page S39). For the angled dyads A1, A2, and A3, risetimes of 17 ± 3 , 25 ± 3 , and 21 ± 3 ns, respectively, are measured (Figure 2, Table 1). Thus, electron transfer is faster in the linear than in the angled isomers, but this is not particularly surprising because the former

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have a more extensively π -conjugated bridging system than the latter. The most significant observations are clearly the distance independence of k_{ET} along the series of angled isomers (A1, A2, A3) and the faster electron transfer in the control compound C1 compared to A1, A2 and A3 (Table 1).

Table 1. Through-bond (r_{tb}) and through-space (r_{ts}) distances between the triarylamine N-atom and the Ru-atom in the 6 dyads from Scheme 1, along with time (τ_{ET}) and rate constants (k_{ET}) for electron transfer between TAA and photogenerated Ru(III). The through-space distances represent the lowest energy close-contact conformation of donor and acceptor calculated for geometry optimized structures using the MMFF level of theory (SI page S40).

| dyad | r _{tb} / Å | r _{ts} / Å | τ_{ET} / ns | k _{ET} / s ⁻¹ |
|------|---------------------|---------------------|-------------------------|-----------------------------------|
| L1 | 35.9 | | < 10 | > 10 ⁸ |
| L2 | 52.8 | | < 10 | > 10 ⁸ |
| A1 | 38.6 | 7.6 | 17±3 | (5.9±1.2)·10 ⁷ |
| A2 | 55.4 | 8.9 | 25±3 | (4.0±0.5)·10 ⁷ |
| A3 | 72.3 | 9.0 | 21±3 | (4.8±0.8)·10 ⁷ |
| C1 | 38.6 | 19.5 | < 10 | > 10 ⁸ |

Prior studies on linear oligo-fluorene bridged donor-acceptor systems already reported very shallow distance dependences of through-bond electron transfer rates.^[9a, 9b] However, the distance dependence of k_{ET}, commonly captured by the so-called distance decay constant (β), is not a bridge-specific property but rather it is governed by the entire combination of donor, bridge, and acceptor.^[17] It is therefore difficult to make a general statement about the distance dependence of k_{ET} for a given bridging structure. However, the hopping mechanism identified previously in some linear fluorene-bridged systems is likely to be operative in our dyads as well, [9a, 9b] because the oxidation potentials of the Ru photosensitizer and the oligo-fluorenes are very close to one another, ca. 1.3 V vs. SCE in both cases (SI page S33). Consequently, after photo-generation of $[Ru(bpy)_3]^{3+}$, hole transfer to the molecular bridge is expected to be possible, and such a hopping mechanism is in line with the distance insensitivity of k_{ET} observable in the angled series.^[9a, 18] However, for the angled dyads A1 - A3, through-space or through-solvent charge transfer pathways between TAA and [Ru(bpy)₃]³⁺ should also be considered, especially in view of the MMFF results discussed above.

The control compound C1 was specifically designed to probe the possibility of through-space electron transfer, because it cannot adopt conformations with short TAA-[Ru(bpy)₃]²⁺ contacts (Figure S9 on page S41) due to the different lengths of the donor and acceptor branches (n = 0, m = 2, Scheme 1). According to MMFF calculations the shortest possible through-space N_{TAA}-Ru distance in C1 is 19.5 Å. If direct through-space tunneling from the donor to the acceptor were to contribute significantly to k_{ET} in any of the angled isomers (A1, A2, A3), one would expect the charge transfer rate to decrease for compound C1 (see SI page S42 for further details). The opposite phenomenon is instead observed (k_{ET} > 10⁸ s⁻¹ for C1, Table 1 and SI page S39), and thus it seems plausible to conclude that

through-space tunneling plays only a minor role in A1, A2 and A3 if any at all. This is not overly surprising, as the flash-quench measurements were undertaken in acetonitrile, a highly polar solvent that can efficiently stabilize charges and therefore minimize the energy difference between open and closed conformers compared to what is obtained with the (crude) in vacuo MMFF calculations (see above and SI page S40-S41). For example, we recently reported electron transfer across a close-contact ion-pair which was significantly disrupted by even modestly polar solvents (CH2Cl2, THF).[19] Given that the coulombic attraction in an ion-pair is significantly greater than for an ion-dipole system such as presented herein, it is unlikely that enough molecules exist in a close-contact conformation in acetonitrile. Electron tunneling through solvent is very inefficient compared to through-bond hopping. For example, the distance decay constant (β) for tunneling through 2-methyl-THF is 1.6 Å⁻ ¹,^[20] whilst for hopping processes the (phenomenological) β values are usually below 0.1 Å^{-1,[18]} Thus, through-space and through-solvent tunneling seem to play a negligible role in our dyads, and through-bond pathways seem to be dominant.

The greater rate of electron transfer in C1 compared to A1, A2 and A3 further gives implications for the mechanism of charge transfer. The linear series of complexes (L1, L2) clearly demonstrates that the hopping mechanism of hole transfer across fluorene bridges is extremely rapid. In the angled series however, the hole must tunnel through the spirobifluorene corner unit. This is likely to be the rate-limiting step of the hole transfer despite some spiroconjugation,^[21] and is corroborated by the more rapid rate of charge transfer in C1, which has the TAA electron donor located much closer to the spirobifluorene unit than any of the angled dyads A1, A2 and A3, leading to stronger coupling between TAA and electronic spirofluorene. Consequently, the tunneling rate through the spirobifluorene corner is expected to increase in C1, as observed experimentally. However, hole tunneling between Ru(III) and its neighboring fluorene unit (rather than electron tunneling between TAA and the spirobifluorene corner) is likely to remain the initial key step of long-range charge transfer between Ru(III) and TAA (see SI pages S42-S43 for further details).

In summary, charge transfer in both the linear and angled dyads seems to be dominated by through-bond (hole) hopping mechanisms over distances of up to 70 Å. A lack of dependence of the electron transfer rate on the number of fluorene bridging units in the angled isomers is attributed to a rate-limiting hole tunneling step through the spirobifluorene corner. Despite the possible formation of conformers with relatively short donoracceptor contacts as a consequence of the relatively flexible long fluorene bridges, no evidence for direct through-space or through-solvent tunneling was observed, presumably due to the ability of acetonitrile to disrupt the ion-dipole attraction that would lead to such phenomena and the comparatively large β values for tunneling through solvent.^[2g, 20] Thus, the kinetics of charge transfer occurring over 40-70 Å can be governed by a single molecular corner unit. This key insight is relevant for the construction of multi-dimensional electron transfer systems or molecular grids [22] in which individual molecular wires are connected by spirobifluorene or other corner-forming units.^[23]

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The kinetics of hole hopping over 40-70 Å are governed by a ratedetermining tunneling step across a spirobifluorene corner unit. This finding is relevant for the construction of multi-dimensional electron transfer systems or grids in which individual molecular wires are interlinked at 90° angles. Hauke C. Schmidt, Christopher B. Larsen, Oliver S. Wenger*

Page No. – Page No.

Electron Transfer around a Molecular Corner

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