

# Product Branching Fractions in the Reaction of NH (ND) ( $^3\Sigma^-$ ) with NO

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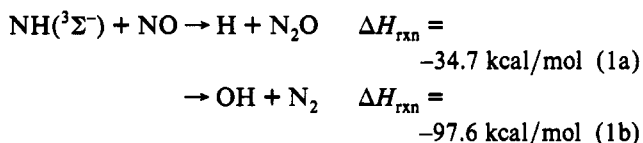
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We have carried out a combined experimental and theoretical study of the reaction of NH (ND) ( $^3\Sigma^-$ ) with NO aimed at understanding the product distribution from that reaction. The reaction was studied at room temperature using the discharge flow technique with mass spectrometric detection of the reaction products. Measured product branching fractions at room temperature for production of  $\text{N}_2\text{O} + \text{H}$  (D) were  $0.8 \pm 0.4$  for NH ( $^3\Sigma^-$ ) + NO and  $0.87 \pm 0.17$  for ND( $^3\Sigma^-$ ) + NO ( $1\sigma$  statistical errors). Stationary points on the HNNO  $^2\text{A}'$  potential energy surface were characterized using the Gaussian 2 *ab initio* method. The initial addition of NH( $^3\Sigma^-$ ) to NO on the  $^2\text{A}'$  surface is predicted to proceed without a barrier to form *trans*-HNNO; reaction to produce the *cis* isomer is predicted to have a barrier of 2.9 kcal/mol. The *cis*- and *trans*-HNNO are predicted to be at  $-48.9$  and  $-56.0$  kcal/mol relative to the separated reactants. Transition states with energies of  $-25.3$  and  $-17.9$  kcal/mol were located for dissociation of the *cis* isomer into  $\text{H} + \text{N}_2\text{O}$  and  $\text{OH} + \text{N}_2$ , respectively. The transition state for interconversion of the isomers was calculated to be at approximately  $-30.8$  kcal/mol. The *trans*-HNNO was found to isomerize to the *cis* form before decomposing. The potential energy surface calculated explains the major features of the reaction.

## Introduction

The reaction of NH( $^3\Sigma^-$ ) with NO is a moderately exothermic, radical-radical reaction with two thermodynamically open channels:<sup>1,2</sup>



It has been claimed to be one of the principal sources of  $\text{N}_2\text{O}$  from combustion,<sup>3,4a</sup> although other flames studies have suggested that the primary product of the reaction is  $\text{N}_2$ , and little  $\text{N}_2\text{O}$  is formed. Its reaction rate coefficient has been studied before near room temperature,<sup>5–9</sup> in shock heated gas,<sup>10,11</sup> and in flame studies.<sup>12–14</sup> Reported room temperature rates are in the range  $k(298) = (3.8\text{--}5.8) \times 10^{-11} \text{ cm}^3/(\text{molecules s})$ . The reaction rate coefficient is reported to be temperature independent in the temperature range 269–377 K,<sup>8</sup> high-temperature reaction rate coefficients of  $k(2200\text{--}3350 \text{ K}) = 2.8 \times 10^{-10} \exp(-6400/T) \text{ cm}^3/(\text{molecule s})$ <sup>10</sup> and  $k(3500 \text{ K}) = (1.3 \pm 0.1) \times 10^{-10} \text{ cm}^3/(\text{molecules s})$ <sup>11</sup> have been reported. Recently, Patel-Misra and Dagdigian<sup>15</sup> have measured nascent OH product distributions arising from the reaction. There have been some direct studies examining the products from the reactions of NH( $^3\Sigma^-$ ) with NO.<sup>8–11</sup> Harrison and co-workers looked for OH product with laser-induced fluorescence but were not successful in detecting it. Despite this negative result, they concluded that  $\text{N}_2 + \text{OH}$  was probably the major product of the reaction.<sup>8</sup> Yamasaki et al. report exclusive production of  $\text{N}_2 + \text{OH}$  at room temperature and claim that  $\text{N}_2\text{O} + \text{H}$  arises from the reaction of NH( $^1\Delta$ ) with NO.<sup>9</sup> Mertens et al. report the branching fraction into  $\text{N}_2 + \text{OH}$  to be  $0.27 \pm 0.10$  for  $2350 \leq T \leq 3040 \text{ K}$ ,<sup>10</sup> which is in good agreement with Yokoyama et al.'s branching fraction of  $0.32 \pm 0.07$  at 3500 K.<sup>11</sup> Patel-Misra and Dagdigian's work<sup>15</sup> establishes OH as a direct product of the reaction but is mute as to its branching fraction.

There have been a number of theoretical studies of the HNNO  $^2\text{A}'$  potential energy surface. Melius and Binkley<sup>16</sup> used the BAC-

MP4 method to calculate reactants, products, and intermediates involved in the NH + NO reaction. Those calculations have since been revised due to newer bond additivity correction procedures; results of the newer calculations are reported in Miller and Melius<sup>4</sup> who also report of calculations of the product branching fraction for the reaction of NH with NO, based on the BAC-MP4 structures and geometries. Fueno et al.<sup>17</sup> used MRD-CI//HF/4-31G(d,p) to calculate energies for stationary points on the HNNO  $^2\text{A}'$  surface, as well as the energy of the *cis* and *trans* isomers of HNNO  $^2\text{A}''$ , which were predicted to lie higher than the  $^2\text{A}'$  states of the two isomers. Harrison and MacLagan<sup>18</sup> reported calculations of geometries at HF/6-31G(d), HF/6-311G(d,p), and MP2/6-31G(d) levels, with energies calculated using both MP2 and MP4 methods. More recently, Walch<sup>19</sup> has calculated a number of points on the HNNO  $^2\text{A}'$  surface, using CASSCF/CCI. The results of these studies vary widely, although they all find that the transition state for production of  $\text{H} + \text{N}_2\text{O}$  is lower in energy than the one leading to  $\text{OH} + \text{N}_2$ .

The present study has sought to address questions concerning the products of the NH( $^3\Sigma^-$ ) + NO reaction by a combined experimental and theoretical approach. We have carried out experiments using the discharge flow technique to directly measure products of the reaction and have used the Gaussian 2 method to characterize stationary points on the HNNO potential energy surface.

## Experimental Section

Experiments were carried out in our double-injector, discharge-flow apparatus. The basic apparatus has been described in detail earlier.<sup>20</sup> The flow tube, a 22-mm-i.d. quartz tube, 530 mm long, was suspended vertically in a vacuum chamber constructed from 4 5/8-in. Conflat hardware. This study utilized a double injector consisting of a 6-mm-o.d. quartz tube mounted concentrically within a 13-mm quartz injector. This assembly was mounted on the flow tube axis. The flow tube ended either at the tip of a 2-mm-i.d. conical sampling nozzle for the mass spectrometer or approximately 1 cm above a 1-mm flat sampling nozzle. The flow tube was pumped out through the outer vacuum jacket, using an oil-free DryStar mechanical pump. Reagent flow rates were measured and controlled by MKS 1359 and Tylan

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FC-2900 transducers connected to MKS 246 flow controllers. Pressures were measured with an MKS 390HA capacitance manometer and controlled with a variably-open butterfly valve connected to an MKS 252 exhaust valve controller. Total pressures of 0.75 and 2.5 Torr were used, with typical flow velocities of 1–4 m/s. Typical initial reagent concentrations were  $5 \times 10^{12} \leq [\text{NO}] \leq 1 \times 10^{13}$  molecules/cm<sup>3</sup> and  $5 \times 10^{11} \leq [\text{NH}_3] \leq 2 \times 10^{12}$  molecules/cm<sup>3</sup>.

Reactants and products were detected mass spectrometrically, using modulated molecular beam sampling of the flow tube. The mass spectrometer was doubly-differentially pumped, with a 1- or 2-mm nozzle, 1-mm skimmer (Beam Dynamics), and 1.5-mm collimating orifice. An electron energy of 18 eV and an emission current of 0.75 mA were used in this study.  $^{15}\text{N}^{14}\text{N}^{18}\text{O}$ ,  $^{15}\text{N}^{14}\text{N}$ , and  $\text{D}_2^{18}\text{O}$  (or  $\text{H}_2^{18}\text{O}$ ) ions were pulse counted, using an SRS 400 gated photon counter interfaced to an IBM AT computer. Modulated beam sampling with ac detection of the ion signals considerably improved the S/N, as well as offering a large measure of immunity from changes in background gas composition.

NH radicals were formed by the reaction of  $\text{NH}_3$  with excess F atoms.<sup>21,22</sup> F atoms were formed by microwave discharge of a highly dilute  $\text{F}_2/\text{He}$  mixture using a Beenakker-type microwave cavity operating at 30–40 W of microwave power.

Reagents used were  $\text{NH}_3$  (Matheson),  $\text{ND}_3$  (MSD Isotopes),  $^{15}\text{N}^{18}\text{O}$  (Isotec),  $\text{N}_2\text{O}$  (Matheson),  $\text{N}_2$  (Scott Specialty Gases, 99.9995%), 5%  $\text{F}_2/\text{He}$  (Spectra Gases), and He (Matheson, 99.9999%). All reagents were used without further purification. Before experiments with  $\text{ND}_3$  the flow system was deuterated by passing  $\text{D}_2\text{O}$  through it overnight.

### Computational Details

Calculations were carried out using the Gaussian 90<sup>23</sup> and Gaussian 92<sup>24</sup> programs. We used the G2 method, as outlined by Pople and co-workers.<sup>25</sup> G2 uses a series of calculations to approximate a QCISD(T)/6-311+G(3df,2p)//MP2/6-31G(d) calculation with an additional "higher-order correction" based on the number of paired and unpaired electrons. In brief, G2 energies were obtained by following these steps: (1) Vibrational frequencies and zero-point energies were determined by scaling the results of HF/6-31G(d)//HF/6-31G(d) calculations. (2) Geometries were optimized at MP2/6-31G(d) level. (3) Energies were calculated using MP4/6-311G(d,p)//MP2/6-31G(d). (4) Corrections due to basis set incompleteness were evaluated by MP4/6-311+G(d,p)//MP2/6-31G(d), MP4/6-311G(2df,p)//MP2/6-31G(d), and MP2/6-311+G(3df,2p)//MP2/6-31G(d). (5) Corrections due to further electron correlation were evaluated by QCISD(T)/6-311G(d,p)//MP2/6-31G(d). (6) "High-level corrections" were evaluated based on the number of  $\alpha$  and  $\beta$  valence electrons.

The method has been shown to yield atomization energies with a average absolute deviation of 0.9 kcal/mol for a large number of species containing first-row elements. We have carried out additional studies of the applicability of the G2 method to transition-state structures.<sup>26</sup> We found that it performed well in predicting transition-state properties, although its performance for transition states could be enhanced by using QCISD/6-311G-(d,p) optimized geometries and frequencies instead of the MP2 optimized geometries and scaled Hf frequencies of the G2 method. We have used this G2//QCISD/6-31G(d,p), or G2Q, method to characterize the transition states leading to  $\text{N}_2 + \text{OH}$  and  $\text{H} + \text{N}_2\text{O}$ .

### Results and Discussion

**Experiment.** We measure the branching fraction of the title reactions directly, using mass spectrometric detection of the stable

reaction products,  $\text{N}_2\text{O}$  and  $\text{N}_2$ . NH (ND) is mixed with an excess of NO, driving the reaction rapidly to completion. The resulting gas mixture is sampled using molecular beam sampling and analyzed using mass spectroscopy. Relative detection efficiencies for the various species were measured directly using calibration mixtures containing known partial pressures of each of the stable species detected.

We use the reaction of ammonia with excess fluorine atoms as our NH radical source. This source was used by Hack et al.<sup>22</sup> in a study of the reaction of  $\text{NH}(^3\Sigma)$  with  $\text{O}_2(^1\Delta)$ . They found that it was possible to make a sizable excess of NH over  $\text{NH}_2$  by proper choice of the  $[\text{NH}_3]/[\text{F}]$  ratio. There are, unfortunately, constraints which limit the maximum  $[\text{NH}_3]_0$  and hence the maximum  $[\text{NH}]$  possible. The first is the well-known production of solid  $\text{NH}_4\text{F}$  if the concentrations of  $\text{NH}_3$  and HF become too high.<sup>27</sup> We also observed production of  $\text{N}_2\text{H}_x$  species for  $[\text{NH}_3]_0 > 2.5 \times 10^{12}$  molecules/cm<sup>3</sup>. To avoid these problems, we used low concentrations of ammonia,  $5 \times 10^{11} \leq [\text{NH}_3]_0 \leq 2 \times 10^{12}$  molecules/cm<sup>3</sup>. The  $[\text{NH}_3]/[\text{F}]$  ratio was adjusted by adjusting  $[\text{F}]$  to achieve a maximum in the  $\text{N}_2\text{O}$  signal.

The use of isotopically-labeled  $^{15}\text{N}^{18}\text{O}$  allowed us to separate signal due to the  $\text{NH} + \text{NO}$  reaction from most background signals. This did not help to separate out contributions due to small amounts of  $\text{NH}_2$  reacting with NO producing  $\text{N}_2 + \text{H}_2\text{O}$ . Approximately 30% of the signal at  $m/e = 29$  ( $^{15}\text{N}^{14}\text{N}$ ) was due to this reaction. To correct for it, we measured the  $m/e = 20$  ( $^{22}$ )/ $m/e = 29$  ( $\text{H}_2^{18}\text{O}$  ( $\text{D}_2^{18}\text{O}$ )/ $^{15}\text{N}^{14}\text{N}$ ) ratio with  $[\text{F}] \ll [\text{NH}_3]$ , ensuring that no NH was present. This ratio, coupled with our measurement of  $m/e = 20$  ( $^{22}$ ), allowed us to correct for this side reaction. In work with  $\text{NH}_3$  there was a further complication, since  $m/e = 20$  contains contributions from  $\text{H}_2^{18}\text{O}$  and HF. We corrected for this by measuring  $m/e = 20$  with NO on and off, but the additional source of error shows up as a larger uncertainty in our branching fraction measurement for NH relative to that for ND. The resulting product branching fractions, measured at room temperature, for production of  $\text{N}_2\text{O} + \text{H}(\text{D})$  were  $0.84 \pm 0.4$  for  $\text{NH}(^3\Sigma^-) + \text{NO}$  and  $0.87 \pm 0.17$  for  $\text{ND}(^3\Sigma^-) + \text{NO}$  (1 $\sigma$  statistical errors).

There have been reports that the excess-F +  $\text{NH}_3$  system produces  $\text{NH}(^1\Delta)$  in addition to  $\text{NH}(^3\Sigma)$ .<sup>28</sup> In order to rule out the  $\text{NH}(^1\Delta) + \text{NO}$  reaction as an interference in this work, we substituted  $\text{N}_2$  for He as carrier in a series of experiments.  $\text{N}_2$  quenches  $\text{NH}(^1\Delta)$ , with a rate of  $k_q = (6-9) \times 10^{-14}$  cm<sup>3</sup>/(molecules s),<sup>29</sup> fast enough to ensure that all the  $\text{NH}(^1\Delta)$  produced in the source is quenched to  $\text{NH}(^3\Sigma)$  before addition of NO. The substitution did not affect the measured branching fraction, demonstrating that the product branching fraction measured here is unaffected by the possible presence of  $\text{NH}(^1\Delta)$ .

**Theory.** Addition of  $\text{NH}(^3\Sigma)$  to  $\text{NO}(^2\Pi)$  gives rise to 12 surfaces. In  $C_1$  symmetry these are two doubly-degenerated  $^2A$  surfaces and two 4-fold-degenerate  $^4A$  surfaces. Since connecting the quartet surfaces with the ground-state products violates spin conservation, we will not consider the quartet surfaces further. Fueno et al.<sup>17</sup> have calculated structures for the *cis*- and *trans*-HNNO equilibrium geometries and find them to be planar, of  $^2A'$  and  $^2A''$  symmetries ( $C_2$  point group). Patel-Misra and Dagdigan's<sup>15</sup> analysis of the OH product from the  $\text{NH}(^3\Sigma) + \text{NO}$  reaction leads them to also conclude that the transition state for the  $\text{N}_2 + \text{OH}$  product channel is planar. Fueno et al.<sup>17</sup> found the  $^2A'$  surface to be more stable than the  $^2A''$  surface, in agreement with Melius<sup>16</sup> and later work by Walch.<sup>19</sup> Walch<sup>19</sup> has examined the correlations between the HNNO surfaces and products. He concluded that the  $^2A''$  surface does not correlate with  $\text{H} + \text{N}_2\text{O}$  and is unfavorable toward formation of  $\text{N}_2 + \text{OH}$ , while the  $^2A'$  could lead to both product channels. We have

TABLE 1: G2 Energies<sup>a</sup>

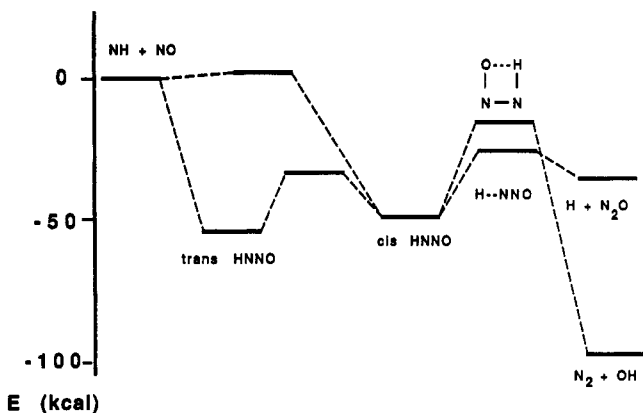
species	$E[\text{MP4}/6\text{-}311\text{G(d,p)}]^b$	$\Delta(+)^c$	$\Delta(2\text{df})^c$	$\Delta(\text{QCI})^c$	$\Delta(\text{G2})^c$	$\Delta(\text{HLC})^c$	$\Delta(\text{ZPE})$	G2 energy <sup>b</sup>
N <sub>2</sub> O	-184.311 97	-7.37	-99.18	19.53	9.13	-40.00	10.99	-184.437 13
<i>trans</i> -HNNO	-184.821 84	-11.17	-100.00	-5.01	-11.70	-40.19	18.61	-184.971 31
<i>cis</i> -HNNO	-184.815 88	-8.62	-99.51	-2.88	-11.56	-40.19	18.61	-184.960 03
<i>trans</i> -HN-NO	-184.702 99	-12.92	-97.61	-32.37	-10.82	-40.19	14.87	-184.882 03
<i>cis</i> -HN-NO	-184.706 86	-12.53	-95.36	-26.01	-10.79	-40.19	14.29	-184.877 46
<i>cis-trans</i> ts	-184.784 41	-9.05	-102.10	-0.94	-11.60	-40.19	17.06	-184.931 24
NN-OH( <sup>2</sup> A')	-184.753 05	-8.62	-99.19	-8.51	-9.87	-40.19	12.03	-184.907 40
QCISD/6-311G(d,p)	-184.748 63	-9.29	-98.25	-19.15	-9.71	-40.19	14.53	-184.910 69
H-NNO( <sup>2</sup> A')	-184.770 77	-8.16	-102.29	-1.85	-9.84	-40.19	10.77	-184.922 33
QCISD/6-311G(d,p)	-184.767 92	-7.89	-101.84	-6.84	-9.75	-40.19	11.97	-184.922 46

<sup>a</sup> The notation used is as follows:  $\Delta(+)=E[\text{MP4}/6\text{-}311\text{G(d,p)}]-E[\text{MP4}/6\text{-}311\text{G(d,p)}]$ ;  $\Delta(2\text{df})=E[\text{MP4}/6\text{-}311\text{G(2df,p)}]-E[\text{MP4}/6\text{-}311\text{G(d,p)}]$ ;  $\Delta(\text{QCI})=E[\text{QCISD(T)}/6\text{-}311\text{G(d,p)}]-E[\text{MP4}/6\text{-}311\text{G(d,p)}]$ ;  $\Delta(\text{G2})=E[\text{MP2}/6\text{-}311\text{G(3df,2p)}]-E[\text{MP2}/6\text{-}311\text{G(2df,p)}]-E[\text{MP2}/6\text{-}311\text{G(d,p)}]+E[\text{MP2}/6\text{-}311\text{G(d,p)}]$ ;  $\Delta(\text{HLC})=-0.19 \times \text{number of } \alpha \text{ valence electrons}-4.81 \times \text{number of } \beta \text{ valence electrons}$ ;  $\Delta(\text{ZPE})=\text{scaled zero-point vibrational energy}$ ; G2 energy =  $E[\text{MP4}/6\text{-}311\text{G(d,p)}] + \Delta(+)+\Delta(2\text{df})+\Delta(\text{QCI})+\Delta(\text{G2})+\Delta(\text{HLC})+\Delta(\text{ZPE})$ . <sup>b</sup> Energies in hartrees. <sup>c</sup> Energies in millihartrees.

TABLE 2: Calculated G2 Geometries and Frequencies (<sup>2</sup>A' Surface)<sup>a</sup>

	N <sub>2</sub> O	<i>cis</i> -HN-NO	<i>trans</i> -HN-NO	<i>cis</i> -HNNO	<i>trans</i> -HNNO	<i>cis-trans</i> ts	NN-OH	H-NNO
$r(\text{NN})$	1.171	1.902	1.911	1.186	1.196	1.186	1.157	1.129
$r(\text{NO})$	1.192	1.165	1.162	1.210	1.194	1.199	1.454	1.197
$r(\text{NH})$		1.041	1.036	1.038	1.023	1.049	1.203	1.423
$\angle\text{HNN}$		92.8	96.2	111.2	109.4	112.8	94.2	116.8
$\angle\text{NNO}$	180.0	119.5	125.0	143.1	137.9	177.0	96.4	168.6
$\omega_1$	2633	683 <i>i</i>	628 <i>i</i>	3600	3713	3376	2421 <i>i</i>	1286 <i>i</i>
$\omega_2$	1392	3212	3251	1640	1562	1976	1765	1830
$\omega_3$	689	1552	1658	1449	1422	1437	1280	1173
$\omega_4$		936	903	1167	1125	1143	805	759
$\omega_5$		317	348	353	322	905 <i>i</i>	471	375
$\omega_6$		253	361	941	1005	455	958	590

<sup>a</sup> Distances are in angstroms, angles in degrees, and frequencies in cm<sup>-1</sup>.



E (kcal)

Figure 1. Schematic potential energy surface for HNNO(<sup>2</sup>A').

therefore focused on characterizing the <sup>2</sup>A' surface. G2 calculations were utilized for most of the points on the surface; G2(Q) calculations were carried out for the transition states leading to N<sub>2</sub> + OH and N<sub>2</sub>O + H. The surface is shown schematically in Figure 1.

Geometries, frequencies, and energies for the stationary points located on the HNNO <sup>2</sup>A' potential energy surface can be found in Table 1–3. By combining these with other calculated G2 energies,<sup>25</sup> we can evaluate the energies for the stationary points on the HNNO <sup>2</sup>A' surface. These values can be found in Table 4. We find very good agreement between our results and experimental energetics for the overall exothermicities of the reaction.

Table 4 also lists the energetics calculated previously by other workers. The quality of previous calculations, as judged from the calculated exothermicities of the reaction, varies widely. The studies by Fueno et al.<sup>17</sup> and Harrison and MacIagan<sup>18</sup> utilized small basis sets and low levels of theory and only qualitatively reproduce the overall reaction exothermicities. There is uniformity in predicting that the transition state for the H + N<sub>2</sub>O channel is below that the N<sub>2</sub> + OH channel and general agreement

TABLE 3: Calculated G2Q Geometries and Frequencies (<sup>2</sup>A' Surface)<sup>a</sup>

	NN-OH	H-NNO
$r(\text{NN})$	1.215	1.152
$r(\text{NO})$	1.423	1.191
$r(\text{NH})$	1.245	1.547
$\angle\text{HNN}$	89.2	111.3
$\angle\text{NNO}$	96.1	167.4
$\omega_1$	2035 <i>i</i>	1401 <i>i</i>
$\omega_2$	2113	2173
$\omega_3$	1665	1288
$\omega_4$	1015	764
$\omega_5$	665	411
$\omega_6$	919	619

<sup>a</sup> Distances are in angstroms, angles in degrees, and frequencies in cm<sup>-1</sup>.

that both channels are thermodynamically open. We agree with Walch<sup>19</sup> and disagree with BAC-MP4<sup>4,16</sup> in predicting that the *trans*-HNNO isomer is more stable than the *cis* isomer. Unfortunately, there is no experimental data on HNNO with which to compare. We find that Walch's<sup>19</sup> energies are uniformly 2.5–4.5 kcal/mol higher than our G2(Q) results. Also, his energy for (N<sub>2</sub>O + H) relative to (NH + NO) is 3.0 kcal/mol higher than experiment. We believe that the source of this error lies in his calculation of  $E(\text{NH}+\text{NO})$  relative to the rest of the surface. An analogous situation exists in the case of the NH<sub>2</sub> + NO surface, where Walch has chosen to use the experimental heat of reaction to position of NH<sub>2</sub> + NO asymptote.<sup>30</sup> If we lower his reported energies by 3.0 kcal, the agreement between his work and the present work becomes quite good.

If we wish to describe the product branching fraction for the NH + NO reaction, we need to first consider the products of decomposition of the initially formed HNNO <sup>2</sup>A' adducts. The presence of a barrier in the channel forming *cis*-HNNO suggests that substantial amounts of *trans*-HNNO will be formed. Geometry rules out a 1,3-hydrogen transfer, followed by formation of OH and N<sub>2</sub>, for the *trans*-HNNO isomer. We also found that extension of the H–N bond leads to a lowering of the barrier to

TABLE 4: Energetics of HNNO Species<sup>a</sup>

<i>cis</i> -NNO	<i>trans</i> -HNNO	<i>cis</i> -HN-NO	<i>trans</i> -HN-NO	N <sub>2</sub> -OH ts	H-N <sub>2</sub> O ts	<i>cis-trans</i> ts	N <sub>2</sub> + OH	N <sub>2</sub> O + H	reference
-48.9	-56.0	2.9	0.1	-17.9 <sup>a</sup>	-25.3 <sup>a</sup>	-30.8	-96.9	-34.5	this work
-46.2	-51.5	6.3	3.2	-15.4	-21.4			-31.7	19
-51.9	-48.8			-22.4	-28.2		-99.1	-37.6	4
-30.9				3.0	-6.9		-90.8	-28.9	18 <sup>b</sup>
-29.5				3.6	-3.8		-91.6	-20.5	18 <sup>c</sup>
-60.3	-60.5			-9.0	-12.2	-28.7	-92.4	-30.9	17
-50.2	-47.3			-17.3	-27.7		-97.5	-31.4	16
							-97.6	-34.7	experiment <sup>d</sup>

<sup>a</sup> All energies in kcal/mol, relative to NH + NO, 0 K. <sup>b</sup> MP4/6-31G(d)//MP2/6-31G(d). <sup>c</sup> MP4/6-311G(d,p)//HF/6-311G(d,p). <sup>d</sup>  $\Delta H_f(\text{NH})$  from ref 1; all other  $\Delta H_f$  from JANAF tables.<sup>2</sup> <sup>e</sup> G2Q energy.

*cis-trans* isomerization, with the barrier disappearing before the N-H bond breaks. Therefore, we conclude that the only product arising from *trans*-HNNO is *cis*-HNNO. The barrier to *cis-trans* isomerization is low enough, relative to the barriers to dissociation, that we expect that the *cis* and *trans*-HNNO will be in equilibrium, and the product branching fraction will be described by the products of dissociation of *cis*-HNNO. On the basis of the relative locations of the transition states leading to N<sub>2</sub>O + H and N<sub>2</sub> + OH, we would predict that the reaction of NH + NO would form predominantly H + N<sub>2</sub>O products. Miller<sup>31</sup> has performed branching fraction calculations following the methodology used in his theoretical study of this reaction<sup>4</sup> and predicts, using G2 energies, a branching fraction of 0.89 for the N<sub>2</sub>O + H channel, in good agreement with our experimental result. He has additionally calculated a deuterium kinetic isotope effect of 0.96, which agrees with the lack of an experimentally observed isotope effect.

## Conclusion

We have directly measured the product branching fraction for the NH(D)(<sup>3</sup>Σ) + NO → products reaction at room temperature. We find that predominant, but not exclusive, production of N<sub>2</sub>O + H (D). We do not see an observable deuterium kinetic isotope effect in the product branching fractions. We have additionally used G2 and G2Q to characterize a number of stationary points on the HNNO(<sup>2</sup>A') surface. The exothermicities derived from this surface are in good agreement with experimental values. The shape of the G2(Q) surface is in good agreement with the recent CASSCF/CCI results of Walch,<sup>19</sup> especially if we lower his surface 3.3 kcal/mol, to bring the CASSCF/CCI exothermicity for the H + N<sub>2</sub>O channel into agreement with the experimental value.

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## References and Notes

- (1) Anderson, W. R. *J. Phys. Chem.* **1989**, *93*, 530.
- (2) Chase, M. W.; et al. *J. Phys. Chem. Ref. Data* **1985**, *14* (Suppl. No 1).
- (3) Miller, J. A.; Bowman, C. T. *Prog. Energy Combust. Sci.* **1989**, *15*, 287.
- (4) (a) Miller, J. A.; Melius, C. F. *Symp. (Int.) Combust. [Proc.]* **1993**, *24*, 719. (b) Melius, C. F. Private communication.
- (5) Gordon, S.; Mulac, W.; Nangian, P. *J. Phys. Chem.* **1971**, *75*, 2087.
- (6) Hansen, I.; Höinghaus, K.; Zetzsch, C.; Stuhl, F. *Chem. Phys. Lett.* **1976**, *47*, 370.
- (7) Cox, J. W.; Nelson, H. H.; McDonald, J. R. *Chem. Phys.* **1985**, *96*, 175.
- (8) Harrison, J. A.; Whyte, A. R.; Phillips, L. F. *Chem. Phys. Lett.* **1986**, *129*, 346.
- (9) Yamasaki, K.; Okada, S.; Koshi, M.; Matsui, H. *J. Chem. Phys.* **1991**, *95*, 5087.
- (10) Mertens, J. D.; Chang, A. Y.; Hanson, K.; Bowman, C. T. *Int. J. Chem. Kinet.* **1991**, *23*, 173.
- (11) Yokoyama, K.; Sakane, Y.; Fueno, T. *Bull. Chem. Soc. Jpn.* **1991**, *64*, 1738.
- (12) Dean, A. M.; Chou, M.-S.; Stern, D. *Int. J. Chem. Kinet.* **1984**, *16*, 633.
- (13) Vandooren, J.; Sarkisov, O. M.; Balakhnin, V. P.; van Tiggelen, P. *J. Chem. Phys. Lett.* **1991**, *184*, 294.
- (14) Sausa, R. C.; Anderson, W. R.; Dayton, D. C.; Faust, C. M.; Howard, S. L. *Combust. Flame*, in press.
- (15) Patel-Misra, D.; Dagdigian, P. J. *J. Phys. Chem.* **1992**, *96*, 3232.
- (16) Melius, C. F.; Binkley, J. S. *Symp. (Int.) Combust., [Proc.]* **1984**, *20*, 575.
- (17) Fueno, T.; Fukuda, M.; Yokoyama, K. *Chem. Phys.* **1988**, *124*, 265.
- (18) Harrison, J. A.; MacLagan, R. G. A. R. *J. Chem. Soc., Faraday Trans. 1990*, *86*, 3519.
- (19) Walch, S. P. *J. Chem. Phys.* **1993**, *98*, 1170.
- (20) Durant, J. L. *J. Phys. Chem.* **1991**, *95*, 10701.
- (21) Walther, C. D.; Wagner, H. G. *Ber. Bunsen-Ges. Phys. Chem.* **1983**, *87*, 403.
- (22) Hack, W.; Kurzke, H.; Wagner, H. G. *J. Chem. Soc., Faraday Trans. 2* **1985**, *81*, 949.
- (23) Frisch, M. J.; et al. *Gaussian 90*, Revision I; Gaussian, Inc.: Pittsburgh, PA, 1990.
- (24) Frisch, M. J.; et al. *Gaussian 92*, Revision B; Gaussian, Inc.: Pittsburgh, PA, 1992.
- (25) (a) Pople, J. A.; Head-Gordon, M.; Fox, D. J.; Raghavachari, K.; Curtiss, L. A. *J. Chem. Phys.* **1989**, *90*, 5622. (b) Curtiss, L. A.; Jones, C.; Trucks, G. W.; Raghavachari, K.; Pople, J. A. *J. Chem. Phys.* **1990**, *93*, 2537. (c) Curtiss, L. A.; Raghavachari, K.; Trucks, G. W.; Pople, J. A. *J. Chem. Phys.* **1991**, *94*, 7221.
- (26) Durant, J. L.; Rohlfing, C. M. *J. Chem. Phys.* **1993**, *98*, 8031.
- (27) Silver, J. A.; Kolb, C. E. *J. Phys. Chem.* **1982**, *86*, 3240.
- (28) Temps, F. Private communication, 1992.
- (29) Nelson, H. H.; McDonald, J. R.; Alexander, M. H. *J. Phys. Chem.* **1990**, *94*, 3291. Hack, W.; Wilms, A. *J. Phys. Chem.* **1989**, *93*, 3540. Bower, R. D.; Jacoby, M. T.; Blauer, J. A. *J. Chem. Phys.* **1987**, *86*, 1954. Freitag, F.; Rohrer, F.; Stuhl, F. *J. Phys. Chem.* **1989**, *93*, 3170.
- (30) Walch, S. P. *J. Chem. Phys.* **1993**, *99*, 5295.
- (31) Miller, J. A. Private communication.