

Bulky Phenylalkyl Substitutions to Bisthienoisatins and Thienoisindigos

Dongho Yoo,* Akihiro Kohara, Minoru Ashizawa,* Tadashi Kawamoto, Hiroyasu Masunaga, Noboru Ohta, Hidetoshi Matsumoto, and Takehiko Mori*

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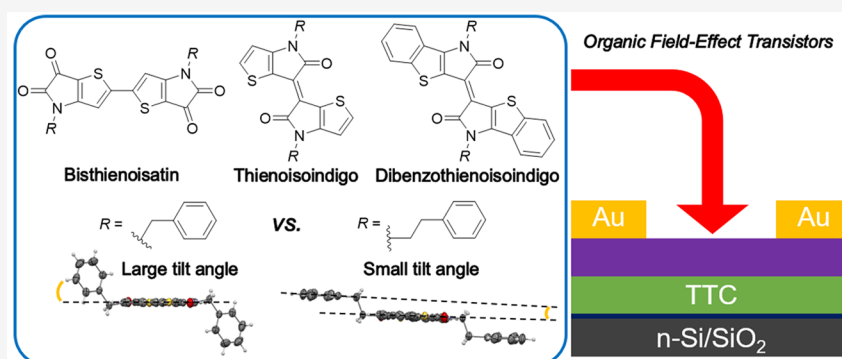
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ABSTRACT: In order to investigate the effects of bulky substituents on the crystal structures and packing modes, *N*-benzyl (Bn) and *N*-2-phenylethyl (EtPh) substituents are introduced in bisthienoisatin (BTI), thienoisindigo (TIIG), and dibenzothienoisindigo (DBTII). These molecules maintain uniform stacking structures, though EtPh-BTI has a two-dimensional slipped-herringbone structure. The benzyl groups are largely distorted from the molecular core, and thin films of Bn-TIIG and Bn-DBTII show poor quality due to the two kinds of molecular orientations. In contrast, the 2-phenylethyl-substituted molecules enable suitable molecular packing owing to the ethylene spacer and show relatively good thin-film qualities as well as much improved transistor properties. The BTI derivatives show only electron transport, but other compounds exhibit ambipolar transistor properties. In particular, EtPh-TIIG and EtPh-DBTII show maximum hole mobilities of about 0.04–0.05 cm² V^{−1} s^{−1} together with moderate electron mobilities.

■ INTRODUCTION

The development of organic semiconductors (OSCs) has drastically progressed in the recent 30 years for applications to organic field-effect transistors (OFET),¹ organic photovoltaics,² and organic light-emitting diodes.³ Among them, OFETs have been investigated continuously for a long period of time.⁴ From the beginning of research in OFETs, p-type OSCs have been extensively designed starting from pentacene,⁵ and the mobility has been improved from an order of 10^{−5} to over 10 cm² V^{−1} s^{−1}.⁶ The development of n-type and ambipolar OSCs is lagging in comparison to p-type OSCs, but the research trend has largely shifted to these OSCs, and many promising small molecules and novel polymer semiconductors have achieved electron mobilities of up to 10 cm² V^{−1} s^{−1}.^{7–12} Recently, a great number of studies have been devoted to process and interface engineering in order to improve the mobility through the device fabrication method.¹³ In addition, molecular engineering has been investigated by modifying molecular structures through the substitution of functional groups and the extension of π skeletons of promising OSCs showing high mobility.^{13,14} An exploration of high mobility has

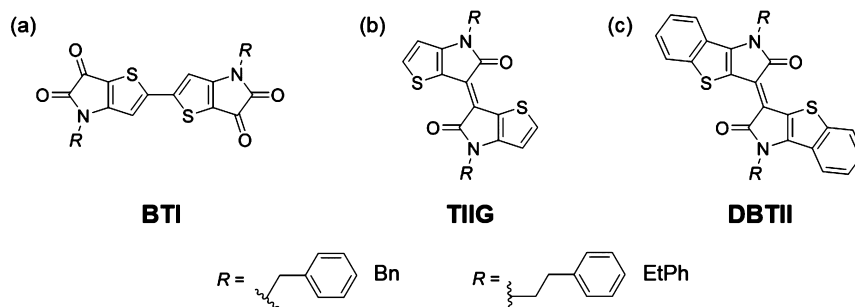
been attempted regardless of the transporting carrier types.¹⁵ One of the ways to achieve high mobility is to control the molecular packing by introducing bulky substituents. In some previous works, the introduction of bulky substituents has been found to be effective in improving the OFET properties in OSCs such as pentacene,^{16,17} [1]benzothieno[3,2-*b*][1]-benzothiophene (BTBT),^{18,19} naphthalenediimide (NDI),^{20,21} and 5,5'-bithiazolidinylidene-2,4,2',4'-tetrathione.²² The bulky substituents ordinarily produce steric hindrance to destroy the intermolecular interaction, but in some cases, bulky groups realize molecular packings favorable to charge transport such as herringbone and brickwork structures.

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Scheme 1. Chemical Structures of (a) BTI, (b) TIIG, and (c) DBTII, Where R = Benzyl (Bn) and 2-Phenylethyl (EtPh)



Previously, we have prepared bisthienoisatin (BTI; Scheme 1a) by the oxidative coupling of thienoisatins. This molecule has a linear planar structure and shows n-type OFET properties owing to the four electron-withdrawing carbonyl groups.²³ Also, we have prepared N-unsubstituted thienoisindigo (TIIG; Scheme 1b) by the dimerization of thienoisatins, which has a planar molecular structure due to an intramolecular oxygen–sulfur (S⋯O) interaction.²⁴ In addition, TIIG forms intermolecular sulfur–sulfur (S⋯S) interactions and hydrogen bonds with adjacent molecules, so that this molecule shows moderate ambipolar OFET properties.²⁴ The derivatives with N-alkyl groups also have been investigated extensively as a repeating unit of polymers and small-molecule OSCs because the TIIG backbone is a promising skeleton.^{25–33} Dibenzothienoisindigo (DBTII; Scheme 1c), in which benzene rings are condensed with TIIG, is also a planar molecule and exhibits p-dominant ambipolar OFET properties on a tetratetracontane (TTC) modified substrate, showing much improved mobilities due to the extended π skeleton.²⁵ However, the above three kinds of compounds have been reported only in the form with straight and branched N-alkyl chains except for N-unsubstituted TIIG. This has prompted us to introduce bulky phenylalkyls to BTI, TIIG, and DBTII cores for the purpose of investigating the effects of bulky substituents on the crystal structures and packing modes of molecules. In previous reports, a 2-phenylethyl group bearing an ethylene spacer has been effective in constructing a preferable molecular packing for carrier transport.^{21,22} Thus, we have prepared N-2-phenylethyl (EtPh)-substituted BTI, TIIG, and DBTII and investigated the correlation between the crystal structures and the OFET properties. For comparison with the 2-phenylethyl-substituted molecules, N-benzyl (Bn)-substituted molecules have also been prepared (Scheme 1).

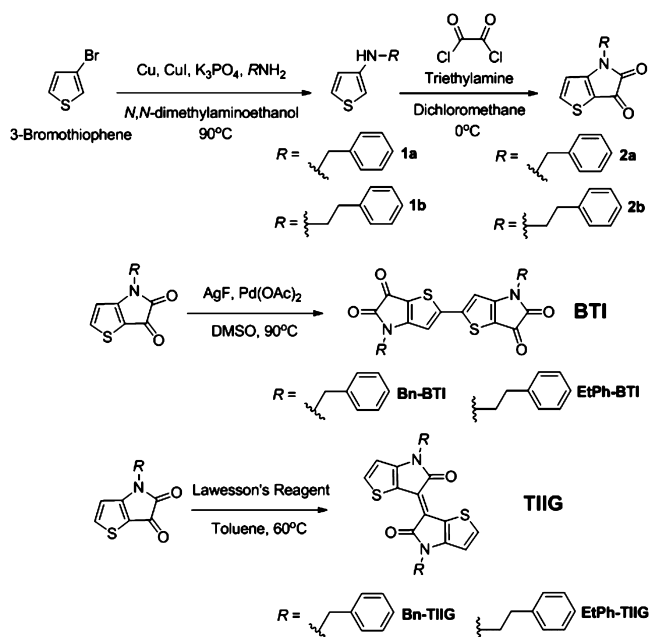
EXPERIMENTAL SECTION

Syntheses. BTI was synthesized from 3-bromothiophene via introduction of phenylalkyl (1), cyclization to thienoisatin (2), and the AgF-mediated oxidative coupling of 2 with Pd(OAc)₂ as a catalyst (Scheme 2).²³ TIIG was synthesized from 2 by a homocoupling reaction using Lawesson's reagent (Scheme 2). DBTII was prepared by the same route as for TIIG from 3-bromobenzo[b]thiophene (Scheme 3).

Crystal Preparation. Single crystals of EtPh-BTI (260 °C), Bn-TIIG (230 °C), Bn-DBTII (320 °C), and EtPh-DBTII (310 °C) were obtained from sublimation under an N₂ atmosphere (flow rate, 50 mL/min; pressure, 10 Pa). Crystals of Bn-BTI were obtained from a chloroform solution, and single crystals of EtPh-TIIG were obtained from a toluene solution.

X-ray Diffraction. Single-crystal X-ray structure analyses were carried out by a Rigaku R-Axis RAPID II imaging plate

Scheme 2. Synthetic Scheme to BTI and TIIG



diffractometer with Cu K α radiation from a rotating anode source with a confocal multilayer X-ray mirror (Rigaku VM-spider, $\lambda = 1.54187$ Å). The structures were solved by the dual-space method (SHELXT)³⁴ and refined by the full-matrix least-squares method by applying anisotropic temperature factors for all non-hydrogen atoms using the SHELXL programs.³⁵ The hydrogen atoms were placed at geometrically calculated positions. A twin analysis of Bn-BTI was conducted by PLATON because crystals of Bn-BTI have non-merohedral twins.³⁶

The thin-film X-ray diffraction (XRD) patterns were taken by using a Philips X'Pert-Pro-MRD instrument with Cu K α radiation ($\lambda = 1.540598$ Å). Two-dimensional grazing incidence wide-angle X-ray scattering (GIWAXS) patterns were obtained at SPring-8 on beamline BL40B2.³⁷ The samples were irradiated at a fixed incidence angle on the order of 0.125°, and the GIWAXS patterns were recorded with a 2-D image detector (PILATUS3 S 2M). The wavelength of the X-ray beam was 1 Å (energy of 12.39 keV), and the camera length was 350.4 mm.

RESULTS AND DISCUSSION

Electronic Properties. The highest occupied molecular orbital (HOMO) and lowest unoccupied molecular orbital (LUMO) levels are estimated by using cyclic voltammetry and ultraviolet visible (UV–vis) absorption spectroscopy. Bn-BTI and EtPh-BTI show two reversible reduction waves similar to those of the previous alkyl-BTI (Figure 1a),²³ whereas Bn-DBTII and EtPh-DBTII exhibit two reversible oxidation

Scheme 3. Synthetic Scheme to DBTII

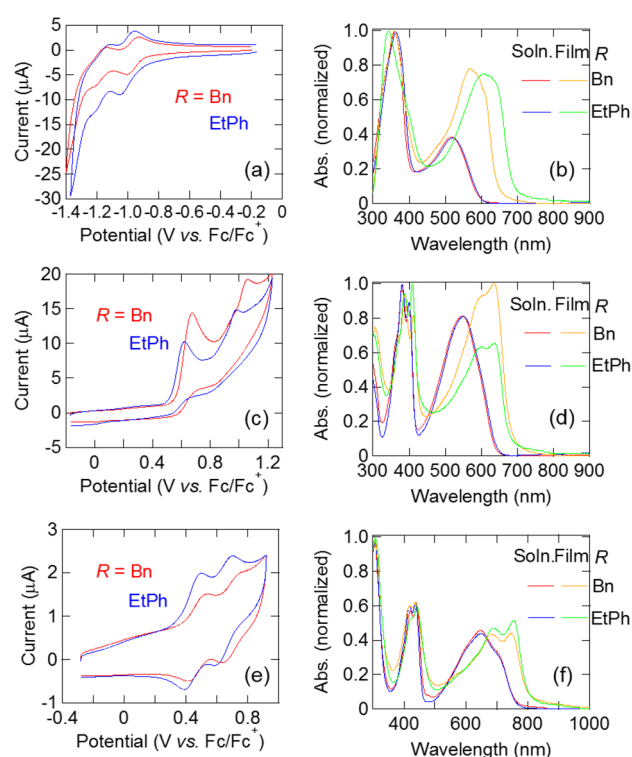
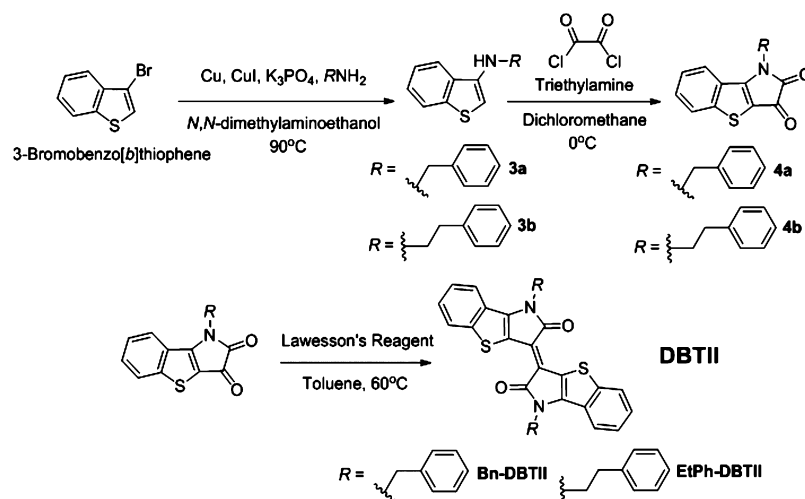


Figure 1. Cyclic voltammograms of (a) BTIs, (c) TIIGs, and (e) DBTIIs. UV-vis absorption spectra of (b) BTIs, (d) TIIGs, and (f) DBTIIs.

waves (Figure 1e). However, Bn-TIIG and EtPh-TIIG show two irreversible oxidation peaks (Figure 1c); thus the HOMO

levels are estimated from the onset potentials (E_{onset}). In other molecules, the HOMO or LUMO levels are estimated from the first half-wave oxidation (DBTIIs) or reduction (BTIs) potentials ($E_{1/2}$) as given in Table 1. The HOMO–LUMO gaps (E_g^{opt}) are extracted from λ_{onset} values of the solution absorption spectra. The LUMO levels of BTIs are about -3.8 eV, which denotes relatively low lying LUMO levels due to the four electron-withdrawing carbonyl groups. TIIGs and DBTIIs are expected to show ambipolar OFET properties because the LUMO levels are lower than -3.15 eV and the HOMO levels are higher than -5.6 eV.³⁸ The absorption spectra of the thin films are red-shifted owing to the intermolecular interactions in the solid state (Figures 1b,d,f). In the thin films of TIIGs and DBTIIs, the absorption bands are almost identical regardless of the substituents (Table 2). However, the absorption band of

Table 2. Optical Properties of the Thin Films

compd	R	λ_{onset} (nm)	λ_{max} (nm)	E_g^{opt} (eV)
BTI	Bn	646	359, 569	1.92
	EtPh	687	343, 607	1.80
TIIG	Bn	685	388, 412, 635	1.81
	EtPh	679	388, 409, 636	1.82
DBTII	Bn	798	306, 439, 684, 747	1.55
	EtPh	799	306, 443, 690, 754	1.55

EtPh-BTI is shifted to long wavelength much more than for Bn-BTI. This means intermolecular interactions are strong in EtPh-BTI in comparison with Bn-BTI (see Crystal Structures). Here, the excited energy levels split greatly, resulting in a narrow energy gap in the solid state.

Table 1. Summary of Electrochemical and Optical Properties in the Solutions

compd	R	$E_{1/2}$ (V)	E_{onset} (V)	E_{LUMO} (eV)	λ_{onset} (nm)	λ_{max} (nm)	E_g^{opt} (eV)	E_{HOMO} (eV)
BTI	Bn	−0.97		−3.83	594	360, 518	2.09	−5.92 ^a
	EtPh	−1.00		−3.80	600	364, 521	2.06	−5.86 ^a
TIIG	Bn		0.57	−3.43 ^b	639	380, 400, 548	1.94	−5.37
	EtPh		0.52	−3.39 ^b	643	381, 400, 551	1.93	−5.32
DBTII	Bn	0.48		−3.65 ^b	758	306, 421, 441, 647	1.63	−5.28
	EtPh	0.44		−3.60 ^b	756	306, 418, 438, 649	1.64	−5.24

^aEstimated by subtracting E_g^{opt} from E_{LUMO} . ^bEstimated by adding E_g^{opt} to E_{HOMO} .

Crystal Structures. In all crystals, the molecules are centrosymmetric, and the half-molecule is crystallographically independent. **Bn-DBTII** crystallizes in a triclinic system, and the others crystallize in monoclinic systems. All molecules in the monoclinic structures form stacks in which the molecules are oppositely tilted from the molecules in the adjacent columns (**EtPh-BTI**, **Bn-TIIG**, and **EtPh-DBTII**) or from the adjacent layers (**Bn-BTI** and **EtPh-TIIG**), although all molecules are parallel in **Bn-DBTII**. In general, the molecular cores are perfectly planar, but the phenyl rings are tilted, exceeding 70° in **Bn** and less than 22° in **EtPh** (Figure 2a,e, Figure 3a,e, Figure 4a,e, and Table 7). Crystallographic data are given in Table 3.

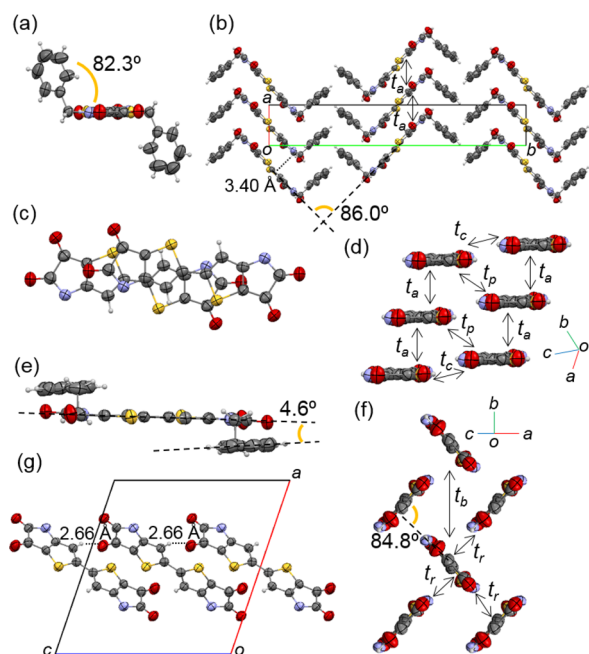


Figure 2. (a) Molecular structure of **Bn-BTI**. (b) Crystal structure of **Bn-BTI** viewed along the c axis. (c) Molecular overlap mode. (d) Crystal structure of **Bn-BTI** viewed along the long axis of the molecular core. (e) Molecular structure of **EtPh-BTI**. Crystal structure of **EtPh-BTI** viewed (f) along the long axis of the molecular core and (g) along the b axis. For clarity, the substituents are omitted in (c), (d), (f), and (g).

Bn-BTI forms a uniform stacking structure along the a axis with an interplanar spacing of 3.40 Å (Figure 2b). The stacked molecules are slipped mainly along the long axis of the molecular core (Figure 2c). Molecules in the same ac layer are parallel, but molecules in the adjacent layers are alternately tilted (86.0°); these layers are stacked along the b axis (Figure 2b). An $S\cdots S$ interaction with a distance of $3.511(2)$ Å and an $S\cdots O$ intermolecular interaction with a distance of $3.263(4)$ Å are observed along the c axis (Figure S5c). Accordingly, the transfer integrals (t_c and t_p) in the transverse direction are comparable to t_a along the stacking axis (Figure 2d and Table 4); these transfer integrals suggest a comparatively isotropic conduction, though the magnitude is not large.

The crystal structure of **EtPh-BTI** looks like a herringbone structure (Figure 2f), but the molecules are largely slipped along the long axis of the molecular core (Figure 2g and Figure S6b). This structure involves side by side intermolecular interactions (Figure 2f), and hydrogen bonds with a distance

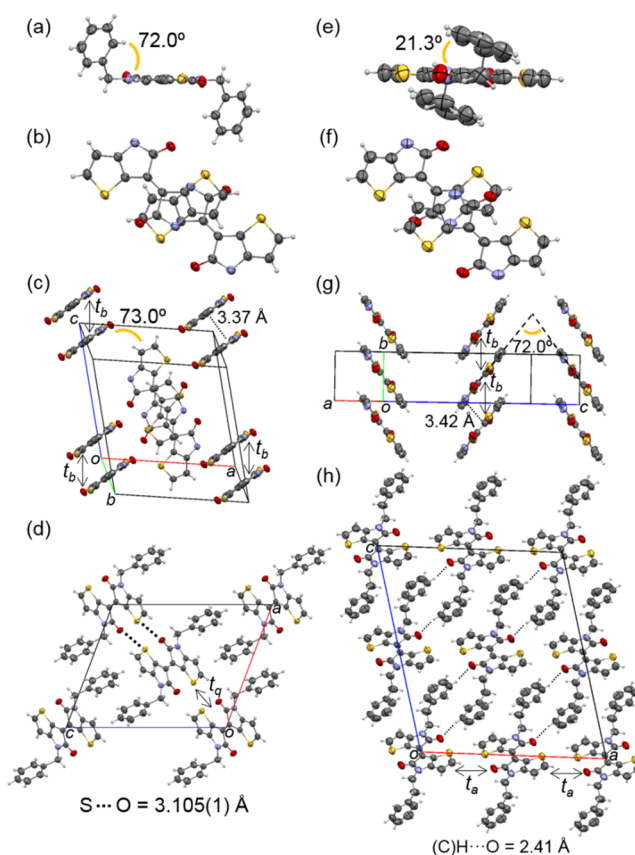


Figure 3. (a) Molecular structure of **Bn-TIIG** and (b) molecular overlap mode. (c) Molecular packing of **Bn-TIIG** and (d) crystal structure viewed along the b axis. (e) Molecular structure of **EtPh-TIIG** and (f) molecular overlap mode. Crystal structure of **EtPh-TIIG** viewed (g) along the molecular long axis and (h) along the b axis. For clarity, the substituents are omitted in (b), (c), (f), and (g).

of 2.66 Å are formed in $C=O\cdots H$ along the c axis (Figure 2g). These interactions are the origins of the largely red shifted solid-state absorption bands in comparison with **Bn-BTI**. The dihedral angle between the adjacent columns is nearly perpendicular (84.8°) (Figure 2f). Nonetheless, the large displacement along the long axis of the molecular core is disadvantageous to charge transport in comparison to the ideal herringbone structure,³⁹ because the transfer integrals are comparatively small (Table 4).

The crystal structure of **Bn-TIIG** is isostructural with N-unsubstituted **TIIG**.²⁴ **Bn-TIIG** forms a uniform stacking structure along the b axis with an interplanar spacing of 3.37 Å (Figure 3c). In **Bn-TIIG** and **EtPh-TIIG**, the molecular long axis is defined by the substituents and is parallel to the N–N direction. This is the perpendicular direction in Scheme 1b, and the intercolumnar direction corresponds to the horizontal direction in Scheme 1b (molecular short axis). In **Bn-TIIG**, molecules in the adjacent columns along the $a + c$ axis are tilted to opposite directions (73.0°) (Figure 3c). In this direction, $S\cdots O$ intermolecular interactions with a distance of $3.105(1)$ Å are formed (Figure 3d); this is similar to the hydrogen bonds of N-unsubstituted **TIIG**.²⁴

EtPh-TIIG forms a uniform stacking structure along the b axis with an interplanar spacing of 3.42 Å (Figure 3g). The adjacent columns along the a axis are parallel (Figure 3h), whereas the molecules in the adjacent layers are alternately tilted by about 72.0° in opposite directions; the layers are

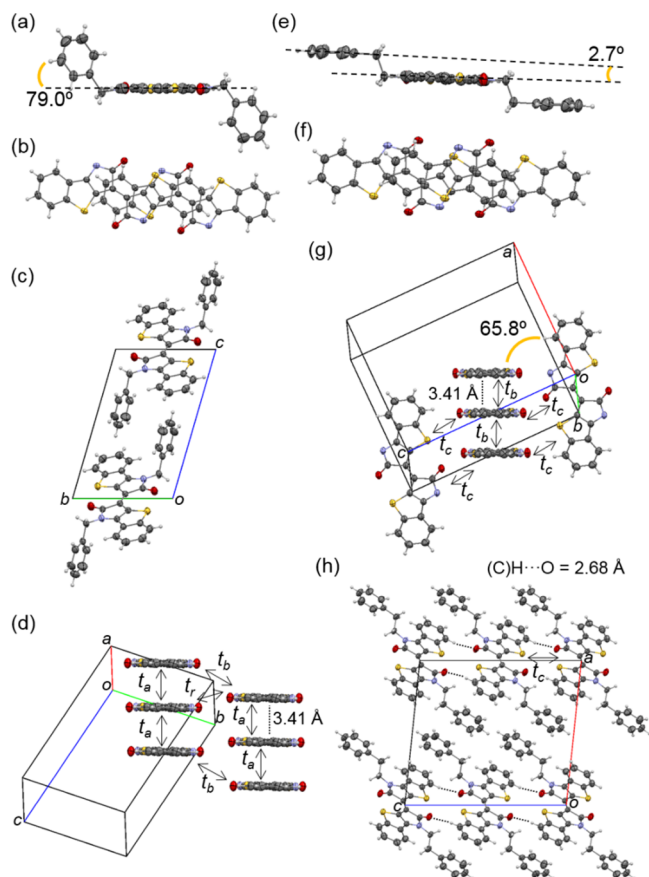


Figure 4. (a) Molecular structure of **Bn-DBTII** and (b) molecular overlap mode. Crystal structure of **Bn-DBTII** viewed (c) along the *a* axis and (d) along the long axis of the molecular core. (e) Molecular structure of **EtPh-DBTII** and (f) molecular overlap mode. (g) Molecular packing of **EtPh-DBTII** and (h) crystal structure viewed along the *b* axis. For clarity, the substituents are omitted in (b), (d), (f), and (g).

Table 4. Electron Transfer Integrals (meV) of BTIs

R	t_a	t_b	t_c	t_p	t_r
Bn	−32.9		−15.3	21.8	
EtPh		8.3			27.8

stacked along the *c* axis (Figure 3g). **EtPh-TIIG** has a large intramolecular torsion angle (21.3°) between the **TIIG** plane and the phenyl ring of **EtPh** in comparison to other **EtPh**-substituted molecules (Figure 3e). Between C=O of the molecular core and a phenyl hydrogen of **EtPh**, a hydrogen bond is formed (distance 2.41 Å) (Figure 3h); thus, this interaction indicates a relatively large intramolecular torsion angle.

In **Bn-TIIG**, the transfer integral along the stacking axis is larger for the LUMO than for the HOMO (Table 5). In

Table 5. Transfer Integrals (meV) of TIIGs

R		t_a	t_b	t_q
Bn	HOMO		−28.7	4.0
	LUMO		42.9	−14.3
EtPh	HOMO	−3.7	−118.1	
	LUMO	−3.4	67.2	

contrast, **EtPh-TIIG** has a particularly large transfer integral for the HOMO. This difference comes from the slip distance along the molecular short axis, though the slip distance along the molecular long axis is the same (Figure 3b,f and Figures S7 and S8). The large displacement in **Bn-TIIG** results in a ring-overatom structure (Figure 3b), though the comparatively small displacement in **EtPh-TIIG** results in a ring-overbond structure (Figure 3f).

In **Bn-DBTII** and **EtPh-DBTII**, the molecular long axis is defined by the molecular core and the substituents block the molecular sides. These molecules form uniform stacking structures with interplanar spacings of 3.41 Å along the *a* and *b* axes, respectively (Figure 4). In **Bn-DBTII**, all molecules are parallel, and an intercolumnar interaction exists along the *b*

Table 3. Crystallographic Data

	Bn-BTI	EtPh-BTI	Bn-TIIG	EtPh-TIIG	Bn-DBTII	EtPh-DBTII
chem formula	C ₂₆ H ₁₆ N ₂ O ₄ S ₂	C ₂₈ H ₂₀ N ₂ O ₄ S ₂	C ₂₆ H ₁₈ N ₂ O ₂ S ₂	C ₂₈ H ₂₂ N ₂ O ₂ S ₂	C ₃₄ H ₂₂ N ₂ O ₂ S ₂	C ₃₆ H ₂₆ N ₂ O ₂ S ₂
formula wt	484.54	512.60	454.56	482.61	554.68	582.73
cryst syst	monoclinic	monoclinic	monoclinic	monoclinic	triclinic	monoclinic
space group	<i>P</i> 2 ₁ / <i>n</i>	<i>P</i> 2 ₁ / <i>c</i>	<i>P</i> 2 ₁ / <i>n</i>	<i>C</i> 2/ <i>c</i>	<i>P</i> $\bar{1}$	<i>P</i> 2 ₁ / <i>c</i>
<i>a</i> (Å)	4.91789(9)	13.6380(3)	12.6820(3)	19.5477(6)	6.37399(12)	14.2197(3)
<i>b</i> (Å)	30.6056(6)	7.18687(16)	5.66961(13)	5.60960(19)	8.36881(15)	6.28460(11)
<i>c</i> (Å)	7.66360(14)	13.0238(3)	15.4096(4)	22.1929(7)	12.9434(2)	15.7764(3)
α (deg)	90	90	90	90	104.1072(7)	90
β (deg)	100.2339(9)	108.7680(13)	110.6364(12)	104.1203(17)	98.5380(7)	95.9370(7)
γ (deg)	90	90	90	90	100.3249(7)	90
<i>V</i> (Å ³)	1135.13(4)	1208.65(5)	1036.89(4)	2360.02(13)	645.27(2)	1402.30(4)
<i>Z</i>	2	2	2	4	1	2
total no. of rflns	13117	13557	11296	12872	7637	15470
no. of unique rflns (<i>R</i> _{int})	2062 (0.0576)	2214 (0.0353)	1899 (0.0552)	2162 (0.0328)	2332 (0.0498)	2575 (0.0407)
<i>D</i> _{calc} (g/cm ³)	1.418	1.408	1.456	1.358	1.427	1.380
<i>R</i> 1 ^a	0.0671	0.0486	0.0445	0.0516	0.0476	0.0403
<i>wR</i> 2 ^b	0.2239	0.1711	0.1194	0.1767	0.1318	0.1133
GOF	1.130	1.140	1.081	1.124	1.175	1.153
<i>T</i> (K)	298	275	275	275	298	298

^a*R*1($F^2 > 2\sigma(F^2)$). ^b*wR*2(F^2) for all observed reflections.

axis (Figure 4c). The phenyl group of **Bn-DBTII** stretches toward the *c* axis facing the phenyl group of the next layer molecule (Figure 4c). The molecular cores construct a one-dimensional conduction path along the stacking axis because there are no obvious intercolumnar short contacts (Figure 4d). In **EtPh-DBTII**, molecules in the adjacent columns aligned along the *c* axis are alternately tilted by about 65.8° in opposite directions (Figure 4g). Along this axis, intermolecular hydrogen bonds are formed between C=O of the molecular core and the hydrogen of the condensed benzene ring (Figure 4h). The overlap modes of **Bn-DBTII** and **EtPh-DBTII** are approximately the same (Figure 4b,f). Along the short axis of the **DBTII** plane, **Bn-DBTII** is slightly slipped while **EtPh-DBTII** is not slipped (Figures S9 and S10). Consequently, as reflected in the transfer integrals, these molecules show the same tendency that overlaps in the HOMO are much larger than overlaps in the LUMO along the stacking direction (Table 6).

Table 6. Transfer Integrals (meV) of **DBTII**s

R		t_a	t_b	t_c	t_r
Bn	HOMO	57.6	−4.3		−1.2
	LUMO	−33.1	−6.3		1.0
EtPh	HOMO		78.7	−9.0	
	LUMO		−42.8	−3.1	

In summary, only **EtPh-BTI** constructs a two-dimensional conduction path with a herringbone structure, while the other molecules form uniform stacking structures (Table 7). Among

Table 7. Summary of Torsion Angle and Short Contacts with π – π Distances

compd	R	between cores and phenyl rings (deg)	S...O interaction (Å)	hydrogen bond distance C=O...H (Å)	interplanar π – π distance (Å)
BTI	Bn	82.3	3.263(4) (inter)		3.40
	EtPh	4.6		2.66	
TIIG	Bn	72.0	3.105(1) (inter), 2.859(1) (intra)		3.37
	EtPh	21.3	2.859(2) (intra)	2.41	3.42
DBTII	Bn	79.0	2.810(1) (intra)		3.41
	EtPh	2.7	2.806(1) (intra)	2.68	3.41

them, all molecules are parallel in **Bn-DBTII**, whereas the adjacent columns are alternately tilted in **Bn-TIIG** and **EtPh-DBTII**. In **Bn-BTI** and **EtPh-TIIG**, the adjacent columns are parallel but the adjacent layers are alternately tilted in the opposite directions. The phenyl rings of the benzyl and 2-phenylethyl groups are not oriented parallel to the molecular core, but the tilt angles of the phenyl rings in the 2-phenylethyl groups are considerably smaller than those in the benzyl groups because the ethylene spacer in the 2-phenylethyl group enables a more appropriate orientation of the phenyl ring (Table 7). In addition, hydrogen bonds are observed only in the 2-phenylethyl-substituted molecules. Therefore, the ethyl-

ene spacer is more effective in forming an undistorted molecular structure.

Only **Bn-BTI** forms an S...S interaction, and **Bn-BTI** and **Bn-TIIG** have intermolecular S...O interactions. In **TIIG**s and **DBTII**s, intramolecular S...O interactions which guarantee the molecular core planarity are observed (Figures S7a–S10a).

Thin-Film Properties. By measurement of thin-film XRD and GIWAXS, molecular orientations in the thin films are investigated (Figure 5). **Bn-BTI** and **EtPh**-substituted molecules show sharp XRD peaks, but weak XRD peaks are observed in **Bn-TIIG** and **Bn-DBTII**. The extracted *d* values agree well with the out-of-plane peaks of GIWAXS (Table 8). In some cases, in-plane peaks are observed as well. Many sharp peaks are observed in **EtPh-TIIG** and **EtPh-DBTII**, indicating good crystallinity of these thin films.

In **Bn-BTI**, the *d* spacing is estimated to be 15.3 Å from XRD and GIWAXS ($q_z = 0.41 \text{ Å}^{-1}$). This value corresponds to half of the *b* axis. In addition, in the in-plane GIWAXS, $q_{xy} = 0.82 \text{ Å}^{-1}$ (=7.65 Å) and $q_{xy} = 1.30 \text{ Å}^{-1}$ (=4.82 Å) are observed, which are very close to the values for $c \sin \alpha \sin \beta$ (=7.54 Å) and $a \sin \beta \sin \gamma$ (=4.84 Å), respectively. Therefore, the crystallographic *ac* plane is aligned parallel to the substrate, where the tilt angle of the **BTI** plane from the substrate normal is 43° (Figure S12).

In **EtPh-BTI**, the *d* spacing is estimated to be 18.9 Å from XRD and GIWAXS ($q_z = 0.33 \text{ Å}^{-1}$). Although this value is longer than any axes of the unit cell, the length of the molecular long axis including **EtPh** is 18.6 Å and is close to the *d* spacing. In addition, in the in-plane GIWAXS, the same peak is observed at $q_{xy} = 0.33 \text{ Å}^{-1}$ (=18.9 Å), indicating that the molecular long axis is parallel to the substrate. Therefore, with retainment of the herringbone structure, the face-on and side-on orientations are mixed on the substrate (Figure S13). In the side-on orientation, the tilt angle of the **BTI** plane from the substrate normal is nearly 90° (84.8°). This orientation is superposed with the face-on orientation, in which the **BTI** plane is parallel to the substrate.

The XRD of **Bn-TIIG** shows two kinds of *d* spacings, 11.8 and 7.7 Å, which correspond to $q_z = 0.53 \text{ Å}^{-1}$ (=11.8 Å) and $q_z = 0.82 \text{ Å}^{-1}$ (=7.65 Å) in the out-of-plane GIWAXS. The former corresponds to the value for $a \sin \beta \sin \gamma$ (=11.87 Å), and the latter is close to half of the value for $c \sin \alpha \sin \beta$ (=7.3 Å). Therefore, two types of orientations coexist in the thin films, in which the crystallographic *bc* and *ab* planes are aligned parallel to the substrate. In the former case, the tilt angle of the **TIIG** plane is 22° from the substrate normal; this is basically the side-on orientation (Figure S14). In the latter case, the tilt angle of the **TIIG** plane is 32° from the substrate normal; this is the edge-on orientation (Figure S14). In the in-plane GIWAXS, peaks at $q_{xy} = 0.54 \text{ Å}^{-1}$ (=11.7 Å) and $q_{xy} = 0.82 \text{ Å}^{-1}$ (=7.65 Å) are observed. This observation also supports the coexistence of two kinds of orientations.

The XRD and out-of-plane GIWAXS in **EtPh-TIIG** show the same *d* spacing of 13.8 Å ($q_z = 0.456 \text{ Å}^{-1}$), but this value is larger than half of the *c* axis (11.1 Å). This implies that the molecules in the thin films are less tilted than the molecules in the single crystals. The thin-film tilt angle from the vertical direction to the substrate (β') is calculated from $l = d/\cos \beta'$, where *l* is the molecule length estimated from the crystal structure (19.9 Å) and *d* is the *d* spacing (13.8 Å). This leads to $\beta' = 46^\circ$, which is slightly smaller than the single-crystal value of 48°. In the thin films, the molecules are oriented more

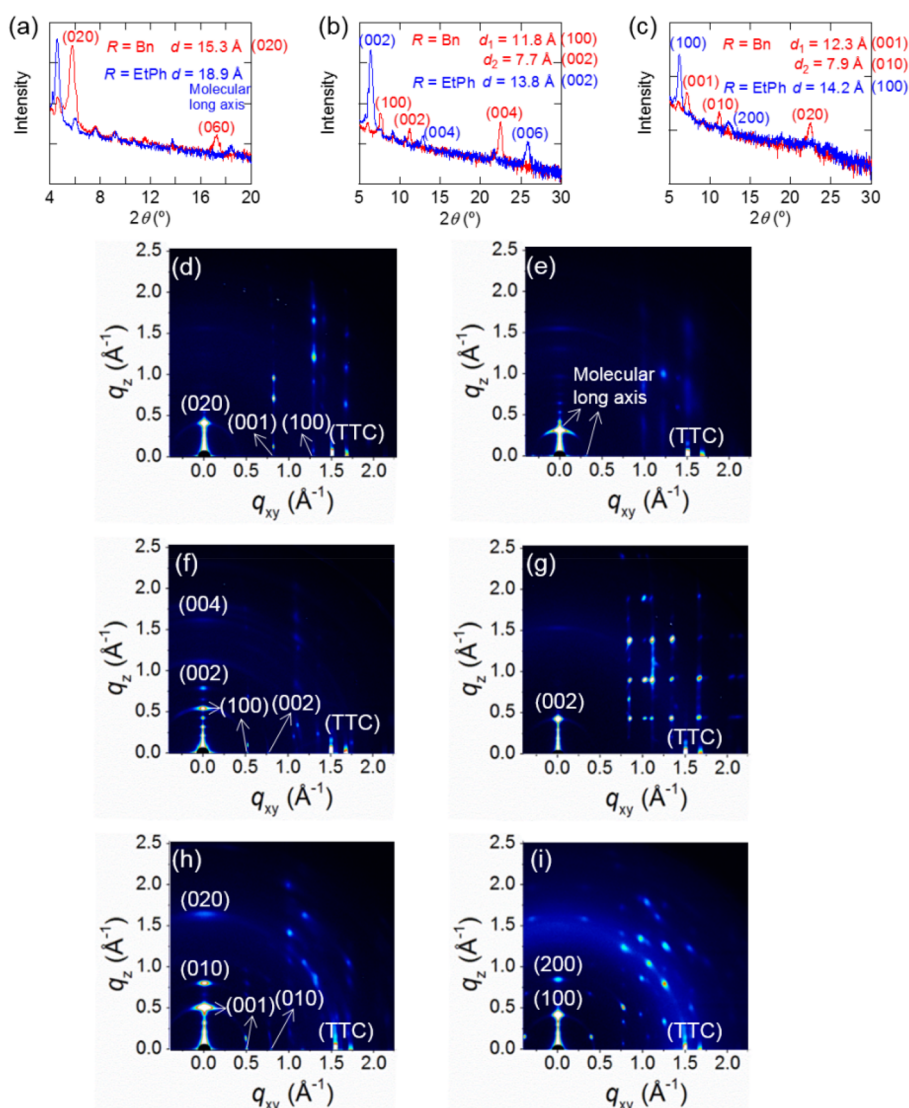


Figure 5. XRD patterns of (a) BTIs, (b) TIIGs, and (c) DBTIIs. GIWAXS patterns of (d) Bn-BTI, (e) EtPh-BTI, (f) Bn-TIIG, (g) EtPh-TIIG, (h) Bn-DBTII, and (i) EtPh-DBTII.

Table 8. XRD and GIWAXS d Values (Å)

compd	R	XRD (hkl)	GIWAXS (hkl)	
			out of plane	in plane
BTI	Bn	15.3 (020), 5.1 (060)	15.3 (020)	7.65 (001), 4.82 (100)
	EtPh	18.9 (molecular long axis)	18.9 (molecular long axis)	18.9 (molecular long axis)
TIIG	Bn	11.8 (100), 7.7 (002), 3.85 (004)	11.8 (100), 7.65 (002), 3.85 (004)	11.7 (100), 7.65 (002)
	EtPh	13.8 (002), 6.9 (004), 3.45 (006)	13.8 (002)	
DBTII	Bn	12.3 (001), 7.9 (010), 3.95 (020)	12.3 (001), 7.91 (010), 3.95 (020)	12.3 (001), 7.91 (010)
	EtPh	14.2 (100), 7.14 (200)	14.2 (100), 7.13 (200)	

closely to the perpendicular direction in comparison to the single crystals (Figure S15).

The XRD of Bn-DBTII shows two d spacings, 12.3 and 7.9 Å, which correspond to $q_z = 0.51 \text{ Å}^{-1}$ ($=12.3 \text{ Å}$) and $q_z = 0.79 \text{ Å}^{-1}$ ($=7.91 \text{ Å}$) in the out-of-plane GIWAXS, respectively. In the single crystal, $c \sin \alpha \sin \beta$ is 12.3 Å and $b \sin \alpha \sin \gamma$ is 7.9 Å; thus, these values are in good agreement with the values of extracted d spacings. In addition, in-plane GIWAXS peaks are observed at $q_{xy} = 0.51 \text{ Å}^{-1}$ ($=12.3 \text{ Å}$) and $q_{xy} = 0.79 \text{ Å}^{-1}$ ($=7.91 \text{ Å}$). Similarly to Bn-TIIG, two types of orientations coexist in the thin films; the crystallographic ab and ac planes are aligned

parallel to the substrate. The tilt angle of the DBTII plane is 55° from the substrate normal in the edge-on orientation and 26° in the side-on orientation (Figure S16).

From the out-of-plane GIWAXS and XRD of EtPh-DBTII, the d spacing is estimated to be 14.2 Å ($q_z = 0.44 \text{ Å}^{-1}$). This value is close to that for $a \sin \beta \sin \gamma$ ($=14.14 \text{ Å}$), indicating that the bc plane is aligned parallel to the substrate. Therefore, the molecules are arranged in an edge-on orientation, and the tilt angle of the DBTII plane is 47° from the substrate normal (Figure S17).

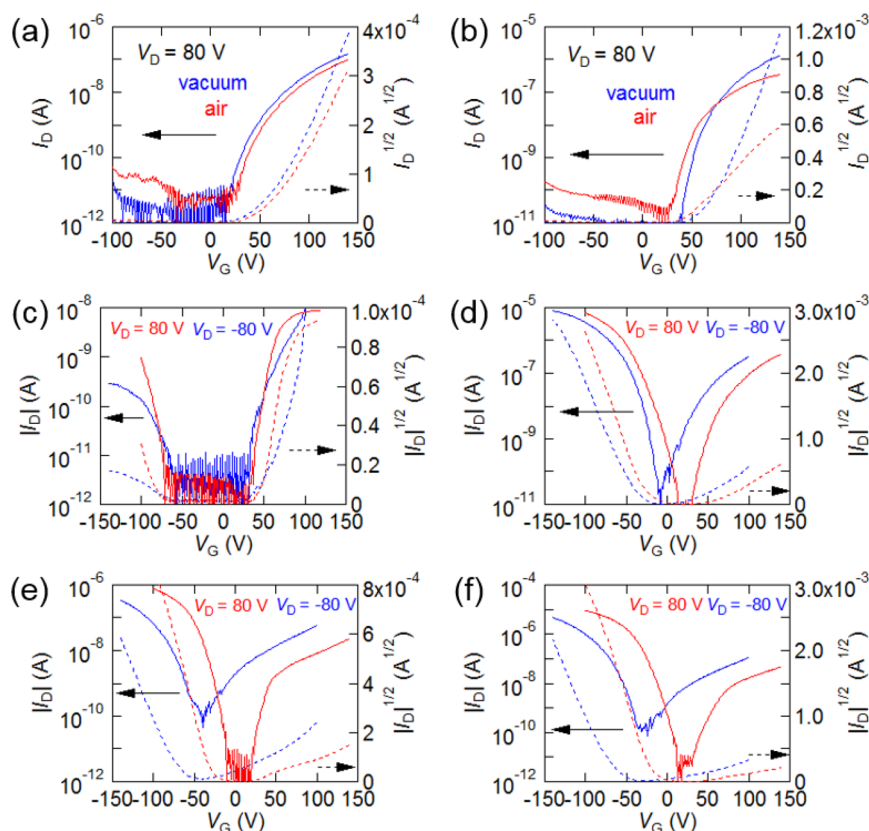


Figure 6. Transfer characteristics of (a) Bn-BTI, (b) EtPh-BTI, (c) Bn-TIIG, (d) EtPh-TIIG, (e) Bn-DBTII, and (f) EtPh-DBTII.

Table 9. OFET Properties of BTIs, TIIGs, and DBTIIs

compd	measurement	hole			electron		
		μ_{av} [μ_{max}] ($\text{cm}^2 \text{V}^{-1} \text{s}^{-1}$)	V_T (V)	I_{on}/I_{off}	μ_{av} [μ_{max}] ($\text{cm}^2 \text{V}^{-1} \text{s}^{-1}$)	V_T (V)	I_{on}/I_{off}
Bn-BTI	under vacuum				5.4×10^{-4} [7.2×10^{-4}]	60	10^5
	in air				2.9×10^{-4} [3.7×10^{-4}]	63	10^4
EtPh-BTI	under vacuum				6.4×10^{-3} [0.01]	67	10^5
	in air				9.6×10^{-4} [1.8×10^{-3}]	50	10^4
Bn-TIIG	under vacuum	6.3×10^{-7} [1.2×10^{-6}]	−41	10^2	9.4×10^{-5} [2.3×10^{-4}]	44	10^4
EtPh-TIIG	under vacuum	0.017 [0.039]	−38	10^5	1.3×10^{-3} [3.5×10^{-3}]	64	10^4
Bn-DBTII	under vacuum	2.4×10^{-3} [4.7×10^{-3}]	−70	10^3	5.3×10^{-5} [8.3×10^{-5}]	35	10^4
EtPh-DBTII	under vacuum	0.027 [0.051]	−68	10^5	9.1×10^{-5} [1.7×10^{-4}]	36	10^3

Transistor Properties. BTIs show n-type transistor characteristics under vacuum and in air because the HOMO levels are too low to show hole transport (Table 1 and Figure 6). EtPh-BTI exhibits improved mobility of more than 1 order of magnitude in comparison to Bn-BTI. This is attributable to the tilt angles from the substrate normal in the thin films. In general, when the molecules are arranged more closely to the substrate normal, the mobility is improved.⁴⁰ Even though face-on-arranged molecules coexist in EtPh-BTI, the side-on-arranged molecules are oriented more perpendicularly to the substrate in comparison to Bn-BTI. However, the mobilities of these two molecules decrease in comparison to the previously reported alkyl-BTIs.²³ The reason is roughness and many grain boundaries in the thin films, as confirmed by AFM (Figure S11a,b). In EtPh-BTI, a two-dimensional conduction path is constructed in the herringbone structure, but one of the conduction paths faces the direction vertical to the substrate in the thin films (Figure S13). Consequently, a one-dimensional conduction path is formed between the electrodes, leading to

the reduced mobility. Upon air exposure, BTIs show slightly decreased OFET properties and the off current increases. These phenomena are consistent with the relatively high LUMO levels in comparison with the air-stable n-type molecules, which have LUMO levels below −4.0 eV.³⁹ However, owing to the passivation effect of the bulky substituents,^{41–44} the air stability is improved in comparison to the previously reported alkyl-BTIs.²³ OFET properties of BTIs, TIIGs, and DBTIIs are given in Table 9.

TIIGs and DBTIIs exhibit ambipolar transistor characteristics under vacuum (Figure 6). Bn-TIIG shows n-dominant ambipolar properties, which is consistent with the transfer integrals (Table 5). Other molecules exhibit p-dominant ambipolar properties, which are attributable to the large transfer integrals between the HOMO and the better donor abilities of those molecules as shown in the relatively high HOMO levels (−5.2 to −5.4 eV, Table 1), which are closer to the work function of the gold electrodes.⁴⁵ Among the four molecules, the Bn-substituted molecules exhibit lower OFET

properties. The Bn-substituted molecules have two kinds of orientations in the thin films, and the randomness is the origin of the ineffective carrier transport. In comparison, improved OFET properties are observed in the EtPh-substituted molecules. In particular, EtPh-TIIG exhibits a maximum hole mobility of about $0.04 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, which is higher than that of N-unsubstituted TIIG.²⁴ This comes from the considerably large transfer integral between the HOMOs of the stacking molecules. The electron mobility is comparable to that of N-unsubstituted TIIG. EtPh-DBTII shows much higher hole mobility due to the extended π skeleton with a reduced reorganization energy for holes (Table S1). However, in comparison to hexyl-DBTII,²⁵ hole and electron mobilities are reduced because of the large tilt angle of the molecular cores from the substrate normal (47°) in the thin films (Figure S17).

CONCLUSIONS

In this paper, we have investigated benzyl (Bn)- and 2-phenylethyl (EtPh)-substituted bithienoisatin (BTI), thienoisindigo (TIIG), and dibenzothienoisindigo (DBTII) molecules. Since the energy levels depend only on the molecular core, these molecules show HOMO and LUMO levels similar to those of the alkyl-substituted BTI, TIIG, and DBTII. The HOMO levels decrease in the order DBTIIs < TIIGs < BTIs, and the LUMO levels decrease in the order TIIGs < DBTIIs < BTIs. Therefore, BTIs show n-type OFET properties and TIIGs and DBTIIs exhibit ambipolar OFET properties, which are in good agreement with the expectations from the energy levels. These molecules form uniform stacking structures, though EtPh-BTI has a slipped herringbone structure, and a few intermolecular interactions are observed owing to the bulky substituents. The benzyl parts have large torsion angles from the molecular core, while the 2-phenylethyl parts have small torsion angles from the molecular core. This is because the latter has an ethylene spacer that is adequate to enable suitable molecular packing. Therefore, the 2-phenylethyl substitution is more advantageous to the carrier transport, and this is evident from the OFET mobilities, which are about $0.01 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ for electrons in EtPh-BTI and 0.04 and $0.05 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ for holes in EtPh-TIIG and EtPh-DBTII, respectively. In contrast, from the thin film XRD and GIWAXS, it is indicated that two kinds of molecular arrangements coexist in the thin films of the benzyl-substituted TIIG and DBTII. This is unfavorable for the carrier transport, which is confirmed by the OFET mobilities of Bn-TIIG and Bn-DBTII. In TIIGs, the dominant carrier polarity in OFET is opposite; Bn-TIIG shows n-dominant properties but EtPh-TIIG shows p-dominant properties. This is caused by the difference in slip distance along the molecular short axis, as reflected in the transfer integrals.

The present work demonstrates how benzyl and 2-phenylethyl substitutions affect the crystal structures and OFET properties in BTI, TIIG, and DBTII. It is confirmed that one-dimensional stacking structures are maintained in almost all cases, even though the bulky substituents are attached at the nitrogen face to the short-axis side of the molecular core. These results give us a hint for designing side substituents of small-molecule and polymer OSCs.

ASSOCIATED CONTENT

Supporting Information

Crystal data (CIF) The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.cgd.0c00087>.

Additional information on material synthesis, electrochemical properties, and optical properties, DFT calculations, thermal properties, crystal structures, thin-film microstructures and morphologies, and thin-film transistor properties (PDF)

Accession Codes

CCDC 1974945–1974950 contain the supplementary crystallographic data for this paper. These data can be obtained free of charge via www.ccdc.cam.ac.uk/data_request/cif, or by emailing data_request@ccdc.cam.ac.uk, or by contacting The Cambridge Crystallographic Data Centre, 12 Union Road, Cambridge CB2 1EZ, UK; fax: +44 1223 336033.

AUTHOR INFORMATION

Corresponding Authors

Dongho Yoo – Department of Materials Science and Engineering, Tokyo Institute of Technology, Tokyo 152-8552, Japan; orcid.org/0000-0003-0886-7533; Email: dongho312@postech.ac.kr

Minoru Ashizawa – Department of Materials Science and Engineering, Tokyo Institute of Technology, Tokyo 152-8552, Japan; orcid.org/0000-0002-6810-256X; Email: ashizawa.m.aa@m.titech.ac.jp

Takehiko Mori – Department of Materials Science and Engineering, Tokyo Institute of Technology, Tokyo 152-8552, Japan; orcid.org/0000-0002-0578-5885; Email: mori.tae@m.titech.ac.jp

Authors

Akihiro Kohara – Department of Materials Science and Engineering, Tokyo Institute of Technology, Tokyo 152-8552, Japan

Tadashi Kawamoto – Department of Materials Science and Engineering, Tokyo Institute of Technology, Tokyo 152-8552, Japan; orcid.org/0000-0002-5676-4013

Hiroyasu Masunaga – Japan Synchrotron Radiation Research Institute (JASRI/SPring-8), Hyogo 679-5198, Japan

Noboru Ohta – Japan Synchrotron Radiation Research Institute (JASRI/SPring-8), Hyogo 679-5198, Japan

Hidetoshi Matsumoto – Department of Materials Science and Engineering, Tokyo Institute of Technology, Tokyo 152-8552, Japan; orcid.org/0000-0002-4949-1184

Complete contact information is available at: <https://pubs.acs.org/doi/10.1021/acs.cgd.0c00087>

Notes

The authors declare no competing financial interest.

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