

isoelectronic with CO. However, other considerations suggest that the chemical behavior of BF might be quite different from that of CO or N₂. In particular, electronic structure calculations indicate³⁰ that individual molecular orbitals of BF are quite asymmetric so that the small overall dipole moment could be misleading.

Nevertheless, Table V presents a comparison of isoelectronic rate coefficient data. Data are not available for N₂ + NO₂ or Cl₂. Indeed, the isoelectronic comparison is completely inappropriate for N₂ + Cl₂ because N₂Cl is not a stable molecule. The entries in the table indicate reasonable isoelectronic trends for the O, O₂, and Cl₂ reactions. Thus, BF and CO + Cl₂ both have negligible activation energy barriers, although the CO + Cl₂ reaction might have an endothermicity to overcome. The available

information³¹ on this reaction actually refers to the inverse reaction. Data for O atom and O₂ molecule reactions show reasonable trends, with N₂ the least reactive and BF the most reactive with either reagent. It is also noteworthy that preexponential terms are approximately the same within each of these isoelectronic series. However, this isoelectronic trends picture works less well for NO₂, with BF + NO₂ perhaps surprisingly fast in comparison with CO + NO₂. As a "pseudo-halogen", however, NO₂ might be expected to experience a stronger attractive interaction upon approach of BF as opposed to CO, and this may account for the much different reactivities.

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Comparative Pulsed CO₂ Laser Pyrolysis. The Decarbonylation Reaction of Propynal

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Comparative transversely excited atmospheric (TEA) CO₂ laser pyrolysis was applied to determine the activation energy of the reaction HC≡C—CHO → HC≡CH + CO. By use of three reference reactions, namely the H₂ elimination of cyclopentene and the dehydrohalogenation of vinyl chloride and vinyl bromide, the activation energy was found to be $E_a = 68 \pm 2$ kcal/mol. This result is in good agreement with $E_a = 69.9$ kcal/mol obtained from a previous ab initio SCF CI calculation of propynal's reaction path. The suitability of the method of comparative IR laser pyrolysis to determine activation energies is contrasted with its limitations when applied to evaluate frequency factors.

Introduction

Propynal, HC≡C—CHO, has shown to be a very suitable intermediate size molecule to explore the microscopic decay path of an electronically excited polyatomic molecule.¹⁻⁷ At moderate excess vibrational energies in the S₁ (nπ*) state, its chemical decay behavior implies the elimination reaction⁸



Since the spectral and dynamic data indicated the S₁ and T₁ potential surfaces not to be dissociative,^{5,7} it was inferred that the dissociation takes place on the potential surface of the electronic ground state. Motivated by this result, extensive ab initio cal-

culations at the CI level were performed to map the S₀ surface and to examine the chemical decay channels of this "intermediate case" test molecule.^{6,7} The lowest reaction path predicted by the calculations is decarbonylation via a concerted mechanism which involves the migration of the aldehyde H atom in the molecular plane. Taking into account the zero-point energies, the barrier height was calculated to be 69.9 kcal/mol.⁷ Thus, excitation into the S₁ vibrational ground state at 74.8 kcal/mol followed by internal conversion should lead to the chemical decay of propynal. This prediction has recently been confirmed by time-resolved experiments in which strong IR emission of nascent CO(v=1) molecules was observed when the UV laser was turned to the 0₀ absorption band.⁹ Based on this finding, an upper limit for the reaction barrier of 72 kcal/mol was estimated.⁹

In the present communication we report on comparative CO₂ laser pyrolyses of propynal. In order to determine the energy barrier of the lowest dissociation channel, sensitized IR laser heating was applied. In conjunction with the reference reactions H₂C=CHCl → HC≡CH + HCl, H₂C=CHBr → HC≡CH + HBr, and cyclopentene → cyclopentadiene + H₂, an activation energy of 68 ± 2 kcal/mol was found.

Experimental Section

In a first series of experiments propynal ($p = 1$ torr) was excited into its C—C stretching mode absorption ($\bar{\nu}_6 = 944 \text{ cm}^{-1}$)¹⁰ by TEA

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TABLE I: Partial and Total Pressures of the Reaction Mixture, Results of the Comparative Pyrolysis Experiments, Arrhenius Parameters of the Reference Reaction (E_{ref} , A_{ref}), Activation Energy, and Frequency Factor of the Probe Reaction (E_{pr} , A_{pr}) with Their Error Limits Obtained from the Experimental Results and the Reference Values

ref	P_{ref} torr	probe	P_{pr} torr	P_{Xe} torr	p_r^a torr	results from Figure 2		E_{ref} kcal/mol	$\log A_{ref}/s^{-1}$	E_{pr} kcal/mol	ΔE_{pr}		$\log A_{pr}/s^{-1}^c$	
						symbols	slope ^b				intercept ^b	slope		ref
cyclo- pentene	1	HCCCHO	1	13	20	7 circle	1.10 ± 0.04	-0.55 ± 0.15	61.0 ± 1.0 ^d	13.35 ^d	67.1	±2.4	±1.1	13.9 ± 0.25
						triangle								
						33 square								
H ₂ CCHCl	2	HCCCHO	1	17	25	12 square	1.01 ± 0.02	1.5 ± 0.15	69.3 ± 2.9 ^e	14.0 ± 0.2 ^e	69.7	±1.4	±2.9	14.8 ± 0.2
						17 triangle								
H ₂ CCHBr	1	HCCCHO	1	13	20	square	1.04 ± 0.03	2.0 ± 0.2	64.9 ± 2.4 ^f	13.3 ± 0.3 ^f	67.7	±1.3	±2.4	14.4 ± 0.35
cyclo- pentene	1	H ₂ CCHBr	1	13	20	circle	0.94 ± 0.03	1.05 ± 0.15	61.0 ± 1.0 ^d	13.35 ^d	64.9	±2.1	±1.1	13.3 ± 0.3

^aTotal pressure including 5 torr of SF₆. ^bWith standard deviation. ^c $\log A$ with $\tau = 10^{-5}$ s; error limits discussed in text. ^dFrom ref 16; $\Delta \log A$ (s⁻¹) = ± 0.2 assumed. ^eFrom ref 17. ^fCf. last line.

peratures, i.e., various laser pulse energies. Moreover, the frequency factors A_1 and A_2 are related by

$$\ln A_1 = B + \frac{E_1}{E_2} \ln A_2 - \left(1 - \frac{E_1}{E_2}\right) \ln \tau \quad (5)$$

Here the constant

$$B = \ln A_1 - \frac{E_1}{E_2} \ln A_2 + \left(1 - \frac{E_1}{E_2}\right) \ln \tau \quad (6)$$

denotes the intercept of the plot $\ln(k_1\tau)$ vs. $\ln(k_2\tau)$ on the $\ln(k_1\tau)$ axis. Thus, the Arrhenius parameters of the reaction of interest are obtained from a plot $\ln(k_1\tau)$ vs. $\ln(k_2\tau)$, provided E_1 is not too different from E_2 .^{12,14,15}

The results of four pyrolysis experiments are depicted in Figure 2. In each case we have selected a concerted, unimolecular elimination reaction as a reference. This reaction type is expected to reduce the possibility of side reactions. Table I summarizes experimental conditions, numerical results (slopes and intercepts) as derived by least-squares fits from Figure 2, the Arrhenius parameters of the reference reactions, and those found for the propynal and vinyl bromide decomposition.

The product yield of each educt pair investigated could be determined from the gas chromatogram in two ways. Either the remainder of educt 1 or that of educt 2 was used as a standard to calculate the absolute yield of C₂H₂. While both ways led to the same result in the case of the educt pair cyclopentene/vinyl bromide, small deviations were observed when propynal was one of the educts and high laser pulse energies were applied. Under these conditions additional products from propynal (<5% of propynal) introduced a small error. Therefore, we used only the educt compound admixed to propynal as an internal standard. The plots derived in this way (Figure 2) are linear throughout the region from low (i.e., $\ln(k\tau) \ll 0$) to high laser pulse energies ($\ln(k\tau) \sim 0$). It is noteworthy that here the product C₂H₂ is not subject to further reactions. In a separate experiment, a mixture of C₂H₂ with the sensitizer SF₆ was found to produce no detectable products even at the highest pulse energies employed in this work.

To assess the error limits of the measured activation energy of propynal, we considered the standard deviation of the slope in Figure 2 together with the uncertainty of the activation energy of the reference reaction. The pertinent data are listed in Table I. The errors of cyclopentene ($\Delta E = \pm 1$ kcal/mol)¹⁶ and vinyl chloride (± 2.9 kcal/mol)¹⁷ were taken from the literature. Since the published kinetic data of vinyl bromide¹⁸ are contradictory,¹⁹ we preferred to reinvestigate the dehydrobromination by laser pyrolysis using the cyclopentene decomposition as a reference

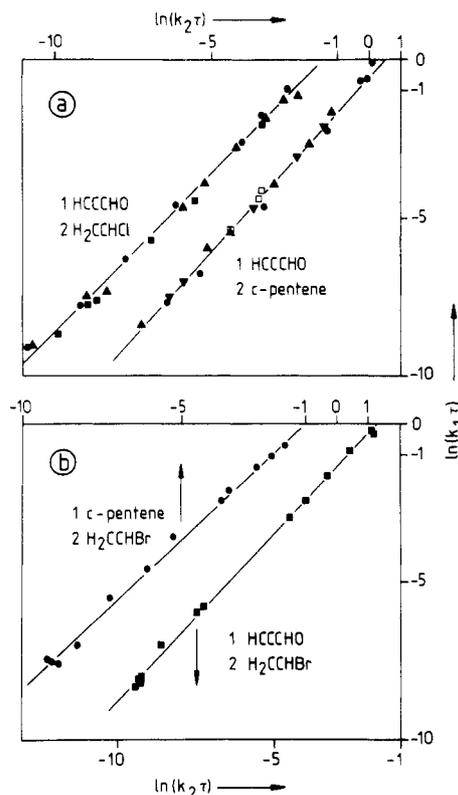


Figure 2. Plots $\ln(k_1\tau)$ vs. $\ln(k_2\tau)$ of the educt pairs propynal/vinyl chloride and propynal/cyclopentene (a) as well as cyclopentene/vinyl bromide and propynal/vinyl bromide (b). Note the two different abscissa scales for $\ln(k_2\tau)$ in (b). The composition of the reaction mixtures is given in Table I.

reaction. This result is also included in Table I.

The values obtained for the propynal reaction with cyclopentene, vinyl bromide, and vinyl chloride agree within the error limits ΔE (slope) and ΔE (reference). Taking into account both error sources, each propynal E value is estimated to be reliable within ± 3 kcal/mol. However, the precision of the method is far better. Using cyclopentene as a common standard for the propynal and vinyl bromide reactions and subsequently vinyl bromide for the propynal decomposition, we obtained for the activation energy of propynal $E = 67.1$ and 67.7 kcal/mol, respectively. According to these results an average value $E = 68 \pm 2$ kcal/mol is considered the most reliable.

The frequency factor of the probe reaction (propynal and vinyl bromide) can be determined from eq 5, if the slope $E_1/E_2 = d \ln(k_1\tau)/d \ln(k_2\tau)$, the intercept B , the frequency factor of the reference reaction, and the reaction time τ are known. The error limits $\Delta \log A$ (probe) with respect to B , A (reference), and τ were calculated with the slope given in Table I. The standard deviation ΔB and the values $\Delta \log A$ (reference) are listed in Table I. The

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reaction time which is not known should at least be 10^{-6} s and shorter than 10^{-3} s. The lower limit is given by the laser pulse duration ($\sim 1.5 \mu\text{s}$), and the upper limit corresponds to the time for heat flow from the irradiated region to the cold surroundings.²⁰ This cooling process was measured in a previous IR laser heating experiment³ where the temperature-dependent phosphorescence of propynal was utilized to monitor the transient temperature. The findings are in good agreement with those of ref 12 and 15 which were based on IR fluorescence measurements. To be on the safe side, we considered the 10^{-6} – 10^{-3} -s range as uncertainty and included it in $\Delta \log A$ (cf. Table I). However, the largest error in A of propynal is due to the uncertainty of the A value of the reference.

Discussion

The laser pyrolysis experiments performed with three different standards provided three sets of Arrhenius parameters for the decarbonylation reaction of propynal. The results are listed in Table I. The activation energy is found to be 67–70 kcal/mol while the values $\log A$ (s^{-1}) lie between 13.9 and 14.8.

To discuss these figures, we first consider the conditions prevailing in our laser pyrolysis experiments and the assumptions underlying the data analysis. In this context the most pertinent factors are a fast thermalization of the laser energy, homogeneity of the gas-phase chemistry, and reaction conditions close to the high-pressure limit.²¹

Possible surface reactions were suppressed by keeping the reaction zone, which was mainly determined by the laser beam dimensions, sufficiently far from the cell walls and by repeated polishing of the NaCl windows. A fast intermolecular randomization providing the same temperature for the reactive probe and reference molecules was favored by using a sensitizer, SF_6 . The direct photon absorption by probe or reference molecules was thus kept negligibly small. If one adapts the results of our previous study on sensitized TEA CO_2 laser heating of propynal³ to the present conditions (1 torr of propynal + standard + 5 torr of SF_6), the thermalization of the absorbed laser energy is expected to be complete within 1 μs . To further improve thermalization and to bring the reaction conditions closer to the high-pressure limit, we added Xe as buffer gas. Its effectiveness was examined with the reaction system propynal/cyclopentene. Due to the relatively large difference of the parameters E and A between propynal and cyclopentene, this system appeared most sensitive to deviations from thermal conditions. The results of the experiments with 13 and 33 torr of Xe agree within experimental error and parallel that without buffer gas as shown in Figure 2. According to this finding, a total pressure of 20–25 torr (including 13–18 torr of Xe) was regarded high enough to provide a fast equilibration of the absorbed laser energy and low enough to furnish a sufficient chemical conversion for ≤ 3 h of laser irradiation time.

For each pair of probe and reference reactions used in this work the activation energies are within 6 kcal/mol ($<10\%$ of E) and the $\log A$ values within 1.1 ($<10\%$ of $\log A$). Similar Arrhenius parameters of probe and reference reaction are very favorable, or even required, for the simple method of the data analysis given with eq 2–6.^{12,15} While method and data handling are appropriate for determination of the activation energies, they might fail in providing accurate $\log A$ values. This situation arises when a strong deviation from the high-pressure limit prevails. Under this condition the activation energy is only slightly decreased but the frequency factor is reduced by several orders of magnitude.²¹ If such a deficiency is considered, eq 2, 5, and 6 may be written as

$$\ln(k_1\tau) = \frac{E_1}{E_2} \ln(k_2\tau) + \ln A_1' - \frac{E_1}{E_2} \ln A_2' + \left(1 - \frac{E_1}{E_2}\right) \ln \tau \quad (7)$$

and

$$\ln A_1 = B + \frac{E_1}{E_2} \ln A_2 - \left(1 - \frac{E_1}{E_2}\right) \ln \tau + \frac{E_1}{E_2} \ln \alpha_2 - \ln \alpha_1 \quad (8)$$

with

$$B = \ln A_1' - \frac{E_1}{E_2} \ln A_2' + \left(1 - \frac{E_1}{E_2}\right) \ln \tau \quad (9)$$

where we have assumed that the ratio of the activation energies E_1/E_2 is the same in the falloff and the high-pressure regime. Deviations in the Arrhenius factor are approximated by the expression $A_1' = \alpha_1 A$ ($A_2' = \alpha_2 A_2$), A relating to the high-pressure value and A' to the effective one. The ratio A_1'/A_2' (or α_1/α_2) is assumed independent of the reaction temperature. Hence, a plot $\ln(k_1\tau)$ vs. $\ln(k_2\tau)$ of eq 7 furnishes E_1/E_2 as before, while the intercept B (eq 9) assumes a different physical meaning when compared with eq 6. Consequently, the determination of the activation energy is independent of this deficiency but not the value of A_1 as is evident from eq 8 and 5. Since it is difficult to establish α_1 and α_2 ,²¹ we are left with an A factor determined from eq 5 subject to a considerable uncertainty. For the special case $E_1 = E_2$ the magnitude of this uncertainty is, according to eq 8, given by $\ln(\alpha_2/\alpha_1)$, the contribution for the reaction in the falloff regime. Taking this into account, it might not surprise to find the $\log A$ values of propynal determined with eq 5 in a wide range and too high relative to the values reported for unimolecular elimination reactions.²²

Since the experimental value of the activation energy (68 kcal/mol) agrees within experimental error with the calculated energy barrier (69.9 kcal/mol), we preferred to determine the $\log A$ value with theoretical information. Based on the calculated energy, geometry, and vibrational frequencies of the transition state and on spectroscopic data of propynal in the electronic ground state, RRKM calculations with inclusion of tunneling corrections were previously performed.⁷ These rate constants $k(E)$ are now converted to $k(T)$ according to

$$k(T) = \int_0^\infty k(E) N(E) e^{-E/kT} dE / \int_0^\infty N(E) e^{-E/kT} dE \quad (10)$$

where $N(E)$ is the vibrational-state density at the energy E calculated with the same algorithm as used with the RRKM procedure. For temperatures of e.g. 1200 or 1800 K one finds

$$k(1200) = 2 \text{ s}^{-1} \quad k(1800) = 3 \times 10^4 \text{ s}^{-1}$$

With a barrier height of 68 kcal/mol instead of 69.9 kcal/mol, $k(T)$ increases by a factor of 3. Very similar values are obtained when the absolute rate theory is applied in conjunction with the transition-state geometry and its vibrations are taken from the ab initio calculation together with the spectroscopic data of the ground state.⁷ According to the rate constants given above, the preexponential factor between 800 and 1800 K is $\bar{A} = 1.2 \times 10^{13} \text{ s}^{-1}$, a magnitude consistent with the "tight" transition state shown in Figure 3 (*cis*-CTST₁) and with the "narrow maximum" of the potential energy along the minimum-energy path around the transition-state geometry.

Considering the strong discrepancy between the theoretical (10^{13} s^{-1}) and experimental A value ($10^{14.4 \pm 0.5} \text{ s}^{-1}$), it is instructive to estimate the effective reaction temperature during pyrolysis. The values of $\ln(k\tau)$ given in Figure 2 and an assumed $\tau = 10^{-5}$ s provide reaction rate constants at various laser pulse energies. When these rates are inserted into the Arrhenius equation of the standard or the probe reactions using the kinetic parameters from Table I, respectively, one finds $T = 1100$ – 1600 K. Furthermore, at a sample pressure of ~ 20 torr and with a gas kinetic collision rate, the ratio $k_{\text{coll}}/k(1600) \sim 10^8/10^5 \sim 10^3$ is obtained as the average value for these reactions. Compared with $k_{\text{coll}}/k(T)$ of similar reactions,²¹ this value is evidently not sufficient to bring the reaction conditions into the high-pressure limit. Since the

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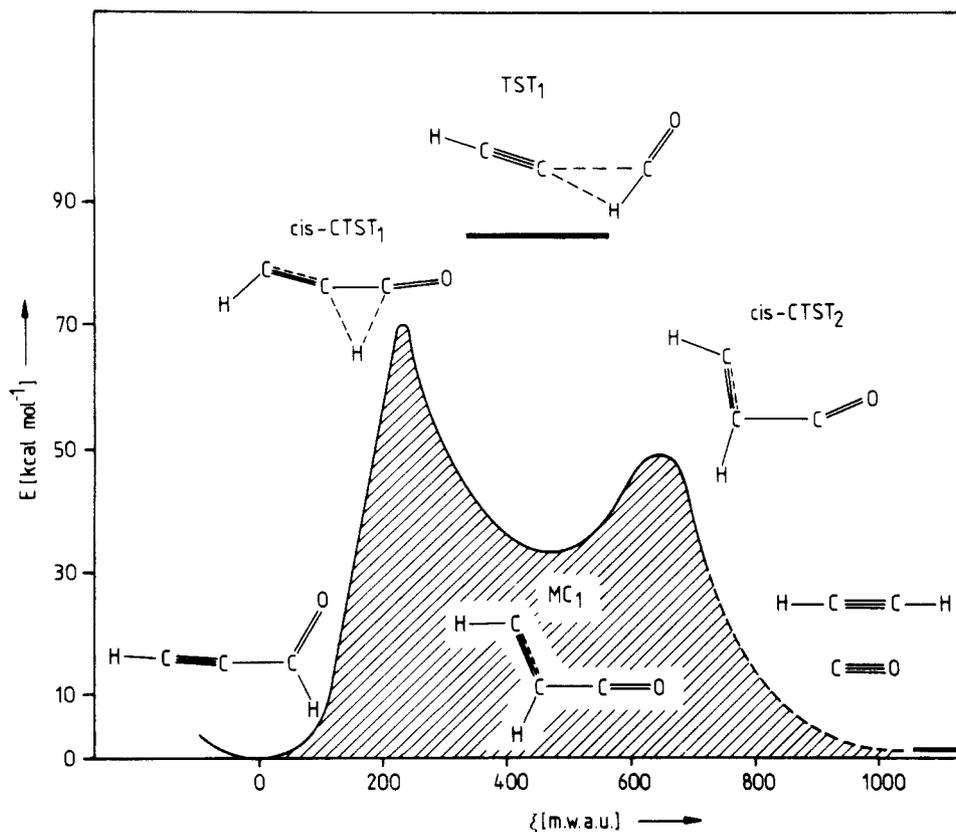


Figure 3. Calculated energy profile on the electronic ground-state surface of propynal along the lowest energy reaction coordinate ξ of the reaction $\text{HCCCHO} \rightarrow \text{CO} + \text{HCCH}$.⁷ The calculated geometries of the two transition states *cis*-CTST₁ and *cis*-CTST₂ and the metastable intermediate MC₁ are portrayed, together with the transition state TST₁ which belongs to the next higher reaction path with the same products. The energy scale refers to the vibrational ground state of propynal as origin; all values have been corrected for the zero-point energy.

available standard reactions show either lower activation energies or considerably higher A values than the propynal decarbonylation, a deviation from the high-pressure limit becomes severe.¹⁵ Unfortunately, an appropriate pressure increase quenches the laser-induced chemical conversion to such a degree that accurate product analysis is no longer feasible.

In the light of the present information as well as of previous, detailed dynamic data the pertinent features of the decarbonylation reaction of propynal are summarized. As displayed in Figure 3, the lowest channel involves two transition states with a maximum barrier height of 68 kcal/mol. Disregarding a slightly higher (~ 2 kcal/mol according to the calculation in ref 6) pathway via the *trans*-CTST₁ transition state, the next higher channel is predicted to possess only one transition state (TST₁) at about 83 kcal/mol (86 kcal/mol without correction for the zero-point energy).^{6,7} This reaction path bears some similarities with that of the formaldehyde decay $\text{H}_2\text{CO} \rightarrow \text{H}_2 + \text{CO}$, the activation barrier of which has been the subject of numerous calculations,²³ its most recent value being 80.9 kcal/mol.²⁴ The TST₁ reaction channel of propynal is characterized by fairly rigid $\text{HC}\equiv\text{C}$ and $\text{H}-\text{C}=\text{O}$ moieties while the lowest path proceeds via an energetically very favorable rearrangement (CTST₁) which involves most bond lengths and angles of the propynal frame. The subsequent potential barrier for CO elimination represented by CTST₂ is considerably lower and plays no role with respect to the reaction rate. The calculated decarbonylation path of propynal is a nearly thermoneutral reaction. Essentially the entire excitation energy goes into the fragments CO and $\text{HC}\equiv\text{CH}$. Moreover, the geometry of the second transition state, CTST₂, preceding the CO elimination, is strongly suggestive of nascent $\text{HC}\equiv\text{CH}$ with a high degree of CH deformational excitation. In accord with this prediction, IR

emission of the hot fragment CO⁹ as well as from the excited H-deformation vibration ($\nu_4 + \nu_5$) in acetylene²⁵ has been monitored after $\text{S}_1 \leftarrow \text{S}_0$ excitation.

Conclusion

The method of comparative pyrolysis with pulsed IR laser excitation has been applied to molecular elimination reactions with activation energies up to 70 kcal/mol. The high reaction temperature thus required could conveniently be achieved by sensitized laser heating without interference from heterogeneous surface reactions. Moreover, the efficient polymerization reaction of propynal was entirely suppressed. Using three different reference reactions, we determined the activation energy of the decarbonylation reaction of propynal to a precision of ± 2 kcal/mol. In spite of carefully chosen experimental conditions, the determination of the frequency factor failed. Comparing this result with the successful attempts performed on reactions having low activation energies, it is very likely that the high pyrolysis temperature prevents to provide the conditions (approach of the high-pressure limit) for an accurate evaluation of the A factor.

The good agreement between the calculated and the experimental activation energy lends further support to the predicted lowest energy path of the reaction $\text{HC}\equiv\text{C}-\text{CHO} \rightarrow \text{HC}\equiv\text{CH} + \text{CO}$. Owing to the low activation energy of 68 kcal/mol, electronic excitation into the S_1 state followed by internal conversion provides for propynal an efficient doorway to enter the reaction channel under collision-free conditions.

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Registry No. $\text{HC}\equiv\text{C}-\text{CHO}$, 624-67-9.

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