

The Infrared Spectrum of C₂N₂ in the Region of the Bending Fundamental ν₅

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The high resolution infrared spectrum of C₂N₂ in the region 200–270 cm⁻¹ has been studied with a Fourier transform spectrometer. In addition to ν₅, the “hot” bands 2ν₅⁰ ← ν₅¹, 2ν₅² ← ν₅¹ and 3ν₅³ ← 2ν₅² have been analysed. Polynomial fits have been applied to obtain the vibrational term values and the rotational constants for the levels involved.

1. INTRODUCTION

Cyanogen (C₂N₂) is a stable, symmetric linear molecule. It has five fundamental vibrations. There are two bending modes, of which one is infrared active (ν₅ = 234 cm⁻¹) and the other Raman active (ν₄ = 503 cm⁻¹). The three stretching vibrations are the Raman active fundamentals ν₁ = 2331 cm⁻¹ and ν₂ = 846 cm⁻¹ and the infrared active fundamental ν₃ = 2157 cm⁻¹. High-resolution spectroscopic studies of this molecule are relatively few (1–7). The results concerning the lowest fundamental ν₅ are clearly incompatible. Accordingly, this investigation was undertaken to get accurate results for the ν₅ band and, at the same time, to obtain some information on the associated “hot” bands, too.

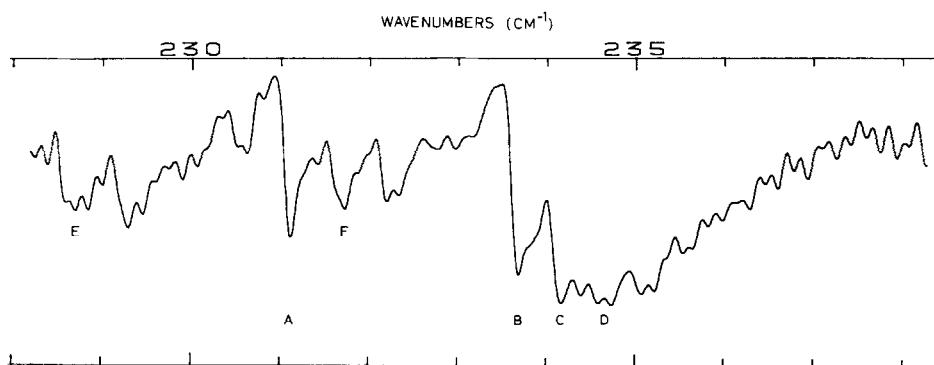


FIG. 1. The *Q*-branch region of ν₅ of C₂N₂ with low resolution. Absorption downward. The *Q* branches of the analyzed bands are marked as follows: (A) 0000⁰2⁰ ← 0000⁰1¹; (B) 0000⁰1¹ ← 0000⁰0⁰; (C) 0000⁰2² ← 0000⁰1¹; and (D) 0000⁰3³ ← 0000⁰2². The *Q*-branch *E* was tentatively attributed to (0000⁰3¹ ← 0000⁰2⁰) and the *Q* branch *F* to (0000⁰3¹ ← 0000⁰2²).

TABLE I
Analyzed Bands of C₂N₂

Transition	J _{MAX} P/R or Q	Measured lines	Standard deviation [10 ⁻⁴ cm ⁻¹]
0000 ⁰ 1 ^{1c} + 0000 ⁰ 0 ⁰	84/82	81	24
0000 ⁰ 1 ^{1d} + 0000 ⁰ 0 ⁰	75	26	22
0000 ⁰ 2 ⁰ + 0000 ⁰ 1 ^{1c}	62/59	52	19
0000 ⁰ 2 ^{2c} + 0000 ⁰ 1 ^{1c}	59/59	48	26
0000 ⁰ 2 ^{2d} + 0000 ⁰ 1 ^{1d}	61/61	44	28
0000 ⁰ 3 ^{3c} + 0000 ⁰ 2 ^{2c}	57/60	46	26
0000 ⁰ 3 ^{3d} + 0000 ⁰ 2 ^{2d}	59/57	48	30

2. EXPERIMENTAL DETAILS

The sample gas of cyanogen was prepared by allowing sodium cyanide to act on copper (II) sulfate. The gas pressure in the triple pass cell giving a path of 1

TABLE II
Observed Line Positions (cm⁻¹ in Vacuum) in the Fundamental ν₅ of C₂N₂

0000 ⁰ 1 ¹ + 0000 ⁰ 0 ⁰											
J	P _{cc} (J)	R _{cc} (J)	Q _{dc} (J)	J	P _{cc} (J)	R _{cc} (J)	Q _{dc} (J)	J	P _{cc} (J)	R _{cc} (J)	Q _{dc} (J)
5	232.1663	-	-	31	-	-	-	57	-	253.3905	-
6	-	-	-	32	-	244.5711	234.4116	58	216.9283	253.7567	-
7	231.5461	-	-	33	-	244.9143	-	59	216.6701	254.1148	-
8	-	-	-	34	223.5238	-	234.4994	60	216.4038	-	263.0733
9	230.9312	-	-	35	-	-	-	61	-	-	-
10	230.6261	237.2383	-	36	-	-	-	62	-	255.2132	-
11	230.3192	-	-	37	222.6739	246.2949	-	63	-	-	-
12	-	237.8922	-	38	222.3916	246.6472	-	64	-	-	-
13	-	238.2178	-	39	-	246.9868	-	65	215.1017	-	-
14	-	238.5468	-	40	-	-	234.7892	66	-	256.6867	-
15	229.1025	238.8714	-	41	-	-	-	67	-	-	-
16	-	239.2016	233.9040	42	221.2700	248.0410	234.8961	68	-	257.4189	236.7292
17	-	239.5312	-	43	-	248.3916	-	69	-	-	236.8135
18	228.2023	239.8580	-	44	-	248.7398	235.0098	70	213.8114	-	236.9056
19	227.9017	-	-	45	-	249.0934	235.0697	71	-	-	-
20	227.6051	240.5222	234.0010	46	-	-	235.1283	72	213.3098	-	237.0784
21	-	240.8568	-	47	-	-	-	73	-	-	-
22	-	241.1902	-	48	-	250.1598	-	74	212.8081	259.6631	-
23	226.7161	241.5252	234.0888	49	219.3413	-	-	75	-	260.0434	237.3604
24	226.4255	-	-	50	219.0727	-	235.3756	76	212.3015	260.4225	-
25	-	-	234.1510	51	-	251.2263	-	77	-	-	-
26	225.8361	242.5287	234.1825	52	218.5281	251.5874	235.5039	78	211.8113	261.1726	-
27	-	242.8715	234.2184	53	218.2637	251.9468	-	79	-	-	-
28	-	243.2051	234.2586	54	217.9958	252.3075	235.6427	80	211.3145	261.9366	-
29	224.9629	243.5497	234.2923	55	217.7294	252.6670	235.7085	82	210.8258	262.6972	-
30	224.6729	243.8901	-	56	-	253.0286	235.7824	84	210.3398	-	-

m was about 300 Pa. The measurements were performed on the Fourier spectrometer at the University of Oulu (8, 9). The resolution attained was better than 0.02 cm⁻¹. The wavenumber scale was calibrated with the aid of water vapor lines (10), which were simultaneously recorded.

3. RESULTS AND ANALYSIS

The term values of a linear molecule may be written as

$$T_v = G_v + B_v[J(J+1) - l^2] \pm (1/2)q_vJ(J+1) - D_v[J(J+1) - l^2]^2 \pm (1/2)\mu_vJ^2(J+1)^2, \quad (1)$$

where G_v is the vibrational term value, B_v is the rotational constant, D_v is the centrifugal distortion constant, and the l -doubling constants q_v and μ_v are the differences of the B and D constants for the component levels of the l -doublets. The upper sign in Eq. (1) applies to the d levels and the lower sign to the c levels. The polynomial expressions applied to the lines in the P and R branches as well as in the Q branch of rovibrational transitions can be derived from Eq. (1) (11).

TABLE III
Results from the Analysis of the Fundamental ν₅

	This work		
	P and R branch	Q branch	Literature
ν _{BC}	[cm ⁻¹] 233.7271(5)	233.7279(12)	
B'' = B _o	[cm ⁻¹] 0.157124(7)		0.157124(12) [1] 0.15712 [4] 0.157115(50) [6] 0.157135(10) [7]
B' - B''	[10 ⁻³ cm ⁻¹] 0.42855(46)	0.65205(109)	
B' = B ^c (ν ₅ ¹)	[cm ⁻¹] 0.157553(7)		
B' = B ^d (ν ₅ ¹)	[cm ⁻¹] 0.157776(8)*		
D'' = D _o	[10 ⁻⁸ cm ⁻¹] 2.63(8)		3.01(20) [1] 2.7(5) [6] 2.50(8) [7]
D' - D''	[10 ⁻⁸ cm ⁻¹] 0.085(7)	0.268(19)	
D' = D ^c (ν ₅ ¹)	[10 ⁻⁸ cm ⁻¹] 2.72(8)		
D' = D ^d (ν ₅ ¹)	[10 ⁻⁸ cm ⁻¹] 2.90(9)**		
q ₅	[10 ⁻⁵ cm ⁻¹] 22.59(9)		22.43(18) [5]
μ ₅	[10 ⁻⁹ cm ⁻¹] -2.2(3)		

* Calculated with the aid of B_o = 0.157124(7) cm⁻¹.

** Calculated with the aid of D_o = 2.63(8) · 10⁻⁸ cm⁻¹.

The central part of the spectrum with low resolution is shown in Fig. 1. This spectrum has been transformed from a small part of the interferogram points only to visualize better the Q branches.

Table I presents an overall summary of the different bands analyzed. For each of these, the observed maximum J value, the number of measured lines, and the standard deviation of the data for the whole band are given.

TABLE IV
Observed Line Positions of the "Hot" Bands

3.1. The Fundamental ν₅

The observed line positions of the fundamental ν₅ are listed in Table II. The molecular constants from unconstrained polynomial fits are given in Table III. The error limits quoted in this table as well as in all the others are standard deviations in the units of the last digits given. The *l*-doubling constants q₅ and μ₅ were determined from the relation

$$F_{5d}(J) - F_{5c}(J) = q_5 J(J+1) + \mu_5 J^2(J+1)^2. \quad (2)$$

The values of *F*(*J*) were calculated from the line positions given in Table II, and from the results for *B*₀ and *D*₀ obtained in this work. The results have been compared with selected literature data. The ground-state rotational constants *B*₀ and *D*₀ and the *l*-doubling constant *q*₅ are in good agreement with the literature values. The result for ν₅ is greater than the calculated literature values 233.0722(15) cm⁻¹ (5) and 233.15(15) cm⁻¹ (1, 4). The calculated value 233.7 cm⁻¹ reported by Jones (12), however, is consistent with the present result. Wang *et al.* (7) observed (1000⁰0⁰ ← 0000⁰0⁰) at 2330.512(2) cm⁻¹ and Picard (6) measured 2096.765(1) cm⁻¹ for the difference transition (1000⁰0⁰ ← 0000⁰1¹). These values yield ν₅ = 233.747(3) cm⁻¹, which is reasonably near to the present result, although the difference is larger than the error limits would allow.

3.2. The "Hot" Bands

All the "hot" bands analyzed are listed in Table I. The observed line positions are presented in Table IV. Table V shows the results obtained from the fits where the *B*" and *D*" constants were constrained to the values given by the present findings. In the error limits for the *B* and *D* of the upper state, the uncertainty of the constrained value has been added to the standard error (one standard deviation) given by the fit. The rotational constants *B* of the levels 0000⁰ν₅' (with *l* = ν₅ and ν₅ = 0, 1, 2, 3) grow regularly as ν₅ increases. The values were fitted according to the expression

TABLE V
Results from the "Hot" Bands of C₂N₂

Transition	ν _{BC} [cm ⁻¹]	B' [cm ⁻¹]	B" [cm ⁻¹]	B' - B" [10 ⁻³ cm ⁻¹]	D' [10 ⁻⁸ cm ⁻¹]	D" [10 ⁻⁸ cm ⁻¹]	D' - D" [10 ⁻⁸ cm ⁻¹]
0000 ⁰ 2 ⁰ ← 0000 ⁰ 1 ¹ c	231.1524(7)	0.158198(8)	0.157553	0.6455(9)	4.27(10)	2.72	1.555(23)
0000 ⁰ 2 ² c ← 0000 ⁰ 1 ¹ c	234.1926(8)	0.158201(8)	0.157553	0.6480(14)	1.75(12)	2.72	-0.967(40)
0000 ⁰ 2 ² d ← 0000 ⁰ 1 ¹ d	234.1923(11)	0.158205(10)	0.157776	0.4288(15)	3.12(13)	2.90	0.220(39)
0000 ⁰ 3 ³ c ← 0000 ⁰ 2 ² c	234.6628(9)	0.158721(10)	0.158201	0.5199(14)	2.48(16)	1.75	0.728(40)
0000 ⁰ 3 ³ d ← 0000 ⁰ 2 ² d	234.6626(10)	0.158731(11)	0.158205	0.5256(16)	2.32(18)	3.12	-0.801(49)

B" and D" constrained (this work).

TABLE VI
Some Vibrational Term Values (cm^{-1}) of C_2N_2

Level	$G_0(v) [\text{cm}^{-1}]$
0000^01^1	233.8849(7)
0000^02^0	464.8796(13)
0000^02^2	468.5526(15)
0000^03^3	704.0110(23)

$$B(0000^0v'_5) = B_0 - \alpha_5^0 v_5 + (\gamma_{55}^0 + \gamma_{ll}) v_5^2. \quad (3)$$

The fit resulted in $\alpha_5^0 = -0.548(5) \times 10^{-3} \text{ cm}^{-1}$ and $\gamma_{55}^0 + \gamma_{ll} = -0.45(15) \times 10^{-5} \text{ cm}^{-1}$.

The levels $2\nu_5^0$ and $2\nu_5^{2c}$ perturb each other through *l*-type resonance, while the $2\nu_5^{2d}$ state is unperturbed (13). The effect of the resonance appears in the ($D' - D''$) values obtained from the fits in the $2\nu_5 \leftarrow \nu_5$ bands. The values of ($D' - D''$) in the resonance diad are roughly equal in absolute value, but have opposite signs, as shown in Table V.

The absorption at 228.7 cm^{-1} , (E) in Fig. 1, was tentatively attributed to $(0000^03^1 \leftarrow 0000^02^2)$ and, respectively, the absorption (F) at 231.7 cm^{-1} to $(0000^03^1 \leftarrow 0000^02^0)$. Some rotational lines could be assigned in accordance with these assumptions, but the results remained so dubious that they will not be reported.

3.3. Vibrational Term Values

The determined term values, as listed in Table VI, were obtained by combining the band centers, B values and q values and by employing the following expression (14) for a band center:

$$\nu_{BC} = G_0(v') - G_0(v'') - B'l'^2 + B''l''^2 \pm (1/2)(q'l'^2 - q''l''^2), \quad (4)$$

where the different ν_{BC} values were weighted with the inverses of their estimated errors squared. The uncertainties in Table VI are the sums of the uncertainties of the various components included in the determinations.

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