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# Shock-induced transformation of $\beta$ - $\text{Si}_3\text{N}_4$ to a high-pressure cubic-spinel phase

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$\beta$ - $\text{Si}_3\text{N}_4$  powders were shock compressed and quenched from 12 to 115 GPa.  $\beta$ - $\text{Si}_3\text{N}_4$  transforms to the spinel-type  $\text{Si}_3\text{N}_4$  ( $c$ - $\text{Si}_3\text{N}_4$ ) by a fast reconstructive process at pressures above about 20 GPa. The yield of  $c$ - $\text{Si}_3\text{N}_4$  recovered from 50 GPa and about 2400 K reaches about 80% and the grain sizes are about 10–50 nm. It is proposed that the fast transformation to  $c$ - $\text{Si}_3\text{N}_4$  occurs by rearrangement of nitrogen stacking layers, which initiates partial breakup of the  $\text{SiN}_4$  tetrahedra and formation of  $\text{SiN}_6$  octahedra at high density. Because of the advantages of massive production and the nanometer characteristics of shock-synthesized  $c$ - $\text{Si}_3\text{N}_4$ , it is possible to investigate the mechanical properties experimentally and to develop new industrial applications. © 2000 American Institute of Physics. [S0003-6951(00)02025-8]

Silicon nitride is an advanced ceramic and has the potential to be used in various applications due to good performance in mechanical, chemical, electronic, and thermal properties. This is due to highly covalent chemical bonding (70% covalence).<sup>1</sup> In addition to the two known polymorphs of  $\alpha$ - $\text{Si}_3\text{N}_4$  and  $\beta$ - $\text{Si}_3\text{N}_4$ , a third cubic form ( $c$ - $\text{Si}_3\text{N}_4$ ) with a spinel structure<sup>2</sup> has been found recently at pressures above 15 GPa and at temperatures over 2000 K, in several tens of nanograms, in a laser-heated diamond anvil cell for reaction times of about 1–10 min. This phase,  $c$ - $\text{Si}_3\text{N}_4$ , is expected to be a hard material because of its greater bulk modulus and shear modulus than those of  $\alpha$ - $\text{Si}_3\text{N}_4$  and  $\beta$ - $\text{Si}_3\text{N}_4$  (Ref. 2) and should also be a semiconductor due to its possible smaller band gap induced by the characteristic structure of octahedral silicon.<sup>3</sup> Because of the highly covalent nature of Si–N bonding, the high-pressure behavior of  $\alpha$ - and  $\beta$ - $\text{Si}_3\text{N}_4$  is of great interest over a wide range of pressure and temperature. Although the high-pressure behaviors of  $\text{Si}_3\text{N}_4$  have been studied at static high pressures<sup>4,5</sup> and dynamic high pressures,<sup>6,7</sup> only the spinel-type high-pressure phase has been observed experimentally as the high-pressure and high-temperature phase.<sup>2</sup> In this structure 2/3 of the silicon atoms are octahedrally coordinated and 1/3 of the silicon atoms remain tetrahedrally coordinated. It would be expected that all silicons have octahedral coordinations with increasing pressure, analogous to stishovite  $\text{SiO}_2$ .

For synthesis of the high-pressure and high-temperature form of  $\text{Si}_3\text{N}_4$ , we investigated the shock-induced transformation of  $\beta$ - $\text{Si}_3\text{N}_4$  powders using a propellant gun, because shock compression generates high-pressure and high-temperature simultaneously. We used two starting  $\beta$ - $\text{Si}_3\text{N}_4$  powders. The first one contained about 5 wt %  $\alpha$ - $\text{Si}_3\text{N}_4$  and less than 1.9 wt % oxygen impurity and the second was pure  $\beta$ - $\text{Si}_3\text{N}_4$  with 0.5 wt % oxygen and did not show the presence of  $\alpha$ - $\text{Si}_3\text{N}_4$  in x-ray analysis [Fig. 1(c)]. The grain sizes of both  $\beta$ - $\text{Si}_3\text{N}_4$  powders were in the submicron range. In the experiments, either pressed  $\beta$ - $\text{Si}_3\text{N}_4$  powder with 50%–60%

of the crystal density or pressed mixtures of  $\beta$ - $\text{Si}_3\text{N}_4$  and copper powders with 50%–80% of the theoretical density were encapsulated in containers of platinum or copper. Detailed descriptions of the shock recovery experiments have been published elsewhere.<sup>8</sup>

Metal flyer plates such as steel and platinum impacted samples in containers. The pressure generated in the sample was calculated by the impedance match method<sup>9</sup> using the measured flyer velocity.<sup>8</sup> The flyer velocities ranged between 1.5 and 2.1 km/s. The shock pressure varied from 12 to 115 GPa, depending on the Hugoniot of the flyer material<sup>10</sup> and impact velocity, and the compression lasted about a microsecond. The bulk temperature for each sample was computed by the thermodynamic method.<sup>11</sup> We used the Birch–Murnaghan equation<sup>12</sup> and thermodynamic parameters<sup>13</sup> for  $\beta$ - $\text{Si}_3\text{N}_4$  and Hugoniot for copper powders<sup>10</sup> to estimate the Hugoniot for mixtures. In our calculation the fusion energy of copper is taken into account but the transformation energy of  $\beta$ - $\text{Si}_3\text{N}_4$  to  $c$ - $\text{Si}_3\text{N}_4$  is not included due to lack of data.

After successful recovery, the container was cut open to take out the sample. To remove the copper matrix nitric acid was used at room temperature. The residue was collected after washing repeatedly with water. The color of the re-

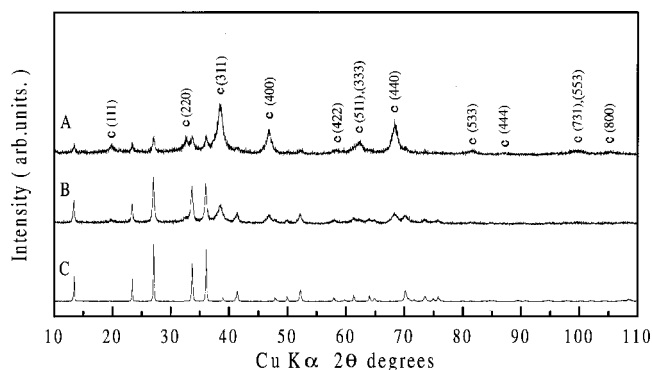


FIG. 1. XRD patterns of two recovered samples (a) from 49 GPa and 2400 K and (b) from 33 GPa and 1770 K and the starting material, pure  $\beta$ - $\text{Si}_3\text{N}_4$  (c). The peaks marked by  $c$  correspond to the cubic spinel-type  $\text{Si}_3\text{N}_4$ . The peak broadening is due to the fine grains.

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TABLE I. XRD data of cubic spinel-type  $c$ - $\text{Si}_3\text{N}_4$ .  $a = 0.774\,40 \pm 0.001\,06$  nm,  $d = 4.012$  g/cm<sup>3</sup>.

$hkl$	$2\theta_{\text{obs}}$ (deg)	$d_{\text{obs}}$ (nm)	$I/I_0$	$d_{\text{cal}}$ (nm)	$I_{\text{cal}}$
111	19.80	0.448 02	11	0.447 17	10
220	32.74	0.273 63	27	0.273 84	27
311	38.49	0.233 70	100	0.233 53	100
400	46.82	0.193 88	35	0.193 63	35
422	58.26	0.158 24	10	0.158 10	10
511, 333	62.33	0.148 85	33	0.149 06	34
440	68.39	0.137 06	56	0.136 92	57
620	77.88	0.122 56	4	0.122 46	4
533	81.56	0.117 93	11	0.118 12	11
444	87.01	0.111 89	5	0.111 79	5
642	97.42	0.102 52	5	0.103 50	5
731, 553	99.71	0.100 77	21	0.100 84	19
800	105.38	0.096 85	11	0.096 816	10
751, 555	119.20	0.089 31	16	0.089 435	15
840	125.74	0.086 55	12	0.086 595	12

sidual powder was gray, compared with the milky white of  $\beta$ - $\text{Si}_3\text{N}_4$ . We used x-ray diffraction (XRD) (Rigaku RINT 2200 V/pc) and analytical electron microscopy (TEM) (JEM 3000F) equipped with energy dispersive x-ray spectroscopy (EDX) to investigate the structure and chemistry of the shock-recovered sample.

Figure 1 illustrates a comparison of typical XRD patterns of the starting  $\beta$ - $\text{Si}_3\text{N}_4$  and two recovered samples from mixtures with copper powders shock compressed at 33 GPa and 1770 K [Fig. 1(b)] and 49 GPa and 2400 K [Fig. 1(a)], respectively. In the recovered samples the peaks of  $\beta$ - $\text{Si}_3\text{N}_4$  become weaker with increasing shock pressure and the newly appearing peaks correspond to the cubic spinel structure. The yield of  $c$ - $\text{Si}_3\text{N}_4$  reaches about 80%. Table I lists XRD data of  $c$ - $\text{Si}_3\text{N}_4$  produced from pure  $\beta$ - $\text{Si}_3\text{N}_4$ , together with the calculation results of the Rietveld method. The lattice parameter of  $c$ - $\text{Si}_3\text{N}_4$  is  $0.774\,40 \pm 0.001\,06$  nm and the calculated density is  $4.012$  g/cm<sup>3</sup>, which is 26% denser than that of  $\beta$ - $\text{Si}_3\text{N}_4$ . The XRD measured lattice parameter of  $c$ - $\text{Si}_3\text{N}_4$  is slightly (0.7%) smaller than the value  $(0.780 \pm 0.003$  nm)<sup>2</sup> measured by electron diffraction and the result  $(0.783\,67$  nm)<sup>3</sup> calculated by first principles. As shown in Fig. 2, EDX analyses of  $c$ - $\text{Si}_3\text{N}_4$  showed the presence of

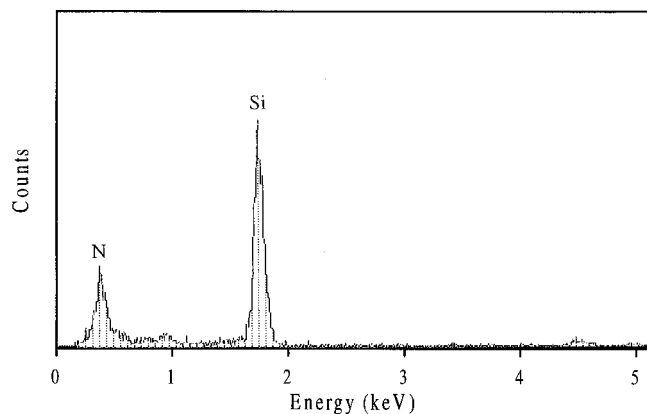


FIG. 2. EDX spectrum of shock-synthesized cubic spinel-type  $\text{Si}_3\text{N}_4$ . Only N and Si are observed and a very weak shoulder can be seen due to oxygen (0.5 keV) near the N peak. High backgrounds near 1 and 4.5 keV correspond to Cu and Ti impurities, respectively.

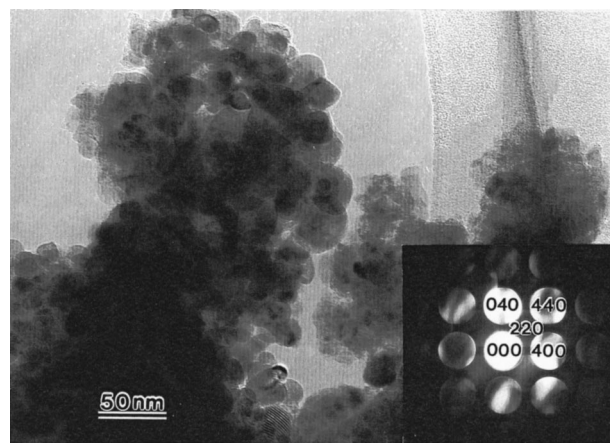


FIG. 3. Microphotographs of TEM and electron nanodiffraction of shock-synthesized, nanometer grains of spinel-type  $\text{Si}_3\text{N}_4$ . The observed  $d$  values are 0.273, 0.193, and 0.138 nm for diffractions (220), (400), and (440), respectively.

Si and N, with the same ratio as the starting  $\beta$ - $\text{Si}_3\text{N}_4$ , and detected a small amount of oxygen as a very weak shoulder on the N peak.<sup>14</sup> Therefore, the slight difference in the lattice parameter is due to the low precision of the measurement by electron diffraction.<sup>2</sup> The grain sizes of  $c$ - $\text{Si}_3\text{N}_4$  are 10–50 nm, which is about one tenth of the initial  $\beta$ - $\text{Si}_3\text{N}_4$ , as shown in Fig. 3 together with an electron nanodiffraction pattern indicating the presence of a cubic, nanometer-size single crystal.

The minimum shock pressure required for the formation of  $c$ - $\text{Si}_3\text{N}_4$  from  $\beta$ - $\text{Si}_3\text{N}_4$  is less than 20 GPa, but the yield from the run at 19 GPa and about 3000 K was the smallest that we could detect in the course of electron microscopy observation. The change to  $c$ - $\text{Si}_3\text{N}_4$  was not detected in samples subjected to shock temperatures of 1270 K at a pressure of 21 GPa, indicating the presence of a threshold minimum temperature. Too high a temperature causes decomposition and/or melting of  $\text{Si}_3\text{N}_4$ . In the case of samples of  $\beta$ - $\text{Si}_3\text{N}_4$  powders, the platinum container was broken and all the samples were exploded out above 80 GPa. This indicates that  $\text{Si}_3\text{N}_4$  decomposed to Si and  $\text{N}_2$  during shock compression, or during the adiabatic release process after shock compression, because of the higher temperatures. Even in the presence of copper in the mixture, the containers of highly porous samples were blown out when the calculated shock temperature was higher than 3000 K at 50 GPa. The  $\alpha$ - $\text{Si}_3\text{N}_4$  phase, coexisting in one of our starting  $\beta$ - $\text{Si}_3\text{N}_4$  powders, survived under shock conditions where most of the  $\beta$ - $\text{Si}_3\text{N}_4$  transformed to  $c$ - $\text{Si}_3\text{N}_4$ , suggesting that  $\alpha$ - $\text{Si}_3\text{N}_4$  has a higher activation barrier for phase transformation. The higher activation energy for the transformation of  $\alpha$ - $\text{Si}_3\text{N}_4$  to  $c$ - $\text{Si}_3\text{N}_4$  has also been demonstrated by the survival of  $\alpha$ - $\text{Si}_3\text{N}_4$  after being subjected to pressures of 37 GPa and temperatures of 3700 K in the diamond anvil cell.<sup>5</sup>

Compared with the duration of static high pressure work,<sup>2</sup> the present shock wave study has verified that the transformation to  $c$ - $\text{Si}_3\text{N}_4$  phase occurs within a microsecond. This suggests that a fast and probably diffusionless mechanism has played an important role. A fast phase transformation of  $\beta$ - $\text{Si}_3\text{N}_4$  to  $c$ - $\text{Si}_3\text{N}_4$ , through a solid–solid reaction, can be explained in terms of the anisotropic compression.

sion of  $\text{SiN}_4$  tetrahedra in  $\beta\text{-Si}_3\text{N}_4$ . In  $\beta\text{-Si}_3\text{N}_4$ ,<sup>1</sup> the structure units consist of slightly distorted tetrahedral  $\text{SiN}_4$  ( $sp^3$  hybridization of Si) and planar  $\text{NSi}_3$  ( $sp^3$  hybridization of N). Tetrahedra  $\text{SiN}_4$  are linked by sharing nitrogen corners so that each nitrogen is common to three tetrahedrons. The puckered rings of alternating Si and N atoms have a stacking sequence of ABAB... and form channels (diameter about 0.15 nm) along the  $c$  direction. The lattice parameters are  $a = 0.2911$  nm and  $c = 0.7595$  nm and the density is  $3.19$  g/cm<sup>3</sup>. On the other hand, in the spinel structure tetrahedral  $\text{SiN}_4$  and octahedral  $\text{SiN}_6$  are the basic units. The N atoms are cubic close packed in a stacking sequence ABCABC.... The  $\text{SiN}_4$  tetrahedron is linked to the  $\text{SiN}_6$  octahedra by sharing nitrogen corners, and each of the  $\text{SiN}_6$  octahedra is joined by sharing edges so that each nitrogen is common to one tetrahedron and two octahedrons.

The basic  $\text{SiN}_4$  tetrahedron unit in  $\beta\text{-Si}_3\text{N}_4$  is slightly distorted, and only the longest bond in the  $\text{SiN}_4$  tetrahedron may shorten significantly by compression. This distortion of the tetrahedral unit has been confirmed by neutron diffraction study at high pressure of  $\text{Si}_2\text{N}_2\text{O} (= \text{Si}_3\text{N}_4 \cdot \text{SiO}_2)$ <sup>16</sup> with a similar structure to  $\beta\text{-Si}_3\text{N}_4$ . The large open channel along the  $c$  axis is destroyed to make new bonds and the coordination number of 2/3 of the Si atoms increases from 4 to 6. As a result, this transformation involves a change of stacking sequence of N layers from ABAB... to ABCABC....

The shock compression technique has verified a fast transformation of  $\beta\text{-Si}_3\text{N}_4$  to  $c\text{-Si}_3\text{N}_4$ , a spinel-type high-pressure phase. This implies that the shock compression technique could be useful for synthesizing a high-pressure  $c\text{-C}_3\text{N}_4$  phase and other related materials such as a spinel solid solution between  $\text{Si}_3\text{N}_4$  and  $\text{C}_3\text{N}_4$ . As demonstrated by shock synthesis of diamond in industry, it has been thought that shock compression has several advantages over static compression for material synthesis. For example, higher pressures can be achieved more easily, simultaneously with high temperatures, and a large amount of samples are treated in each synthesis. Therefore, shock-synthesized  $c\text{-Si}_3\text{N}_4$  could become an important industrial material. The shock synthesis of high-pressure  $\text{Si}_3\text{N}_4$  has important implications for making hypothetical, hard materials of  $\beta\text{-C}_3\text{N}_4$  with the  $\beta\text{-Si}_3\text{N}_4$  structure<sup>17</sup> and cubic  $\text{C}_3\text{N}_4$  with the structure of high-pressure  $\text{Zn}_2\text{SiO}_4$  (willemite II).<sup>18</sup> The presence of stable  $\text{C}_3\text{N}_4$  forms has not been confirmed experimentally

yet, although the phases as a possible hard material have been predicted theoretically.<sup>17,18</sup> In this sense the transformation to the high-pressure cubic spinel form of  $\text{Si}_3\text{N}_4$  can be used as a guide for making  $\text{C}_3\text{N}_4$ .

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- <sup>14</sup>When a significant amount of oxygen is dissolved as  $\text{SiO}_2$  in  $c\text{-Si}_3\text{N}_4$ , the lattice parameter becomes smaller because the radius of N is larger than that of O. A similar observation can be seen in the Alon ( $\text{Al}_2\text{O}_3\text{-AlN}$ ) spinel solid solution.<sup>15</sup> Our sample contains only a little oxygen, which does not affect the lattice parameter significantly. The lattice parameter of  $c\text{-Si}_3\text{N}_4$  produced from  $\beta\text{-Si}_3\text{N}_4$  with 1.9 wt % oxygen is measured to be  $0.775\,05 \pm 0.000\,69$  nm.
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