

# Probing Interligand Electron Transfer in the $^1\text{MLCT } S_1$ Excited State of $\text{trans-Mo}_2\text{L}_2\text{L}'_2$ Compounds: A Comparative Study of Auxiliary Ligands and Solvents

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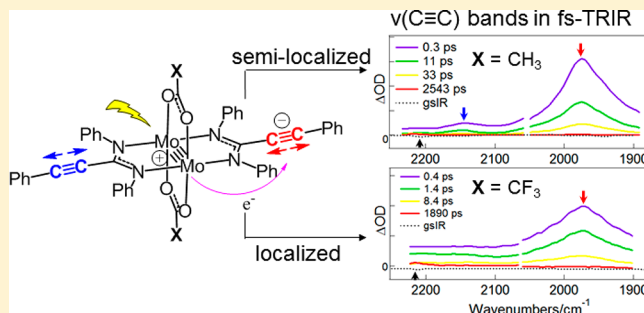
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## Supporting Information

**ABSTRACT:** The interligand charge dynamics of the lowest singlet metal-to-ligand charge-transfer states ( $^1\text{MLCT } S_1$  states) of a series of quadruply bonded  $\text{trans-Mo}_2(\text{NN})_2(\text{O}_2\text{C-X})_2$  paddlewheel compounds are investigated, where NN is a  $\pi$ -accepting phenylpropiolamidinate ligand and  $\text{O}_2\text{C-X}$  ( $\text{X} = \text{Me}$ ,  $\text{tBu}$ ,  $\text{T}^i\text{PB}$ , or  $\text{CF}_3$ ) is an auxiliary carboxylate ligand. The compounds show strong light absorption in the visible region due to MLCT transitions from the  $\text{Mo}_2$  center to the NN ligands. The transferred electron density was followed by femtosecond time-resolved infrared (fs-TRIR) spectroscopy with vibrational reporters such as the ethynyl groups on the NN ligands. The observed fs-TRIR spectra show that these compounds have asymmetric  $^1\text{MLCT } S_1$  excited states where the transferred electron mainly resides on a single NN ligand. The presence of interligand electron transfer (ILET) is suggested to explain the shape of the  $\nu(\text{C}\equiv\text{C})$  bands and the influence of auxiliary ligands and solvents on the interligand electronic coupling. The ILET in the  $^1\text{MLCT } S_1$  state is shown to be sensitive to the functional groups on the auxiliary ligands while being less responsive to changes in solvents.



## INTRODUCTION

The photophysics of transition metal (TM) complexes present an opportunity for the application of fundamental research in inorganic chemistry.<sup>1–3</sup> TM complexes with high-energy and long-living charge-transfer excited states have found wide applications in the area of solar energy conversion and photocatalysis.<sup>4,5</sup> However, some fundamental questions on the nature of the excited-state charge dynamics of TM complexes remain, such as a clear description of the charge distribution in compounds with equivalent ligands or metal centers.<sup>1,6,7</sup> The location of the positive or negative charge in such charge-transfer excited states can be difficult to assign. A simple description using integer point charges can be insufficient or inaccurate when considering electron transfer and other interactions between the identical redox centers. In previous works, the charge distribution in charge-transfer excited states of TM compounds has been studied by spectroscopic methods such as luminescence, Stark spectroscopy, and transient absorption spectroscopy as well as by computational studies.<sup>8–10</sup> In recent years, femtosecond time-resolved infrared

(fs-TRIR) spectroscopy has become a more reliable and direct method to study excited-state photophysics.

In the Chisholm group, we have focused on the lowest metal-to-ligand charge-transfer singlet states ( $^1\text{MLCT } S_1$  states) of  $\text{trans-M}_2\text{L}_2\text{L}'_2$  and  $\text{M}_2\text{L}_4$  paddlewheel compounds ( $\text{M} = \text{Mo}$  or  $\text{W}$ ) containing metal–metal quadruple bonds and  $\pi$ -accepting ligands. The MLCT states of these  $d^4$ – $d^4$  metal dimers were shown to display a variety of charge-distribution patterns with respect to the transferred electron.<sup>11–14</sup> The distribution of the electron ranges from being totally localized, semilocalized, and totally delocalized over the acceptor ligands. For example, the transferred electron was found to be localized on one  $(\text{NPh})_2\text{CC}\equiv\text{CPh}$  ligand in the  $^1\text{MLCT } S_1$  state of  $\text{Mo}_2[(\text{NPh})_2\text{CC}\equiv\text{CPh}]_4$ ,<sup>15</sup> whereas it was found to be semilocalized over two  $(\text{N}^i\text{Pr})_2\text{CC}\equiv\text{CPh}$  ligands in the  $^1\text{MLCT } S_1$  state of  $\text{trans-Mo}_2[(\text{N}^i\text{Pr})_2\text{CC}\equiv\text{CPh}]_2(\text{O}_2\text{CMe})_2$  and fully delocalized over two  $(\text{NPh})\text{C}(\text{S})\text{C}\equiv\text{CPh}$  ligands in

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the  $^1\text{MLCT } S_1$  state of  $\text{trans-Mo}_2[(\text{NPh})\text{C}(\text{S})\text{C}\equiv\text{CPh}]_2(\text{O}_2\text{C-T}^i\text{PB})_2$ .<sup>11,16</sup>

We have previously suggested that the observed charge-distribution patterns in the  $^1\text{MLCT } S_1$  states of  $\text{trans-M}_2\text{L}_2\text{L}'_2$  and  $\text{M}_2\text{L}_4$  compounds are the results of an interligand electron transfer (ILET) process on the vibrational time scale ( $\sim 10^{-12}$  s), where “fast” ILET is associated with MLCT states delocalized over multiple ligands, while “slow” ILET is associated with MLCT states localized on one ligand.<sup>16</sup> Due to the mediator role of the dimetal center, the electron transfer can be greatly influenced by the electronic properties of the bimetallic core and its interactions with the acceptor ligands.<sup>12,16</sup>

Here, a series of  $\text{trans-Mo}_2\text{L}_2\text{L}'_2$  amidinate compounds (shown in Figure 1) were prepared with different carboxylate

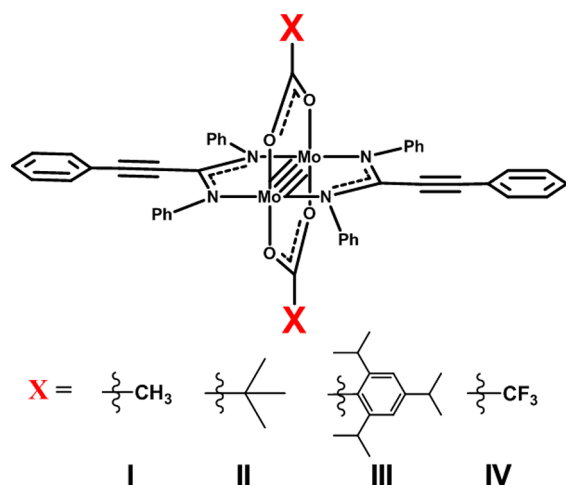


Figure 1. General structure of I–IV.

auxiliary ligands, namely,  $\text{trans-Mo}_2(\text{NN})_2(\text{O}_2\text{C-Me})_2$  (I),  $\text{trans-Mo}_2(\text{NN})_2(\text{O}_2\text{C-}^t\text{Bu})_2$  (II),  $\text{trans-Mo}_2(\text{NN})_2(\text{O}_2\text{C-T}^i\text{PB})_2$  (III), and  $\text{trans-Mo}_2(\text{NN})_2(\text{O}_2\text{C-CF}_3)_2$  (IV), where NN = *N,N'*-diphenylphenylpropiolamidinate, and Me = methyl,

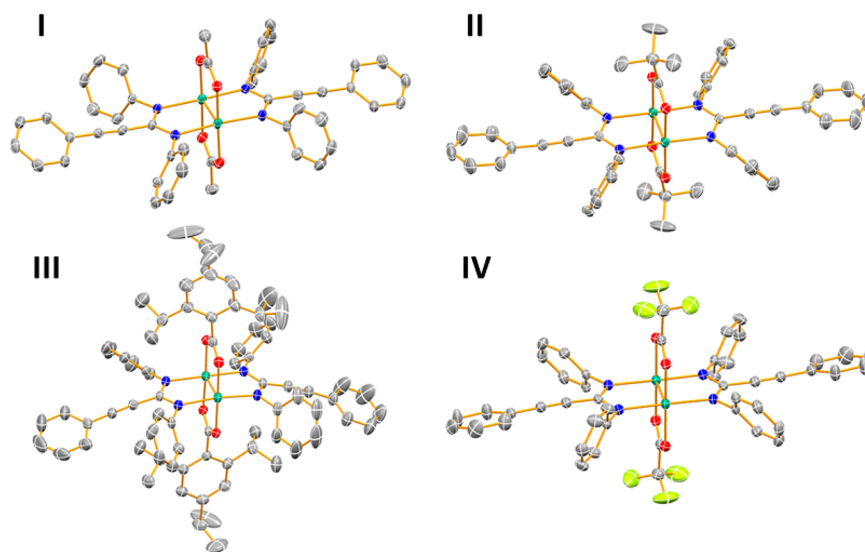


Figure 2. Crystal structures of I–IV with thermal ellipsoids drawn at the 50% probability level. Hydrogens and solvent molecules were omitted for clarity. Green = molybdenum, blue = nitrogen, red = oxygen, gray = carbon, and yellow = fluorine. (The plot of III was adapted from ref 15 with permission from American Chemical Society.)

$^t\text{Bu}$  = *tert*-butyl,  $\text{T}^i\text{PB}$  = 2,4,6-triisopropylbenzyl, and  $\text{CF}_3$  = trifluoromethyl. Preparation and characterization of compound III were reported previously.<sup>15</sup> These compounds show strong MLCT transitions to the NN ligands but not the carboxylate ligands.

The excited-state photophysics of I–IV and the interligand charge dynamics between NN ligands were examined with respect to perturbations exerted by the auxiliary carboxylate ligands. Considering previous reports of charge localization in these compounds (the dipole moment of the excited states is a function of electron localization), the influence of solvent dipoles on the charge distribution was also investigated.<sup>11,15</sup>

## RESULTS AND DISCUSSION

**Synthesis and Crystal Structure.** Compounds I–IV were prepared using a common synthetic method where two equivalents of the NN ligand (lithium *N,N'*-diphenylphenylpropiolamidinate) were reacted with one equivalent of the appropriate  $\text{Mo}_2\text{L}_4$  carboxylate compound.<sup>17–19</sup> The compounds were purified through solution processes, and the purity of each was confirmed by NMR and MALDI-TOF mass spectrometry. The structures of I–IV obtained by single-crystal X-ray crystallography are displayed in Figure 2.

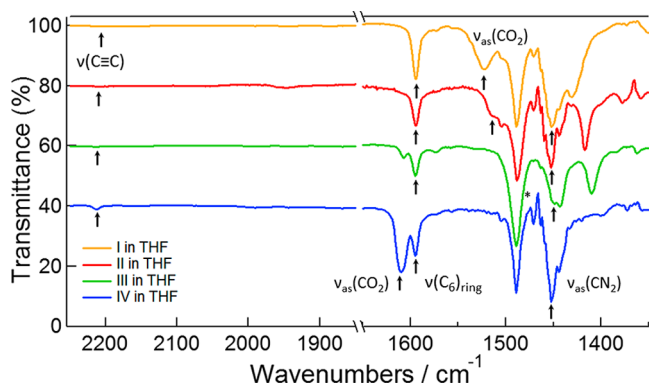
The core structures of I–IV have effective  $D_{2h}$  symmetry. All of the molecules adopt the *trans* configuration, in which the plane of the carboxylate ligands is perpendicular to the plane of the NN ligands. Those on the NN ligands in I–IV are found to be coplanar with the  $\text{CN}_2\text{-Mo}_2\text{-N}_2\text{C}$  plane. The phenyl-acetylenyl groups in III are slightly tilted due to the arrangement of the phenyl groups on the nitrogen atoms. All of the phenyl groups on the donor nitrogen atoms are effectively twisted ( $\sim 50\text{--}80^\circ$ ) away from the  $\text{CN}_2$  plane.

All four compounds display Mo–Mo bond lengths of around 2.10 Å, typical of Mo–Mo quadruple bonds,<sup>20</sup> and average Mo–N and Mo–O bond lengths of around 2.13 and 2.12 Å, respectively. The Mo–Mo and Mo–O bond lengths in IV are slightly longer (by  $\sim 0.02$  Å) than those in I–III, likely due to the influence of the strongly electron-withdrawing  $\text{-CF}_3$

groups.<sup>21</sup> The intramolecular distance between the C≡C bond centers is around 9.7 Å in I–IV.

In the solid-state, each Mo center in II and IV was found to be coordinated by a dimethyl sulfoxide (DMSO) solvent molecule, while only one metal in I and III was so coordinated. The crystal structures of I–IV including coordinating solvent molecules are displayed in Figures S9 and S10 in the Supporting Information.

**Ground-State Infrared Spectroscopy.** The C≡C bonds, phenyl rings, and CN<sub>2</sub> moieties on the NN ligands and CO<sub>2</sub> moieties on the carboxylate ligands are convenient IR reporter groups for the excited-state charge distribution and interligand charge dynamics. I–IV were studied in solution by FT-IR spectra to identify their signature vibrational bands. The IR spectra of I–IV in tetrahydrofuran (THF) are shown in Figure 3, and the IR spectra of III in toluene, dichloromethane



**Figure 3.** FT-IR spectra of I–IV in THF solution. The spectra lines are plotted with an offset from each other on the Y axis. \*The  $\nu_{as}(\text{CO}_2)$  of III was masked.

(DCM), and dimethylformamide (DMF) are shown in Figure S11 in the Supporting Information. Frequency calculations were used to assist the assignments of above-mentioned vibrational bands. The calculated IR spectra of I–IV are shown in Figure S12 in the Supporting Information.

The vibrational bands belonging to the NN ligands appear at similar frequencies in the ground states of I–IV. A weak band at around 2210  $\text{cm}^{-1}$  was found for the  $\nu(\text{C}\equiv\text{C})$  stretch; two sharp bands, one at around 1595  $\text{cm}^{-1}$  and one at around 1490  $\text{cm}^{-1}$ , were found for  $\nu(\text{C}_6)_{ring}$  stretches of two different modes, and a strong band at around 1450  $\text{cm}^{-1}$  was found for the asymmetric  $\nu(\text{CN}_2)$  stretch. The irregular shape of the  $\nu(\text{CN}_2)$  bands is due to strong solvent absorption in the region. The weaker, neighboring bands at 1605 and 1480  $\text{cm}^{-1}$  in III were attributed to the  $\nu(\text{C}_6)_{ring}$  stretching bands in the  $-\text{T}^{\text{PB}}$  moiety.

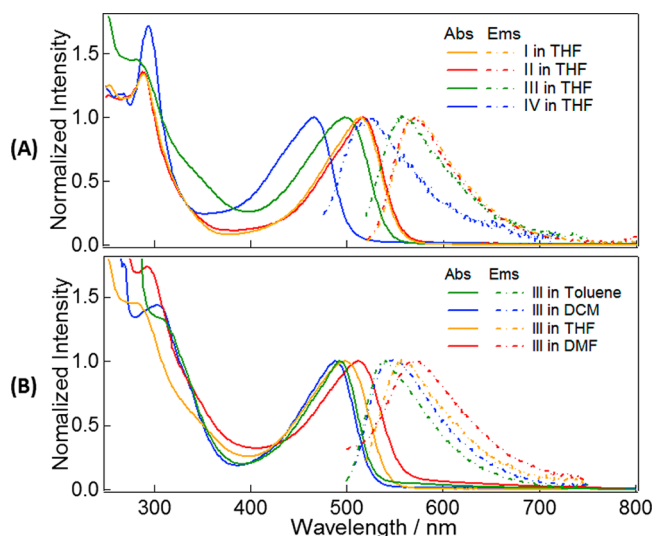
In comparison, the asymmetric  $\nu(\text{CO}_2)$  stretching bands of I–IV appear over a wider range of frequencies due to different functional groups. A strong  $\nu(\text{CO}_2)$  band was found at 1523  $\text{cm}^{-1}$  for I, at 1515  $\text{cm}^{-1}$  for II, and at 1623  $\text{cm}^{-1}$  for IV. The  $\nu(\text{CO}_2)$  band for III likely occurs near 1490  $\text{cm}^{-1}$  (as spotted in the ground-state bleach in the fs-TRIR spectra of III), where it is masked by the strong  $\nu(\text{C}_6)_{ring}$  band.<sup>15</sup> The high energy of the  $\nu(\text{CO}_2)$  vibrations in IV, again, is a result of the electron-withdrawing nature of the  $-\text{CF}_3$  groups which depopulated the electron density in the  $\pi^*$  orbital of the carboxylate ligand.

Compound III in solvents of differing polarities was also studied by FT-IR spectroscopy, as shown in Figure S11. The

frequencies of the  $\nu(\text{C}\equiv\text{C})$  and  $\nu(\text{C}_6)_{ring}$  bands did not shift significantly with respect to change of solvents. The  $\nu(\text{CN}_2)$  bands, alternatively, shifted slightly to a higher energy from the toluene to DCM to THF solutions.

#### Electronic Absorption and Emission Spectroscopy.

Compounds I–IV display intense absorption and weak fluorescence emission in solution at room temperature. The UV–vis absorption and emission spectra of I–IV collected in THF are shown in Figure 4A, and the spectra of III in toluene,



**Figure 4.** Normalized absorption (Abs) and emission (Ems) spectra of I–IV, (A) I–IV in THF solution and (B) III in toluene, DCM, THF, and DMF solutions.

DCM, THF, and DMF are shown in Figure 4B. Quantitative parameters like the MLCT maxima, molar extinction coefficients ( $\epsilon$ ), and fluorescence quantum yields ( $\Phi_{\text{FL}}$ ) are summarized in Table 1.

As shown by Figure 4 and Table 1, the MLCT transitions of I–IV were significantly influenced by both the carboxylate ligands and the solvents. In THF, II featuring the *tert*-butylcarboxylate ligands shows the lowest-energy MLCT band at 517 nm, and IV featuring the trifluoroacetate ligands shows the highest-energy MLCT band at 486 nm, with an energy difference of  $\sim 1240 \text{ cm}^{-1}$ . The molar extinction coefficient of II is almost double that of IV.

With respect to solvent environments, III shows the lowest-energy MLCT absorption in DMF and the highest-energy MLCT absorption in DCM (not toluene). DMF significantly broadens the MLCT band and shifts it by  $\sim 980 \text{ cm}^{-1}$ . The observed solvatochromism here is in accordance with strong dipole–dipole interactions, as well as factors like solvent coordination (Figure S9 and S10) and also  $\pi$ – $\pi$  stacking interactions.

The emission spectra of I–IV in the visible range are all identified as fluorescence from the <sup>1</sup>MLCT S<sub>1</sub> states. The quantum yields are generally well-below 1% due to charge-transfer transitions and strong competition from intersystem crossing. The fluorescence bands of I–IV resemble the MLCT absorption bands with respect to both their shapes and trends in energy. These compounds also show strong phosphorescence from the lowest metal-centered triplet states (<sup>3</sup> $\delta\delta^*$  T<sub>1</sub> states) in the near-infrared region ( $\sim 800$ – $1100 \text{ nm}$ ), but they are not discussed here as they are outside the scope of this

Table 1. Photophysical Data of I–IV for the MLCT Transitions

compounds	solvent	absorption $\lambda_{\max}/\text{nm}$	$\epsilon/\text{M}^{-1}\text{cm}^{-1}$	fluorescence $\lambda_{\max}/\text{nm}$	$\Phi_{\text{FL}}/\%$	Stokes shift/ $\text{cm}^{-1}$
I	THF	514	37 600	573	0.20	2000
II	THF	517	41 600	570	0.23	1800
III	toluene	494	29 900	541	0.13	1800
	DCM	489	33 000	545	0.09	2100
	THF	498	29 400	557	0.19	2360
IV	DMF	514	24 600	571	0.17	1940
	THF	466	23 900	525	0.03	2410

paper. The Stokes shifts, which measure the energy difference between the vertical absorption and emission transitions, range from 1800 and 2410  $\text{cm}^{-1}$  for I–IV. IV shows the largest Stokes shift and lowest fluorescence quantum yield.

#### Femtosecond Time-Resolved Infrared Spectroscopy.

The  $^1\text{MLCT}$   $S_1$  states of I–IV were further studied using femtosecond time-resolved infrared (fs-TRIR) spectroscopy. The spectra display vibrational bands that are present in the  $^1\text{MLCT}$   $S_1$  state and the  $^3\delta\delta^*$   $T_1$  state, as seen for similar compounds.<sup>11,22,15,16</sup> The shifts in the vibrations, whether to higher or lower energy from the ground state, provide information on the charge-density distribution and charge dynamics in the excited states.

The  $S_1$  lifetimes of I–IV found in the fs-TRIR spectra are between 10 and 21 ps, as shown in Table 2. The kinetic traces

Table 2.  $^1\text{MLCT}$   $S_1$  State Lifetimes of Compounds I–IV

compounds	solvents	$\tau(S_1)/\text{ps}$
I	THF	$19.0 \pm 0.4$
II	THF	$21.4 \pm 0.9$
III	toluene	$15.5 \pm 0.3$
	THF	$19.1 \pm 0.4$
IV	DMF	$15.2 \pm 0.9$
	THF	$10.5 \pm 2.3$

and exponential fitting plots of I–IV are displayed in S14 and 15. In all cases, a short lifetime component ( $<1$  ps) and a longer lifetime component were found for the  $^1\text{MLCT}$   $S_1$  state features, where the former was attributed to intramolecular relaxation processes.

**fs-TRIR Spectra of I–IV in THF Solution.** The fs-TRIR spectra of I–IV were collected in THF solution at room temperature with excitation of I–IV into their MLCT bands, I–III at 515 nm, and IV at 400 nm. The fs-TRIR spectra of I–IV in the range of 1350–2250  $\text{cm}^{-1}$  with selected time delays are shown in Figure 5. For each compound, the inverted ground-state IR spectrum is also plotted as a dotted line for reference.

The fs-TRIR spectra of I–IV at early time delays (purple and blue traces) are associated with the  $^1\text{MLCT}$   $S_1$  states. Strong absorption bands of the  $\nu(\text{C}\equiv\text{C})$  and  $\nu(\text{C}_6)_{\text{ring}}$  stretches were observed; these stretches show significant shifts to lower wavenumbers compared to the ground state, in accordance to the displacement of electron density from the  $\text{Mo}_2$  center to the  $\pi^*$  antibonding orbitals of the NN ligands.

In the  $^1\text{MLCT}$   $S_1$  states of I–III, each compound displays two transient  $\nu(\text{C}\equiv\text{C})$  features, a strong absorption band at  $\sim 1970$   $\text{cm}^{-1}$  and a weak absorption band at  $\sim 2150$   $\text{cm}^{-1}$ . The two  $\nu(\text{C}\equiv\text{C})$  stretches are shifted by  $-235$  and  $-65$   $\text{cm}^{-1}$  in I,  $-230$  and  $-64$   $\text{cm}^{-1}$  in II, and  $-230$  and  $-70$   $\text{cm}^{-1}$  in III. Each compound also shows one  $\nu(\text{C}_6)_{\text{ring}}$  absorption band at  $\sim 1560$

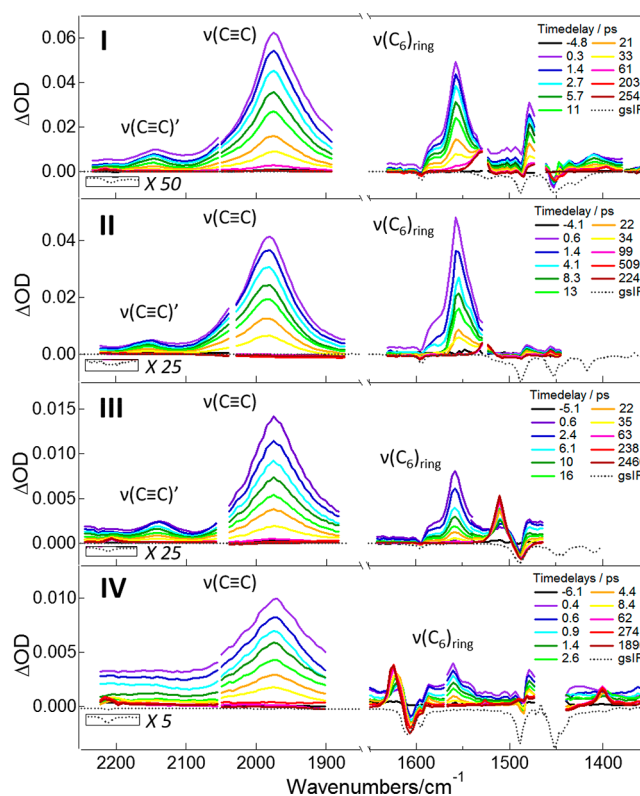


Figure 5. fs-TRIR spectra of I–IV in THF solution at room temperature with inset showing time delays. Inverted ground-state IR (gs-IR) spectra are shown by dotted lines.

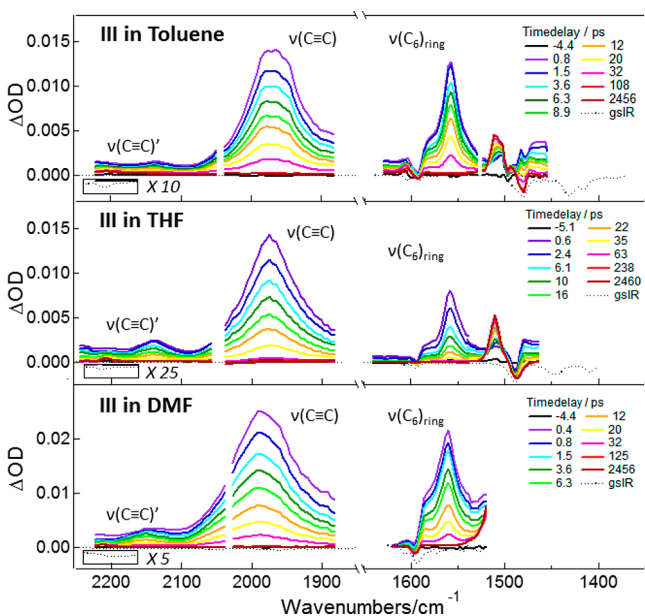
$\text{cm}^{-1}$  together with a ground-state bleach band at  $\sim 1595$   $\text{cm}^{-1}$ . The  $\nu(\text{C}_6)_{\text{ring}}$  stretches are shifted around  $-40$   $\text{cm}^{-1}$ . The observation of two  $\nu(\text{C}\equiv\text{C})$  bands in I–III, as suggested in previous works, corresponds to two unequally reduced NN ligands in the  $^1\text{MLCT}$   $S_1$  state.<sup>11,15,16</sup>

The fs-TRIR spectra of IV, however, are quite different. The most noticeable feature is that only one  $\nu(\text{C}\equiv\text{C})$  transient feature was observed at picosecond time delays. The strong  $\nu(\text{C}\equiv\text{C})$  band of IV is observed at  $1970$   $\text{cm}^{-1}$ , with a shift of  $-230$   $\text{cm}^{-1}$ . This  $\nu(\text{C}\equiv\text{C})$  band (fwhm  $\approx 120$   $\text{cm}^{-1}$ ) is much broader compared to those of the other three compounds (fwhm  $\approx 90$   $\text{cm}^{-1}$ ). The expected weaker  $\nu(\text{C}\equiv\text{C})$  band, however, was missing or too weak to be detected. This presence of a single  $\nu(\text{C}\equiv\text{C})$  band that is  $\sim 230$   $\text{cm}^{-1}$  shifted indicates that only one NN ligand was reduced in the  $^1\text{MLCT}$   $S_1$  state of IV.

In tens of picoseconds, both the  $\nu(\text{C}\equiv\text{C})$  and the  $\nu(\text{C}_6)_{\text{ring}}$  bands of the  $^1\text{MLCT}$   $S_1$  states decay away and features corresponding to the  $\nu(\text{CO}_2)$  and  $\nu(\text{CN}_2)$  stretches of the  $^3\delta\delta^*$   $T_1$  states emerge (red traces). Most apparent are the asymmetric  $\nu(\text{CO}_2)$  vibrations which shift  $\sim 15$   $\text{cm}^{-1}$  to higher

energy from the ground state (due to reduced back-bonding donation to the  $\pi^*$  orbital from the  $\text{Mo}_2\delta$  bond). A weak  $\nu(\text{C}\equiv\text{C})$  absorption band was also observed at  $\sim 2200\text{ cm}^{-1}$  in each of I–IV, close to the ground-state frequency. No  $\nu(\text{C}_6)_{\text{ring}}$  bands were spotted in the  ${}^3\delta\delta^* \text{ T}_1$  state. These spectral changes clearly depict the  ${}^1\text{MLCT } S_1 \rightarrow {}^3\delta\delta^* \text{ T}_1$  intersystem crossing process, in which the electron transfers back from the NN ligands to the  $\text{Mo}_2$  center.

**fs-TRIR Spectra of III in Solvents of Different Polarities.** In addition to THF at room temperature, the excited states of III were examined in toluene, DCM, and DMF solution to study the solvent influence. The spectra of III in toluene, THF, and DMF are shown in Figure 6. The DCM solution yielded no useable spectra, presumably due to photochemical decomposition. The fs-TRIR spectra of III in DMF solution were not collected in the  $1400\text{--}1500\text{ cm}^{-1}$  region due to strong solvent absorption.



**Figure 6.** fs-TRIR spectra of III in toluene, THF, and DMF solution with inset showing time delays. Inverted ground-state IR (gs-IR) spectra are shown by dotted lines.

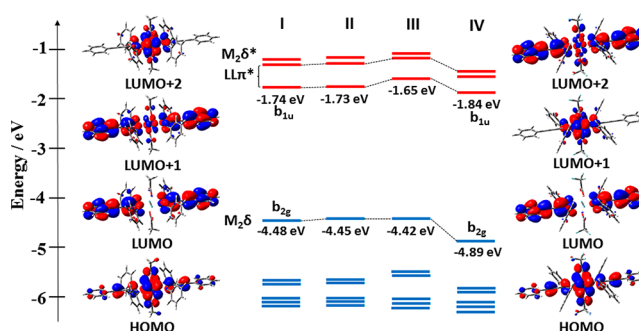
We have previously established that the  ${}^1\text{MLCT } S_1$  state of III is not symmetrical and that there should be a significant dipole moment associated with the excited state species.<sup>15</sup> That means, strong dipole–dipole interactions between the  ${}^1\text{MLCT } S_1$  state of III and the solvent molecules are expected. Interestingly, similar to the fs-TRIR spectra of III in THF, two  $\nu(\text{C}\equiv\text{C})$  bands (one weak, one strong) were observed for the  ${}^1\text{MLCT } S_1$  state of III in toluene and DMF. The  $\nu(\text{C}\equiv\text{C})$  bands are shifted by  $-75$  and  $-240\text{ cm}^{-1}$  in toluene and by  $-60$  and  $-220\text{ cm}^{-1}$  in DMF, which are quite similar compared to shifts of  $-70$  and  $-235\text{ cm}^{-1}$  in THF.

The  ${}^1\text{MLCT } S_1$  lifetime of III is  $15.5\text{ ps}$  in toluene and  $15.2\text{ ps}$  in DMF. The band shapes and early  ${}^1\text{MLCT}$  kinetics in toluene, THF, and DMF are different according to Figure 6 and the exponential fitting shown in Figure S15. However, the electron-density distribution in the  ${}^1\text{MLCT } S_1$  state of III are qualitatively the same in all three solvents. The shifts in DMF solution are smaller than those of the other two solutions,

which is surprising given that a polar solvent is expected to stabilize charge-transfer states.

**Density Functional Theory Calculations.** The electronic structures and electronic transitions of I–IV were investigated by density functional theory (DFT) and time-dependent density functional theory (TD-DFT) methods, respectively, using the B3LYP functional and the Gaussian 09 suite.<sup>23,24</sup>

The frontier molecular orbital (MO) energy diagram and MO surface plots of I–IV in the gas phase are shown in Figure 7.



**Figure 7.** MO energy levels and orbital surface plots of I–IV. LUMO+2, LUMO+1, LUMO, and HOMO are shown for I (left) and IV (right), respectively.

The electronic structures of  $\text{Mo}_2$  paddlewheel compounds are best described by interactions between the  $\text{Mo}_2$  quadruple bond orbitals and combinations of the ligand  $\pi^*$  orbitals.<sup>12</sup> With the NN ligands in the *trans* orientation, there are two possible NN  $\pi^*$  orbital combinations: the in-phase combination which does not interact with the  $\text{Mo}_2\delta$  bond orbital and the out-of-phase combination which does.

As shown in Figure 7, the highest occupied molecular orbital (HOMO,  $b_{2g}$ ) of I–IV is mainly composed of the  $\text{Mo}_2\delta$  orbital and the out-of-phase  $\text{LL}\pi^*$  orbital of the NN ligands. The lowest unoccupied molecular orbital (LUMO,  $b_{1u}$ ) of I–IV is mainly made up of the in-phase  $\text{LL}\pi^*$  orbital of the NN ligands. The HOMOs and LUMOs of I–IV appear rather similar except for the energy level difference. No significant contribution from the metal center or the carboxylate ligands was observed in the LUMO.

In compound IV, the  $\text{Mo}_2\delta$ -based HOMO is lower in energy by  $\sim 0.4\text{ eV}$  compared to those of I–III, which is attributed to the strong electron-withdrawing ability of the  $-\text{CF}_3$  groups. A similar effect is seen on the  $\text{Mo}_2\delta^*$ -based orbital, which moves lower than the out-of-phase NN  $\text{LL}\pi^*$  orbital. In compound III, the  $\text{LL}\pi^*$ -based LUMO is higher in energy by  $0.1\text{ eV}$  than those of I and II, which is attributed to steric influences of the  $-\text{T}^i\text{PB}$  groups over the NN ligands. The lower HOMO in IV and the higher LUMO in III lead to the increased energies of the MLCT transitions compared to I and II.

Time-dependent density functional theory (TD-DFT) calculations were used to examine the electronic transitions of I–IV. The results show that the MLCT transitions of I–IV can be described fully as HOMO to LUMO transitions. This is convenient since we can look at the composition of the LUMO for information on the excited-state electron population. The lowest two transitions of I–IV in the visible range are described in Table 3.

As shown in Table 3, the MLCT transitions of I–IV are predicted to be fully allowed transitions. The predicted energies

**Table 3. Electronic Transitions in I–IV As Predicted by TD-DFT<sup>a</sup>**

compounds	transition	origin	$\lambda$ / nm	$f$
I	$S_0 \rightarrow S_1$	HOMO $\rightarrow$ LUMO	530	1.0329
	$S_0 \rightarrow S_2$	HOMO $\rightarrow$ 0.92 LUMO+2 ( $\delta^*$ )	493	0.0003
II	$S_0 \rightarrow S_1$	HOMO $\rightarrow$ LUMO	534	1.0050
	$S_0 \rightarrow S_2$	HOMO $\rightarrow$ 0.92 LUMO+2 ( $\delta^*$ )	492	0.0002
III	$S_0 \rightarrow S_1$	HOMO $\rightarrow$ LUMO	528	0.8746
	$S_0 \rightarrow S_2$	HOMO $\rightarrow$ 0.92 LUMO+2 ( $\delta^*$ )	496	0.0002
IV	$S_0 \rightarrow S_1$	HOMO $\rightarrow$ 0.91 LUMO+1 ( $\delta^*$ )	489	0.0005
	$S_0 \rightarrow S_2$	HOMO $\rightarrow$ LUMO	475	1.0268

<sup>a</sup> $\lambda$  is the transition wavelength and  $f$  is the oscillator strength.

match reasonably well with those observed in absorption spectra. It is interesting that the MLCT transitions of I–III are predicted as  $S_0 \rightarrow S_1$  transitions, but the MLCT transition in IV is a  $S_0 \rightarrow S_2$  transition where it surpasses the energy of the  $\delta \rightarrow \delta^*$  transition. However, according to the emission spectra and the fs-TRIR spectra, the  $S_1$  states of I–IV are all indeed <sup>1</sup>MLCT in nature. No features of a <sup>1</sup> $\delta\delta^*$   $S_1$  state were observed, and these often have signature lifetimes of  $\sim 50$  ps.<sup>25</sup>

The composition of the frontier MOs in I–IV was studied by natural bonding orbital (NBO) analysis.<sup>26,27</sup> Due to the property differences between the metal orbitals and the ligand orbitals, they are often differentiated from each other. The percentage contributions of the metal orbitals are summarized in Table 4.

**Table 4. Metal-Orbital Contribution (%) to the Frontier MOs in I–IV**

compounds	HOMO	LUMO	LUMO+1	LUMO+2
I	65.2	1.2	6.5	76.8
II	65.1	1.1	6.6	76.7
III	64.6	1.3	6.0	76.9
IV	62.9	1.3	76.7	3.9

The composition of the HOMO and LUMO provide approximate descriptions of the electron-density distribution in the ground state and upon transitioning to the <sup>1</sup>MLCT state. As shown in Table 4, molybdenum orbitals and the NN ligand orbitals contribute consistently to the HOMO of I–IV, where the metal accounts for over 60% of the to the HOMO. The NN ligand contributions remain effectively the same despite changes in the carboxylate ligands.

As for the LUMO, the NN ligand orbitals contribute almost exclusively to the LUMO of I–IV while the metal orbitals contribute less than 1.5%, consistently. The trivial metal-orbital contributions to the LUMO of I–IV are mainly of  $Mo_2\pi_{sp}$  character, which is in strong contrast to the previously studied *trans*- $Mo_2L_2L'_2$  compound with (NPh)C(O)C $\equiv$ CPh and (NPh)C(S)C $\equiv$ CPh ligands. The latter displayed significant  $Mo_2\delta^*_{4d}$  contribution ( $\sim 3.4\%$  and  $\sim 10\%$ ) to the LUMO.<sup>15,16</sup> The limited metal-orbital contribution indicates that interligand electronic interactions can be rather weak in their MLCT states as there is no direct orbital contact.

## DISCUSSION

Using a donor–acceptor (D–A) description, *trans*- $Mo_2L_2L'_2$  compounds I–IV can be viewed as A–D–A type compounds. When the transferred electron is equally shared by the two

acceptor ligands in the MLCT state (i.e.,  $A^{-1/2}-D^+-A^{-1/2}$ ), only one IR-active  $\nu(C\equiv C)$  stretch is expected since the excited state is still symmetrical.<sup>16</sup> When the transferred electron is totally localized on one acceptor ligand (i.e.,  $A^-D^+-A$ ), two IR-active  $\nu(C\equiv C)$  stretches are expected but only one would be shifted from the ground-state wavenumber value.<sup>16</sup> In the intermediate case where the transferred electron is unequally shared (i.e.,  $A^{-(1-\alpha)}-D^+-A^{-\alpha}$ ), two discrete  $\nu(C\equiv C)$  bands are expected, both of which are shifted from the ground state.<sup>16</sup> The last scenario, which leads to an asymmetrical, semidelocalized excited state, has been studied before. In a paper from 2006, Terenziani et al. formulated the quantum mechanics of charge-transfer states in organic A–D–A and D–A–D type molecules and predicted unequal charge separation under certain circumstances.<sup>28</sup>

In inorganic chemistry, classic mixed valence (MV) compounds offer convenient precedents for dealing with unequal oxidation numbers and open-shell electronic structures of symmetrical molecular systems.<sup>29</sup> The Robin–Day scheme, which is frequently used to describe charge localization/delocalization in the ground state of MV complexes, can be adapted here to describe the electron-density distribution in the <sup>1</sup>MLCT excited states. In the same spirit, a class I <sup>1</sup>MLCT state corresponds to a transferred electron totally localized on one ligand, a class III <sup>1</sup>MLCT state to a transferred electron totally delocalized over two or more ligands, and a class II <sup>1</sup>MLCT state to the intermediate case. In analogy to intramolecular electron transfer in MV compounds, interligand electron transfer (ILET) can be used elegantly to describe interligand interactions in these MLCT excited states.

The observed shifts of the  $\nu(C\equiv C)$  bands in the <sup>1</sup>MLCT  $S_1$  states of I–IV are summarized in Table 5. The electron delocalization over the ligands in <sup>1</sup>MLCT  $S_1$  states are assigned with MV classifications as we described.

**Table 5. Shifts of  $\nu(C\equiv C)$  Bands and MV Assignments of the <sup>1</sup>MLCT  $S_1$  States of I–IV**

compounds	solvent	$\nu(C\equiv C)$ shift in $S_1$ state/ $cm^{-1}$	MV classification
I	THF	–65, –235	class II
II	THF	–64, –230	class II
III	toluene	–75, –240	class II
	THF	–70, –235	class II
IV	DMF	–60, –220	class II
	THF	–230	class I

The <sup>1</sup>MLCT  $S_1$  states of I–III were assigned as class II based on the two  $\nu(C\equiv C)$  bands observed in their fs-TRIR spectra, which suggests that the transferred electron is semilocalized between the two NN ligands. In IV, the single observed  $\nu(C\equiv C)$  band suggests that the <sup>1</sup>MLCT  $S_1$  state is a class I type, and the electron density is present exclusively on one of the two NN ligands. We have previously studied the fs-TRIR spectra of (DAniF)<sub>3</sub>Mo<sub>2</sub>(O<sub>2</sub>CC $\equiv$ CPh),<sup>30</sup> which is a similar compound functionalized with a single MLCT ligand, and the transferred electron is necessarily localized. A single  $\nu(C\equiv C)$  band was observed in the <sup>1</sup>MLCT  $S_1$  state with a shift of 240  $cm^{-1}$ , which agrees well with the fs-TRIR spectra of IV.

The surrounding solvent molecules appear to have limited influence on the mixed valence phenomena in the MLCT excited states. The <sup>1</sup>MLCT  $S_1$  states of III in the nonpolar toluene and the polar DMF solutions are both showing similar  $\nu(C\equiv C)$  band structures, as in THF. The coordinating and

close-contacting solvent molecules still strongly influence the excited-state relaxation as we see from the fluorescence quantum yields and relaxation dynamics. Yet the interligand interactions were hardly affected. With classic MV complexes, numerous studies have been conducted on the influence of solvents on electron transfer in the ground state of class II and III “borderline” compounds.<sup>29,31</sup> However, few simple and widely accepted conclusions were reached.

The conditions for localized and delocalized electron transfer in these excited-state mixed valence systems are not clear yet, and possibly involve multiple determining factors that control the interligand coupling. On the basis of our results, the metal–ligand orbital mixing in the frontier orbitals appears to be essential.  $\text{Mo}_2\delta^*$  contribution to the single-occupied LUMO in the  $^1\text{MLCT}$   $S_1$  states of the paddlewheel compounds, in particular, contributes to the strong interligand interactions.<sup>16</sup> The localized character of the  $^1\text{MLCT}$   $S_1$  states of I–IV is consistent with very limited contribution of the  $\text{Mo}_2\delta^*$  orbital in their LUMO.

The more localized character of the  $^1\text{MLCT}$   $S_1$  state of IV is intriguing compared to the other three. The DFT investigations on the ground states of I–IV fall short to reveal significant factors that would contribute to different long-range interligand interactions. Excited-state calculations of large TM complexes have not yet been efficient and reliable and thus were not employed in this study.

Since the influence on the mixed valence of I–IV comes from the functional groups on the  $\text{O}_2\text{C}-\text{X}$  ligands, the localization could be because of Coulombic interactions related to their electron-donating or electron-withdrawing abilities. The cross configuration of the auxiliary ligands and the MLCT ligands makes it tempting to compare the structure of compounds I–IV to that of a field-effect transistor (FET). It is possible that the ILET in *trans*- $\text{Mo}_2\text{L}_2\text{L}'_2$  compounds can be modulated just like a FET by Coulombic interactions through changing the “gate” potentials.

Due to the ultrashort lifetimes of the charge-transfer  $S_1$  states of TM complexes, charge-distribution patterns in the  $^1\text{MLCT}$  states of TM complexes have only been studied in a few systems (for example,  $\text{Ru}(\text{bpy})_3^{2+}$ ,  $\text{Fe}(\text{bpy})_3^{2+}$ , and  $\text{Re}(\text{bpy})(\text{CO})_3\text{Cl}$ ).<sup>32–34</sup> There have been reports on interligand interactions in the  $^3\text{MLCT}$   $T_1$  states of TM complexes that have much longer lifetimes, such as *trans* Pt(II) bis-acetylides and Re(I) heteroleptic polypyridine complexes.<sup>35,36</sup> The Pt(II) bis-acetylides bear similarities to the compounds studied here, and the current literature appears to favor the idea that the  $^1\text{MLCT}/^1\text{LLCT}$  excited states of these compounds are delocalized while the  $^3\text{MLCT}/^3\text{LLCT}$  excited states are localized.<sup>8,37,38</sup>

Dereka et al. recently reported a fs-TRIR study of the charge-transfer  $S_1$  state ( $\tau \approx 1$  ns) of a bis-cyanophenyl pyrrolopyrrole, an A–D–A type organic compound having two terminal  $\text{C}\equiv\text{N}$  reporter groups.<sup>39</sup> The charge-distribution pattern in its  $S_1$  state showed strong solvent dependence, where it displayed a single shifted  $\nu(\text{C}\equiv\text{N})$  band in nonpolar solvents (assigned to a delocalized  $S_1$  state) and two unequal  $\nu(\text{C}\equiv\text{N})$  bands in polar solvents (assigned to an asymmetrical, localized  $S_1$  state). While the same solvent dependence was not observed, the resemblance between the dual  $\nu(\text{C}\equiv\text{N})$  bands observed in this A–D–A type organic compound and the dual  $\nu(\text{C}\equiv\text{C})$  bands observed in I–III is remarkable.

Finally, the real-time symmetry breaking from the presumably delocalized Franck–Condon states of compounds I–IV

was not directly observed. The  $^1\text{MLCT}$   $S_1$  states of I–IV are already asymmetrical by  $\sim 200$  fs according to the fs-TRIR spectra, which is the time resolution of our instrument.

## CONCLUSION

The excited-state photophysics of the  $^1\text{MLCT}$   $S_1$  states of a series of *trans*- $\text{Mo}_2(\text{NN})_2(\text{O}_2\text{C}-\text{X})_2$  paddlewheel compounds were examined in solution at room temperature. The transferred electron from the dimetal center to the accepting ligands was followed with  $\nu(\text{C}\equiv\text{C})$  stretches located on the ligands using fs-TRIR spectroscopy. Although having two equivalent NN ligands, the  $^1\text{MLCT}$   $S_1$  states of these compounds are not symmetrical, and the transferred electron is shared unequally over the NN ligands.

The auxiliary carboxylate ligands  $\text{O}_2\text{C}-\text{X}$  displayed significant influences over the interligand interactions between the NN ligands. The observed  $\nu(\text{C}\equiv\text{C})$  patterns are explained in the construct of excited-state mixed valence in an A–D–A type donor–acceptor system. An interligand electron transfer (ILET) process is suggested to exist in the  $^1\text{MLCT}$   $S_1$  states of these compounds, which are modulated by the auxiliary ligands. DFT studies have shown evidence of the involvement of metal orbitals in the interligand electronic coupling.

The demonstration of control over excited-state charge dynamics through the choice of auxiliary ligands is unusual. Tuning interligand electronic interactions and excited-state symmetry provides a new perspective for designing more efficient molecular and nanomaterials for photochemical and photophysical applications.

## EXPERIMENTAL SECTION

**General Methods.** The dimolybdenum compounds are air and moisture sensitive. Syntheses and handling of metal compounds and ligands were conducted using standard Schlenk line and glovebox techniques. All of the solvents used were thoroughly dried and degassed. The homoleptic carboxylate precursor compounds were prepared from  $\text{Mo}(\text{CO})_6$  with established methods.<sup>17,20</sup> Details of the preparation of I–IV through exchange reactions are included in the Supporting Information.

NMR spectra of I–IV were recorded on a 400 MHz Bruker DPX Advance 400 spectrometer. All  $^1\text{H}$  NMR chemical shifts are in ppm relative to the protio impurity in chloroform-*d* at 7.26 ppm. Matrix-assisted laser desorption ionization time-of-flight (MALDI-TOF) mass spectra were obtained on a Bruker Microflex mass spectrometer. Dithranol was used as the matrix. Nine peptides with molecular masses ranging from 450 to 3100 Da were used for calibration.

**Crystal Structure Determination.** Suitable crystals of I, II, and IV for X-ray crystallography were grown by the diffusion of hexanes into concentrated THF solutions of I, II, and IV containing a small amount of DMSO. Diffraction data were collected on a Nonius Kappa Apex II CCD diffractometer with  $\text{Mo } K\alpha$  radiation. All work was done at 150 K using an Oxford Cryosystems Cryostream Cooler.

**Photophysical Measurements.** UV–vis absorption and emission spectra of I–IV were collected under air-free conditions using a 1 cm  $\times$  1 cm quartz cell equipped with a Kontes stopcock. The absorption data were recorded on a PerkinElmer Lambda 900 spectrometer, and the emission data were recorded on a SPEX Fluoromax-2 fluorometer.  $\text{Ru}(\text{bpy})_3(\text{PF}_6)_2$  was used as a standard in the quantum-yield determination. FT-IR spectra of I–IV were collected on a PerkinElmer Spectrum GX FTIR spectrometer. Sample solutions were sealed in a PerkinElmer semidemountable cell with  $\text{CaF}_2$  windows and a 0.1 mm Teflon spacer.

Femtosecond time-resolved infrared (fs-TRIR) experiments were performed using a Ti:sapphire oscillator/regenerative amplifier combination (1 kHz, fwhm  $\sim 200$  fs), as previously described.<sup>40</sup> The mid-IR probe/reference beams were generated through a difference

frequency generation (DFG) mechanism and was collected with a Triax 320 spectrometer, where they were dispersed with a grating onto separate arrays of a liquid nitrogen-cooled HgCdTe detector ( $32 \times 2$  pixels). Samples of I–IV were prepared to have  $\sim 1$  absorption at their MLCT peaks. Sample solutions were kept air-free in a PerkinElmer semidemountable cell with a 0.1 mm Teflon spacer and two 4.0 mm CaF<sub>2</sub> windows. The excitation-beam power at the sample was tuned to 1–2  $\mu\text{J}$ .

All time-resolved spectra were plotted in Igor Pro 6.0. Kinetic traces were fitted with a global fitting package with a sum of exponentials,  $S(t) = \sum A_i \exp(-t/\tau_i) + C$ , where  $A_i$  is the amplitude,  $\tau_i$  is the lifetime, and  $C$  is an offset.

**Computational Methods.** Compounds I–IV were studied using density functional theory (DFT) methods with the B3LYP functional in the Gaussian09D01 suite. The 6-31G(d) basis set was used for the H, C, N, O, and F atoms, and the Stuttgart–Dresden (SDD) effective core potentials were used for Mo atoms.<sup>41</sup> Atomic coordinates from crystal structures were used as the initial input for structure optimization. Force constant and vibrational frequency analyses were performed on each compound to make sure that all structures are optimized to a global minimum. Isosurface contour plots of I–IV were created with Gaussview 5.0.8 with isovalues at 0.02. NBO 3.1 was used for the analysis of molecular orbital composition.<sup>26</sup> The polarizable continuum model (PCM) was used to account for the solvent effects, and default parameters were used for all the solvents.<sup>42</sup> A summary of solvent polarity moments, HOMO/LUMO energy levels and calculated MLCT wavelengths of III are presented in Figure S4 and Table S3 in the Supporting Information.

The vibration frequency analyses of I–IV were performed and used in the FT-IR spectra interpretation. All frequencies were scaled by a factor of 0.961 as suggested.<sup>43</sup>

## ■ ASSOCIATED CONTENT

### ■ Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.inorgchem.7b01164.

Crystal structure refinement information on I, II, and IV, NBO analysis, NMR, and MALDI-TOF spectra of I–IV (PDF)

### ■ Accession Codes

CCDC 1541002–1541003 and 1541005 contain the supplementary crystallographic data for this paper. These data can be obtained free of charge via [www.ccdc.cam.ac.uk/data\\_request/cif](http://www.ccdc.cam.ac.uk/data_request/cif), by emailing [data\\_request@ccdc.cam.ac.uk](mailto:data_request@ccdc.cam.ac.uk), or by contacting The Cambridge Crystallographic Data Centre, 12 Union Road, Cambridge CB2 1EZ, UK; fax: +44 1223 336033.

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### ■ Notes

The authors declare no competing financial interest.

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