

an increasing energy separation of levels—a consequence of the positive anharmonicity. This is particularly evident in the ground state with the relatively broad minimum of the ground-state potential.

One motivation behind the present work has been the pursuit of accurate spatially resolved temperatures of flames based on sequence structure peak areas. Franck–Condon factors for the sequence structure were generated from the normalized coefficients used in the diagonalization procedure and scaled with room temperature Boltzmann factors. These total combined factors along with their associated Franck–Condon factors are given in Table II and demonstrate the good agreement between the theoretical and experimentally acquired peak areas. This agreement indicates little contribution of kinetic terms such as geometric saturation or ac stark shifts to the overall dependences of the peak areas and demonstrates the feasibility of nonintrusively determining spatially resolved temperatures.

Sequence structure similar to that of the  $A \leftarrow X$  transition has also been observed in the hot-band region of the  $B \leftarrow X^{9,10}$  state. Analysis of the spectra in Figure 5 of ref 9 (one-photon absorption

around 260 nm) yielded the same ground-state potential as the present study and, within experimental uncertainty, a similar excited-state potential as that of the A state. This suggests a minor deviation in the influence of the B state's electronic manifold (composed of nitrogen's  $3p_{x,y}(+)$  orbitals) over the A state's manifold (composed of nitrogen's  $3s(+)$  orbitals) on the torsional vibration.

### Conclusion

Franck–Condon factors and energy positions for a number of levels of the  $\nu_{13}$  torsional mode have been determined for the  $A \leftarrow X$  transition of Dabco. This data has been shown useful in measuring vibrational levels over a wide temperature range in difficult environments. The technique is simple and ions, as in the present study, or fluorescence (centered around 300 nm)<sup>14</sup> can be monitored with gated electronics. Further work on preheating and pyrolysis of Dabco in atmospheric and low-pressure flames is now in progress.

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## H/D Isotope Exchange in the $D_2$ – $H_2O$ System under the Influence of Ultrasound

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Water was irradiated by 300-kHz ultrasound under atmospheres of argon, deuterium, and mixtures of the two gases, and the isotopic composition of the gas was determined. The yields of  $H_2$  and HD increase with increasing deuterium content of the gas mixture, reach maxima at 35%  $D_2$  ( $H_2$ ) and 60%  $D_2$  (HD), and fall off at higher  $D_2$  concentrations. The yield of the  $D_2$  consumption peaks at 50%  $D_2$ . The absolute yields are greater by almost one order of magnitude than those of ordinary sonolytic reactions such as  $H_2O_2$  formation. The results are explained by reactions of the H and OH radicals generated in the sonolysis of water. In the absence of deuterium, a large percentage of these radicals recombine.  $D_2$  at low concentrations scavenges mainly OH radicals, the principal products being  $H_2$  and HDO. At higher  $D_2$  concentration,  $D_2$  also scavenges H atoms, the principal products now being HD and HDO. The decrease in all yields at higher  $D_2$  concentrations is explained by the lower temperatures reached in the gas bubbles, which are formed by cavitation.

### Introduction

The chemical effects of ultrasonic waves in aqueous solutions are brought about by the high temperatures and pressures which exist for a short time during the nearly adiabatic compression phase of oscillating or collapsing gas bubbles.<sup>1</sup> When water is irradiated under argon,  $H_2$  and  $H_2O_2$  are generated and a smaller amount of  $O_2$  is formed. Dissolved substances often are oxidized or reduced, the yields generally being of the same order of magnitude as that of the  $H_2O_2$  formation in pure water.<sup>2</sup> The initial step of the sonolysis of water is the thermal dissociation



Both H and OH recently were detected in ESR spin-trapping experiments.<sup>3</sup>  $H_2$  and  $H_2O_2$  are formed via the combination of radicals. It was previously shown that the yields of some chemical reactions are substantially increased when gaseous radical scavengers such as oxygen or methane are present in moderate concentrations in an argon atmosphere.<sup>4-6</sup> In the present work,

deuterium was used as a scavenger, and the D–H exchange between  $D_2$  and  $H_2O$  was investigated by measuring the isotopic composition of the gas after irradiation. In these experiments, water was irradiated under various mixtures of argon and deuterium. At  $D_2$  contents greater than 10%,  $H_2O_2$  was no longer a product, i.e. in most of these experiments H–D exchange was the only process occurring.

The  $D_2$ – $H_2O$  exchange reaction under the influence of  $\gamma$ -rays has previously been reported.<sup>7</sup> It was found to be initiated by the free radicals from the radiolysis of water. It is known from a number of studies that the reactions in sonolysis occur to a large extent in the gaseous phase. This has also to be taken into consideration in discussing the sonolytically initiated exchange reaction.

### Experimental Section

Water was deaerated by evacuation with a diffusion pump. Deaerated water (100 mL) was saturated with the desired  $D_2$ –Ar mixture. The glass irradiation vessel had a volume of 66.5 cm<sup>3</sup> and could be closed during irradiation. It was completely filled with the  $D_2$ –Ar solution displacing the argon used to purge the vessel. Finally, 15 cm<sup>3</sup> of solution was displaced by the gas mixture. By this procedure, 51.5 cm<sup>3</sup> of  $H_2O$  was in the vessel

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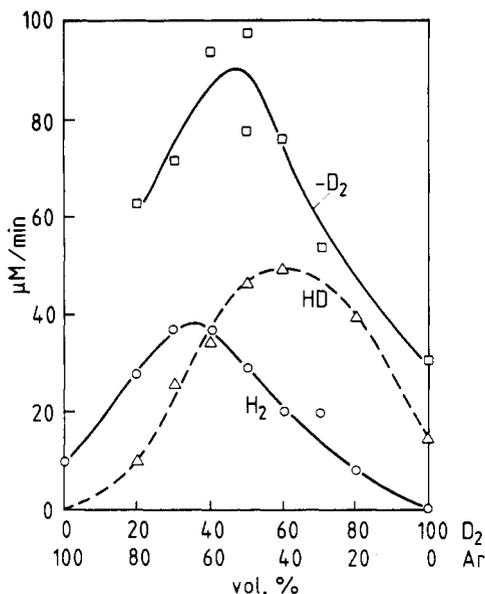
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**Figure 1.** Rates of formation of  $H_2$  and HD and of the consumption of  $D_2$  in the ultrasonic irradiation of water under various mixtures of argon and deuterium.

with  $15\text{-cm}^3$  gas phase. After irradiation,  $15\text{ cm}^3$  of water was added to the solution and the gas pushed into a syringe. The gas finally was injected into a mass spectrometer for isotopic analysis. All these procedures were made using standard Van Slyke and syringe techniques.<sup>8</sup> Much care was taken to make sure that no traces of air were present during irradiation and analysis.

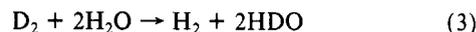
The ultrasound was delivered by the 300-kHz quartz generator previously described.<sup>6</sup> The power dissipation in the liquid was 12 W as determined from the rise in temperature in an experiment without a cooling water bath. During the irradiation the increase in temperature was less than  $1^\circ\text{C}$ .

## Results

The rates of formation of HD and  $D_2$  and that of the consumption of  $D_2$  are plotted in Figure 1 as functions of the composition of the argon–deuterium mixture. The conversion of  $D_2$  was between 5 and 15% in these experiments. The rate of  $D_2$  consumption had therefore to be calculated as the difference in two large numbers. This explains the rather large scattering of the data for  $-D_2$  in Figure 1. Under pure argon, the rate of  $H_2$  formation was  $10\ \mu\text{M min}^{-1}$ .  $H_2O_2$  was formed at a rate of  $8.0\ \mu\text{M min}^{-1}$  under these conditions, the difference in oxidation equivalents between the two yields being due to  $O_2$  formation. The addition of  $D_2$  caused the  $H_2$  yield to increase. The  $O_2$  and  $H_2O_2$  yields (not shown in the figure) rapidly decreased to practically zero at 10%  $D_2$ . A maximum yield of  $37\ \mu\text{M min}^{-1}$  for  $H_2$  was observed at 35%  $D_2$  in the gas mixture. Under pure  $D_2$ , no  $H_2$  was found.

The HD yield increased in a somewhat different way, a maximum yield of  $49\ \mu\text{M min}^{-1}$  being attained at 60%  $D_2$ . A noticeable HD yield was still observed under pure  $D_2$ . The rate of decrease in  $D_2$  concentration also went through a maximum (at 50%  $D_2$ ). The rate of  $D_2$  consumption was much greater than that of HD formation. We conclude that most of the D atoms consumed appeared in the form of HDO in the liquid. No experiments were undertaken to trace HDO as there are no other deuterium-containing products besides HD and HDO. Note that the rate of  $D_2$  consumption at the maximum ( $90\ \mu\text{M min}^{-1}$ ) is about ten times greater than that of a more usual sonochemical reaction such as  $H_2O_2$  formation.

Two overall processes for isotopic exchange between deuterium and water may be written:



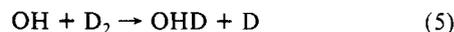
As all three products, i.e.  $H_2$ , HD, and HDO, appear, one has to postulate that both processes, 2 and 3, occur in the ultrasonic irradiation of water under an argon–deuterium atmosphere. In mixtures with comparable percentages of the two gases, processes 2 and 3 are roughly equal. In the irradiation under pure  $D_2$ , where  $H_2$  is not formed, only exchange process 2 takes place.

## Discussion

The maximum yields for  $H_2$  and HD formation and  $D_2$  consumption are much higher than the yields of  $H_2$  and  $H_2O_2$  observed in the irradiation of water under pure argon. This effect is understood in terms of the radical recombination reaction



which occurs very efficiently in the absence of a gaseous scavenger.<sup>5,6</sup> In the presence of sufficient  $D_2$ , all the primary H and OH radicals are scavenged to form the final products of HD exchange with large yields:



At low  $D_2$  concentrations, mainly OH radicals are scavenged (eq 5), the remaining H atoms forming  $H_2$  and HD (eq 7 and 8). With increasing  $D_2$  concentration, reaction 6 also becomes important, the result being that the HD yield becomes more important than that of  $H_2$  at moderate  $D_2$  concentrations. These considerations explain why the maximum of the curve for HD in Figure 1 appears at a higher  $D_2$  concentration than the maximum of the curve for  $H_2$ . Finally, when the  $D_2$  concentration in the gas mixture becomes too high, a lower temperature is reached in the nearly adiabatic compression phase of the gas bubbles. This effect finally leads to a decrease in all yields in Figure 1 at the higher  $D_2$  concentrations.

A rough estimate of the radicals that recombine (eq 4) may be made. The maximum  $H_2$  yield at 35%  $D_2$  in Figure 1 is about four times greater than the  $H_2$  yield in the irradiation under pure argon, and the HD yield at 35%  $D_2$  is about twice the yield of  $H_2$  in the absence of  $D_2$ . We assume that the yield of the primary decomposition reaction (eq 1) is not yet significantly decreased at 35%  $D_2$ . It can then be derived from the above mechanism that about 80% of the radicals generated initially recombine in the absence of  $D_2$ . This corroborates our previous conclusions about the importance of the recombination reaction.<sup>5,6</sup>

We must still discuss another thermal decomposition reaction which may take place in the presence of  $D_2$ . Since the bond strength in D–D is lower than that in H–OH, it seems plausible that deuterium molecules also are dissociated in the gas bubbles:



In fact, it is known that the reducing effects of ultrasound are strongly enhanced in the presence of hydrogen in the argon atmosphere.<sup>9</sup> The D atoms formed contribute to the isotopic exchange by reacting with water molecules according to



This reaction is endothermic by only 0.6 eV, and it cannot be excluded that part of the HD results from it. However, the reasons may be mentioned for reaction 10 not being of great importance. In the first place, the formation of  $H_2$  cannot be explained in this way. Secondly, HD should be formed in a very large yield as the chain reaction of eq 10 and 5 becomes possible. The yield of HD, however, did not substantially exceed that of  $H_2$ , i.e. the yield of

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a product which did not stem from a chain reaction.

We conclude that the isotope exchange between  $D_2$  and  $H_2O$  can be explained by free-radical reactions on the basis of principles which have been derived previously in studies on other chemical reactions. These principles include the dissociation of  $H_2O$  in the collapsing gas bubble, the recombination reaction of radicals in the absence of a gaseous scavenger, and the dependence of these

yields on the composition of the argon- $D_2$  gas atmosphere under which water is sonolyzed.

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