An Electron Spin Resonance Investigation of β-Substituted 1,1-Diphenylethyl Radicals

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A large number of β -substituted 1,1-diphenylethyl derivatives Ph₂CCH₂MR_n have been generated by addition of \cdot MR_n radicals to 1,1-diphenylethylene. The magnitude of the β -proton hyperfine splitting constants together with the observation of linewidth alternation effects on the CH₂ triplet, indicate that the preferred conformation of these radicals is close even though not coincident with that placing the MR_n group in an eclipsed position with respect to the $2p_z$ orbital on C_x. The slight departure of these adducts from an eclipsed geometry has been interpreted as due to the presence on C_x of the two phenyl rings which are somewhat twisted with respect to each other thus producing an asymmetric potential around the radical centre. Although all the 1,1-diphenylethyl derivatives which have been investigated adopt the same conformation, the β -proton splitting cover the range from 6 to 11 G. The two common interpretations invoked to explain the effect of β -substituents on this splitting are discussed.

The conformational preference of β -substituted ethyl radicals has been largely studied by e.s.r. spectroscopy. Simple ethyls of general structure $\dot{C}H_2CH_2MR_n$, have been found to adopt the conformation (I), in which the β -substituent MR_n eclipses the $2p_z$ orbital on C_{α} , when M is an atom from rows 2—4 of the Periodic Table,¹⁻⁸ or the staggered conformation (II) when M is from row 1.^{9,10} Preference for the eclipsed conformation (I) is also exhibited by first row elements when the protons on the α -carbon are replaced by bulky groups, as for instance in the β -substituted 1,1-dimethyl-¹⁰ and 1,1-di-tbutyl-ethyls.¹¹

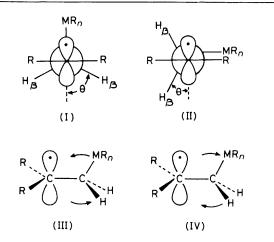
The minimum energy conformation of these radicals can be established from the isotropic hyperfine splitting at the β -protons, $a_{H_{B}}$, by means of the relation (1)¹² where θ is the

$$a_{\rm H_B} = A + B \left< \cos^2 \theta \right> \tag{1}$$

dihedral angle between the symmetry axis of the $2p_z$ orbital on C_{α} and the C_{β} -H_{β} bond. In equation (1) *A* is usually neglected ¹³ and *B* is assumed to be twice the hyperfine splitting at the methyl protons in the related R₂CCH₃ radical, since for a freely rotating methyl group $\langle \cos^2\theta \rangle$ is 0.5. According to equation (1) the β -proton splittings should be *B*/4 in conformation (I) and 3*B*/4 in conformation (II). However, only occasionally $a_{H_{\beta}}$ has been found to correspond to one of these two values. Usually the measured splittings are intermediate between one of them and the *B*/2 value expected for free rotation of the CH₂MR_n group because of torsional oscillations around the minimum energy conformation.

Generally, $a_{H\beta}$ exhibits also a temperature dependence [positive or negative when the preferred conformation is (I) or (II), respectively], since temperature variations induce a change in the populations of the torsional levels determining the ensemble average $\langle \cos^2\theta \rangle$

Interestingly there are a few β -substituted alkyl radicals adopting the eclipsed conformation (I) where the β -proton splitting is lower than *B*/4. Examples include ethyls bearing Br,¹⁴ Cl,^{2,11,15,16} and SR ^{1,17,18} β -substituents. A second remarkable class of radicals is represented by alkyls where the leading atom of the MR_n β -substituent is a Group IVB metal (Si, Ge, or Sn).^{1,11,19} In these radicals $a_{H_{\beta}}$ is invariably larger than *B*/4 even though the absence of temperature dependence of this splitting indicates that the MR_n group is fairly rigidly held in conformation (I).¹¹



The abnormally low values of $a_{H_{\beta}}$ found with bromine, chlorine, or thio-substituents have been interpreted in terms of geometry deformations about the β -carbon due to an attractive interaction between these electronegative groups and the $2p_z$ orbital on the α -carbon.^{1,16,20} This causes the substituent to move towards the radical centre and the β -protons away from it, as exemplified in structure (III), with the result of decreasing the overlap between the C-H bonds and the singly occupied orbital and, therefore, the value of $a_{H_{\beta}}$. This effect is commonly referred as asymmetric bridging.

On the other hand no particular attention has been paid to the remarkably large values of the β -proton coupling in alkyls containing Group IVB substituents. However, there is at least one paper where larger than expected $a_{H\beta}$ splitting observed in trialkylsilyl adducts of 1,1-di-t-butylethylene are explained in terms of structural distortions of the radical due to steric crowding as indicated in (IV).¹⁹

As an alternative to explanations based on the deformation of the molecular skeleton, it has also been suggested that variations of $a_{H\beta}$ may be caused by the different electronegativities of the β -substituents simply by increasing or decreasing the electron-releasing power of the C_{β} -H_{β} bonds toward the singly occupied orbital.^{4,21} In practice this means that the assumption, implicit in equation (1), that a unique value of *B* will describe the conformational dependence of $a_{H\beta}$ in any radical is not strictly correct. Computational support to the latter interpretation has been recently given for the case of alkyls containing silyl and thiyl substituents.¹⁸

A further complication is that these radicals may be nonplanar at the radical carbon centre, especially when they bear alkyl substituents at C_{α} (see for instance the t-butyl radical).²² In consequence, the presence of a more or less electronegative substituent on C_{β} may also induce a change in the degree of pyramidalization of the radical site and therefore in the magnitude of a_{Ha} .

We report here the e.s.r. spectral parameters of a large number of 1,1-diphenylethyls obtained by addition of $\cdot MR_n$ radicals to Ph₂C=CH₂. Since all these adducts adopt a conformation very close to the eclipsed (I) because of the crowding around the α -carbon, they are suitable probes to test the two interpretations discussed above. An advantage of these radicals is that they are likely to be more planar at the radical centre than 1,1-dialkylethyls, as the result of the extended delocalization of the unpaired electron in the two phenyl rings.

Experimental

1,1-Diphenylethylene was a commercial product which was distilled at reduced pressure prior to use. Ethyl t-butyl peroxide ²³ and diacetyl peroxide ²⁴ were prepared according to established methods. 1,1-Diphenyl-2-chloro- and 1,1-diphenyl-2-bromo-ethane were obtained by reaction of Ph₂CHCH₂-OSO₂C₆H₄CH₃ ²⁵ with LiCl or LiBr in hexamethylphosphoramide (HMPA) following the procedure given by Mosher *et al.* for the preparation of neopentyl chloride and neopentyl bromide.²⁶

1,1-Diphenyl-2-chloroethane. A solution of the tosylate (3.52 g) and LiCl (0.34 g) in HMPA (13 ml) and water (10 drops) was heated at 90 °C for 8 h. The mixture was poured into cold water and extracted with chloroform. The organic phase was dried, the solvent removed, and the residue was eluted with light petroleum-diethyl ether (95:5). The product was obtained as an oil (1.2 g) which was vacuum distilled, b.p. 126—128 °C at 2 mmHg (lit.,²⁷ 112 °C at 0.05 mmHg). 1,1-Diphenyl-2-bromoethane. Using the same procedure

1,1-Diphenyl-2-bromoethane. Using the same procedure described above for the chloro-derivative, the title compound (0.6 g) was obtained from the tosylate (3.52 g) and LiBr (0.82 g) in HMPA (10 ml), m.p. 61-63 °C (lit.,²⁸ 63 °C).

Other materials were commercially available and used without further purification.

The radical adducts were generated in degassed solutions of t-butylbenzene by reaction of $Ph_2C=CH_2$ (*ca.* 10%) with $\cdot MR_n$ radicals produced photochemically or thermally. Methyl-cyclohexane or n-pentane were used occasionally as solvents. The radical concentration was found to depend on temperature, stronger e.s.r. signals being generally obtained at higher temperature (*ca.* 370 K).

The \cdot MR_n radicals were generated by the following methods. (i) Hydrogen abstraction from the appropriate precursor with t-butoxyl radicals [reactions (2) and (3)]. This method was

$$(Bu^{t}O)_{2} \xrightarrow{hv} 2Bu^{t}O$$
 (2)

$$Bu^{t}O \cdot + HMR_{n} \longrightarrow Bu^{t}OH + \cdot MR_{n}$$
 (3)

used for $\cdot MR_n = \cdot SiPh_3$, $\cdot GePh_3$, $\cdot SnPh_3$, $\cdot Ccl_3$, $\cdot CHCl_2$, 1adamantyl, $\cdot Cl$, $\cdot COMe$, $\cdot COPh$, $\cdot COC_6F_5$; all gave adducts with Ph₂C=CH₂.

(ii) Photolytic cleavage (4) of the M-M bond of peroxides, disulphides, and diselenides. Adducts were observed for

$$\mathbf{R}_{n}\mathbf{M}\mathbf{M}\mathbf{R'}_{n} \xrightarrow{n^{v}} \mathbf{R}_{n}\mathbf{M}^{\cdot} + \mathbf{R'}_{n}\mathbf{M}^{\cdot}$$
(4)

 $\cdot MR_n = \cdot SMe$, $\cdot SEt$, $\cdot SPh$, $\cdot OEt$, and $\cdot OCOPh$, while no radicals were detected in the temperature range 270–390 K

for $\cdot MR_n = Bu'O \cdot$ and PhSe. The O-O cleavage (5) of diacetyl peroxide was immediately followed by decarboxylation²⁹ and we could observe only the adduct of 1,1-diphenylethylene with methyl radicals.

$$MeC(O)OOC(O)Me \xrightarrow{hv} 2MeC(O)O \cdot \xrightarrow{} 2CO_2 + 2Me \cdot (5)$$

(iii) Photolysis of di-t-butyl peroxide (DTBP) in the presence of organometallic compounds. Through reaction (6) we

$$Bu^{t}O^{\cdot} + X(MR_{n})_{3} \longrightarrow Bu^{t}O^{\dot{X}}(MR_{n})_{3} \longrightarrow Bu^{t}OX(MR_{n})_{2} + \cdot MR_{n} \quad (6)$$

could obtain the Ph radical from triphenylarsine ³⁰ and triphenylborine,¹¹ and Buⁿ and Me₂N radicals from tri-nbutylphosphine and tris(dimethylamino)phosphine ³¹ respectively, all giving adducts with Ph₂C=CH₂. The olefin adduct was also observed with the Bu¹ radical obtained from triphenylphosphine or triethyl phosphite [reaction (7)],³² while

$$Bu^{t}O + PR_{3} \longrightarrow Bu^{t}O\dot{P}R_{3} \longrightarrow O=PR_{3} + Bu^{t}$$
 (7)

no addition of the intermediates phosphoranyl was detected. On the other hand, the diethoxyphosphonyl radical obtained from tetramethyl pyrophosphite ³³ gave an adduct with $Ph_2C=CH_2$ [reaction (8)].

$$Bu^{t}O^{\cdot} + (EtO)_{2}POP(OEt)_{2} \longrightarrow O^{=}\dot{P}(OEt)_{2} + Bu^{t}OP(OEt)_{2} \quad (8)$$

(iv) Iodine or bromine abstraction with $Bu_{3}^{n}Sn$ radicals [reactions (9) and (10)]. This method was applied to generate

$$(Bu^{n}_{3}Sn)_{2} \xrightarrow{hv} 2 Bu^{n}_{3}Sn$$
 (9)

$$Bu^{n}_{3}Sn + XMR_{n} \rightarrow MR_{n} + Bu^{n}_{3}SnX \quad (10)$$

• CF_3 and PhCO radicals from CF_3I and PhCOBr respectively.¹¹ The olefin adduct with trifluoromethyl gave strong e.s.r. signals even at room temperature.

(v) Photolytic or thermal decomposition (11) of $R_nMN=NM'R_n'$.

$$\mathbf{R}_{n}\mathbf{M}\mathbf{N}=\mathbf{N}\mathbf{M}'\mathbf{R}'_{n}\longrightarrow\mathbf{N}_{2}+\mathbf{R}_{n}\mathbf{M}\cdot\mathbf{R}_{n}'\mathbf{M}'\cdot\quad(11)$$

Room temperature photolysis of solution of $Ph_2C=CH_2$ and azobisisobutyronitrile resulted in the appearance of the e.s.r. spectrum of the CMe_2CN radical. Above 330—340 K this spectrum was replaced by that one from the adduct with 1,1-diphenylethylene.

Decomposition of PhNHN=NPh gave only the adduct with Ph· and no evidence for the addition of the aniline radical PhNH was obtained. When photolysing solutions containing Ph(Me)NN=NN(Me)Ph we could observe weak signals due to the phenyl methyl nitroxide PhN(\ddot{O})Me, but no adduct with Ph₂C=CH₂.

Also other attempts to generate adducts with aminyl radicals, other than Me₂N·, were unsuccesful. Photolysis of Ph₂C= CH₂ and Ph₂NH in the presence of di-t-butyl peroxide (DTBP) yields the e.s.r. spectrum of diphenyl nitroxide, while without DTBP a strong and well resolved spectrum of diphenylaminyl Ph₂N· (a_N 8.76, a_o 3.62, a_m 1.50, a_p 4.25 G, g 2.0032) was observed.³⁴ The generation of aminyl radicals by photochemical reaction of secondary amines with olefins in nonpolar solvents has been already observed and interpreted in terms of the formation of an intermediate exciplex followed by hydrogen transfer.³⁵

Room temperature hyperfine splitting constants ($G = 10^{-10}$	$^{-4}$ T) and g fa	actors for the radical a	adducts Ph ₂ CCH ₂ MR	, of 1,1-diphenylethylene
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MR _n	$a_{\mathrm{H}}^{\mathrm{Ph}}(o,m,p)$	a _{HB}	d <i>a</i> нβ/d <i>T</i> (mG K ⁻¹)	aother	g
CH ₃	3,11, 1,26, 3,38	9.04	5		2.002
Bu ⁿ ^a	3.10, 1.26, 3.36	9.58	b	$a(\gamma-CH_2) 0.60$	2.002
				$a(\delta-CH_2) 0.15$	
Bu'	3.08, 1.27, 3.35	9.41	1.5	$a(\delta - CH_3) 0.14$	2.002
1-Adamantyl	3.10, 1.28, 3.37	9.44	1		2.002
Ph	3.10, 1.27, 3.39	9.05	5.5	a(2H) 0.16	2.002
CHCl ₂ ^a	3.09, 1.27, 3.37	9.49	Ь		2.002
CMe ₂ CN ^c	3.06, 1.25, 3.32	9.05	b	a(δ-CH ₃) 0.14	2.002
CCl ₃	3.07, 1.26, 3.35	8.50	1		2.003
CF ₃	3.16, 1.28, 3.44	8.90	1 2 b	a(3F) 0.62	2.002
C(O)Me ^c	3.14, 1.28, 3.42	9.67	Ь	a(3H) 0.13	2.002
C(O)Ph	3.11, 1.26, 3.38	9.34	6		2.002
$C(O)C_6F_5^{a}$	3.03, 1.24, 3.31	9.54	ь 5		2.002
NMe ₂ ^a	3.08, 1.26, 3.36	9.56	5	a(N) 4.32	2.002
OEt	3.12, 1.27, 3.40	9.58	2	$a(CH_2) 0.31$	2.002
				a(CH ₃) 0.22	
OC(O)Ph ^c	3.11, 1.28, 3.41	8.50	5		2.002
SiPh ₃	3.00, 1.22, 3.28	11.08	1		2.002
GePh ₃	2.97, 1.21, 3.26	10.39	Ь		2.002
SnPh ₄	2.92, 1.18, 3.18	10.12	Ь		2.001
$P(O)(OEt)_2^d$	3.11, 1.28, 3.39	9.91	1	a(P) 66.40	2.002
SMe	3.06, 1.27, 3.35	7.65	Ь	a(CH ₃) 0.43	2.003
SEt	3.04, 1.26, 3.32	7.57	4	$a(CH_2) 0.50$	2.003
				a(CH ₃) 0.20	
SPh	3.04, 1.26, 3.33	7.80	2		2.003
Cl ^a	2.98, 1.25, 3.25	6.00	Ь	a(³⁵ Cl) 10.88	2.003

Reaction of Bu¹O· radicals with di-isopropylamine in the presence of Ph₂C=CH₂, gave the adduct [a_o 3.07, a_m 1.28, a_p 3.34, $a_{H_{\beta}}$ 9.41, a_H (6H) 0.15 G, g 2.0028] between the olefin and the carbon-centred radical Me₂CNHCHMe₂. The latter (a_H^{Me} 17.7, $a_N = a_H^{NH} = 2.4$ G) could be observed in a separate experiment in the absence of Ph₂C=CH₂. No e.s.r. evidence for the aminyl or the aminyl adduct was obtained.

The 1,1-diphenyl-2-chloroethyl radical was also produced by hydrogen abstraction from Ph₂CHCH₂Cl with DTBP, thus obtaining a slightly better spectrum than by addition of \cdot Cl to Ph₂C=CH₂. A similar attempt to generate the Ph₂CCH₂Br radical from Ph₂CHCH₂Br was unsuccessful, this being a further example of the impracticability of detecting β -bromoalkyl radicals in the liquid phase by e.s.r.¹⁴

Results and Discussion

The e.s.r. spectral parameters of the radical adducts of 1,1diphenylethylene are reported in the Table. An examination of these data shows that the ring proton couplings are practically independent of the nature of the MR_n β -substituent, values slightly less than the average being observed only for $MR_n = GePh_3$, $SnPh_3$, and Cl. This indicates that delocalization of the unpaired electron on the β -substituent is essentially the same in all the adducts. Also, the g-factors in the majority of cases are constant (g 2.0028 \pm 0.0001), exceptions being found when MR_n is SnPh₃ (2.0017), SR (2.0031), Cl (2.0032), and CCl₃ (2.0033). The β -proton splittings on the other hand embrace a much larger range (6-11 G), and their temperature dependence is virtually negligible for substituents such as SiPh₃, P(O)(OEt)₂, Bu⁴, 1-adamantyl, CCl₃ while it is of the order of 5-6 mG K⁻¹ for CH₃,Ph,NMe₂, C(O)Ph, and OC(O)Ph.

Conformational Behaviour.—The preferred conformation of the Ph₂CCH₂MR_n adducts can be established from $a_{H\beta}$ and equation (1) if B is known. Since in the Ph_2CH_3 radical a methyl proton splitting of 15.6 G has been reported, 36 B can be taken for this class of radicals as 31.2 G. Therefore a lower limit of 7.8 G for $a_{H_{B}}$ is predicted in the eclipsed conformation (I) where θ is 60° for both the CH₂ protons. A β -proton coupling of 23.4 G can instead be calculated for the staggered conformation (II). The experimental values of 6-11 G are thus indicative that all the β -substituted 1,1-diphenylethyls we have generated adopt a conformation very close to the eclipsed (I). Actually, some deviation of the angle φ between the C_B-M bond and the symmetry axis of the $2p_z$ orbital on C_{α} from 0°, as expected for perfect eclipsing, can be inferred from the linewidth of the central multiplet corresponding to $m(CH_2) = 0$ which in many cases is broader than the low- and high-field multiplets with $m(CH_2) = \pm 1$. In Figure 1, which shows as an example the e.s.r. spectrum of the phosphonyl adduct of $Ph_2C=CH_2$, this linewidth alternation effect is clearly observable at room temperature while it disappears at 370 K. The broadening of the central multiplet is more pronounced at low temperature and with bulky substituents, thus when M is a Group IVB element its width increases in the order SnPh₃ < GePh₃ < $SiPh_3 < Bu^t$. Strong broadening is also apparent at room temperature in the spectra of the 1-adamantyl, $C(Me)_2CN$, and ·CCl₃ adducts, while with the ·CH₃, ·CF₃, and ·SMe adducts deviation from the expected 1:2:1 intensity ratio of the three CH₂ multiplets starts to be observed at lower temperatures. This effect is due to an in-phase modulation of the $a_{H_{R}}$ splittings which are non-equivalent in the minimum energy conformation.

The departure of these radicals from an eclipsed geometry is presumably related to the presence on C_{α} of the two phenyl rings which are somewhat twisted with respect to each other as in benzophenone ketyl.³⁷ As a consequence the MR_n group will experience an asymmetric potential and will be compelled to be bent toward the phenyl exhibiting reduced steric hindrance. The conversion of one enantiomer into the

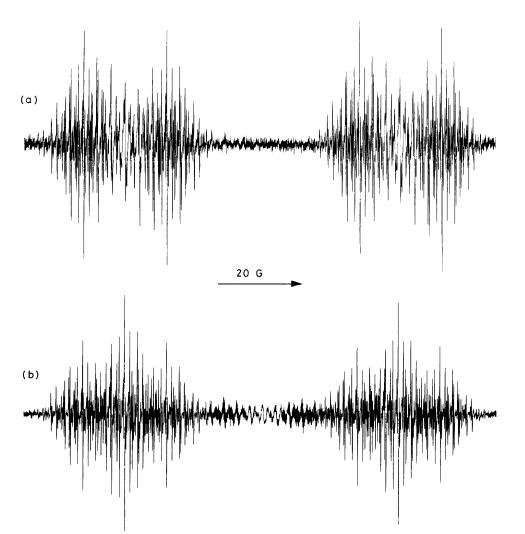


Figure 1. E.s.r. spectrum of the $Ph_2CH_2P(O)(OEt)_2$ radical recorded at room temperature (a) and at 373 K (b). In the central region of the latter one it can also be observed the spectrum of the adduct of $Ph_2C=CH_2$ with the Bu^{t_1} radical

other through a partial rotation of the phenyl groups,³⁷ will force the β -substituent to interchange its position as shown in the Scheme.

If this is the case, we should expect the oscillation of MR_n to be governed by the movement of the two phenyl rings and thus to be characterized by an activation energy typical for the latter process. In principle the activation parameters for the two kind of motions might be determined from the temperature dependence of the line shape if the hyperfine splitting in the slow exchange region can be determined. However, none of the adducts was found to give e.s.r. spectra strong enough to be recorded below 220 K, a temperature which is not sufficiently low to freeze either dynamic process.

Nevertheless, an estimate of the activation energy for the oscillation of the MR_n group may be made without knowing the limiting values a_1 and a_2 of $a_{H_{\beta}}$, by comparing the heights of the central lines of the three multiplets corresponding to $m(CH_2) = \pm 1,0.^{38}$ The widths of these lines are given by equation (12) ³⁹ where $T_{2,0}^{-1}$ is the contribution to the widths

$$T_2^{-1}(m) = T_{2,0}^{-1} + \frac{\gamma_e^2}{8} (m_1 - m_2)^2 (a_1 - a_2)^2 \tau \quad (12)$$

from other relaxation mechanisms and is assumed to be independent of $m(CH_2)$, m_1 and m_2 are the nuclear quantum

numbers of the individual protons, and $\tau = 1/k$, the mean lifetime between jumps. From (12) equation (13) follows or

$$\tau = \left(\frac{T_2^{-1}(0)}{T_2^{-1}(\pm 1)} - 1\right) \frac{8T_{2,0}^{-1}}{\gamma_e^2(a_1 - a_2)^2}$$
(13)

by substituting to the widths the intensities I_m of the lines, (14). Thus, by assuming a temperature dependence for τ of the

$$\tau = \left(\sqrt{\frac{2I_{\pm 1}}{I_0} - 1}\right) \frac{8T_{2,0}^{-1}}{\gamma_e^2(a_1 - a_2)^2}$$
(14)

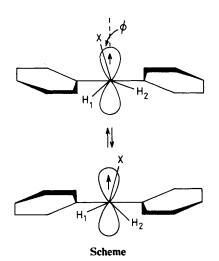
form (15) a plot of log $[(2I_{\pm 1}/I_0)^{\pm} - 1]$ as a function of 1/T will

$$\tau = \tau_0 \exp(E_a/RT) \tag{15}$$

give a straight line with slope $E_a/2.303 R.^{38}$

The determination of the activation energy for the motion of the β -substituent in the triphenylsilyl and phosphonyl adducts of Ph₂C=CH₂, where overlapping of the low- and highfield multiplets with the central line is negligible, gives in both cases E_a 5.5 ± 0.5 kcal mol⁻¹. Although we do not know the corresponding activation energy for the internal rotation of the phenyl rings, presumably this one will not be too different

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from the value measured in the benzophenone ketyl, *i.e.* 6.35 ± 0.2 kcal mol^{-1,37} From the similarity of these two values it may be inferred that the oscillation of the MR_n group is actually governed by the twisting of the phenyls.

Equations (14) and (15) may also provide us with an estimate of the difference $a_1 - a_2$, and therefore of the angle φ (see Scheme), if log $1/\tau_0$ is assumed to be the same as log A for benzophenone ketyl, *i.e.* 13.1. The angle φ may be derived from a_1 and a_2 by rearranging equations (11)-(15). This gives

$$\frac{a_1}{a_2} = \frac{\cos^2(\theta + \varphi)}{\cos^2(\theta - \varphi)}$$
(16)

 φ 6.1 and 4.7° for Ph₂ČCH₂P(O)(OEt)₂ and Ph₂ČCH₂SiPh₃, respectively. Larger values are expected for more bulky substituents such as Bu^t, whose adduct shows selective line broadening even at 370 K; however in no case should φ exceed 15°.

Substituent Effect on the β -Proton Splittings.—In the preceding section it has been shown that the preferred conformation of all the Ph₂CCH₂MR_n adducts is close to the eclipsed (I), with some departure from it of the order of few degrees. The average value of $a_{H\beta}$ we are measuring is then given in terms of equation (1), by (17). For $\varphi < 10^{\circ}$ this quantity does

$$a_{H_{\beta}} = \frac{B}{2} \bigg[\cos^2(\theta + \phi) + \cos^2(\theta - \phi) \bigg]$$
$$\cong B (\cos^2\theta - \phi^2 \cos 2\theta)$$
(17)

not substantially differ from $B \cos^2 \theta$. We may then ignore in the following discussion the small deviation of the MR_n group from a perfectly eclipsed conformation.

On this basis the β -proton splitting is predicted by equation (1) to be 7.8 G in the low temperature limit. Since we are working at relatively high temperatures, torsional oscillations are expected to somehow increase this value. Actually $a_{H_{\beta}}$ lies in the range 9–9.6 G for the large majority of the first row MR_n substituents, the only noticeable exceptions being found with CCl₃ and OC(O)Ph (8.50 G). With second row substituents, on the other hand, the $a_{H_{\beta}}$ values are remarkably different and decrease on going from left to right in the Periodic Table from 11.08 G for SiPh₃ and 6.00 G for Cl.

The scattering of these splittings indicates that equation (1), although extremely valuable in distinguishing between eclipsed and staggered conformations, it is to a certain extent inadequate fully to describe the conformational dependence of the β -proton coupling in these radical adducts, as it has been pre-

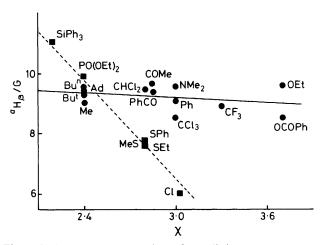


Figure 2. Room temperature hyperfine splitting constants at the β -protons in the Ph₂CCH₂MR radicals against the mutually consistent group electronegativities ⁴⁰ of the MR_n substituents

viously found in related β -substituted ethyls. The two proposed interpretations of this behaviour have been extensively discussed in the Introduction and can be summarized here as follows. (i) Equation (1) is assumed to be valid which ever is the β -substituent and departures of $a_{H\beta}$ from the expected values are due to a change of the $C_{\alpha}C_{\beta}H_{\beta}$ bond angle from sp^3 geometry. (ii) The validity of equation (1) is challenged in the sense that the *B* term should depend on the electronegativity of the β -substituent and perhaps a more complex form of $a_{H\beta}$ on the angle θ should be considered.¹⁸

As far as the latter interpretation is concerned an obvious test of it is to plot the a_{H_B} couplings as function of the electronegativity of the MR_n group to verify whether these two quantities are correlated. The choice of a suitable group electronegativity scale is not an easy one since many sets based on different assumptions are reported in the literature. We used the mutually consistent group electronegativities given by Wells,⁴⁰ and the related plot is shown in Figure 2. It is apparent that no correlation at all exists for first row substituents; actually a least square fitting of these data gives a correlation coefficient of only 0.35. With second row substituents on the other hand a correlation coefficient of 0.98 can be calculated. No better total correlation is obtained if different electronegativity scales, such as that calculated by Huheey,⁴¹ are employed. Hence, from the absence of a common correlation for first and second row MR, groups, it may be inferred that variations of the B term of equation (1) with electronegativity cannot be responsible for the observed changes of a_{H_B} , although the good correlation coefficient found with second-row substituents may have some interesting implications.

Distortions at C_{β} , with the MR_n substituent moving closer to the $2p_z$ orbital on C_{α} when M is a Group VI or VII element and away from it when M belongs to Group IV or V, would explain the trend of the β -proton splittings in the adducts of Ph₂C=CH₂ with second row substituents. However, conclusive evidence for this interpretation is lacking. Moreover, asymmetrical bridging by chlorine has been dismissed in a recent *ab initio* calculation ⁴² performed on the 2-chloroethyl radical where the chloromethyl group has been found essentially tetrahedral. This result is at variance with that of a previous INDO calculation ²⁰ predicting a C_{\alpha}C_{\beta}Cl angle of only 90° in the same radical. It should be mentioned, however, that no explanation of the unusually low value of $a_{H_{\beta}}$ was given in that paper.⁴²

To this purpose it may also be worth comparing the three

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radicals Me2CCH2Cl,7 But2CCH2Cl,11 and Ph2CCH2Cl where the β -proton splittings are similar, *i.e.* 6.18, 6.88, and 6.00 G, respectively, while the ³⁵Cl coupling is much smaller in the diphenyl derivative (21.24, 22.45, and 10.88 G, respectively). If asymmetric bridging is accepted, a measure of it should be given by the ratio a_{Cl}/a_{Hg} , which is 3.44, 3.26, and 1.81 in the three radicals. This means that the amount of deformation at C₆ ought to be considerably lower in the unsaturated 1,1diphenyl-2-chloroethyl. Since bridging is usually associated with the importance of ionic structures such as $R_2C^{++}CH_2$ X^{-43} its extent is expected to be the larger the lower is the ionization potential of the R₂C fragment. As the ionization potential of Ph2CH (7.32 eV)⁴⁴ is slightly less than that of Me2-CH (7.5 eV)⁴⁵ the maximum extent of bridging should be experienced by the Ph₂CCH₂Cl radical. This is just opposite to what can be inferred by the measured hyperfine splitting constants.

In conclusion, it seems that neither interpretation can satisfactorily account for the experimental data and that more experimental and computational work is needed fully to clarify the behaviour of β -substituted ethyl derivatives.

Rate of Trapping.-It is known that photolysis of DTBP in t-butylbenzene leads to the neophyl radical $PhCMe_2CH_2$ which then rearranges to 1-methyl-1-benzylethyl ·CMe₂CH₂-Ph.^{46,47} If the photolysis is carried out at ca. 370 K in the presence of 1,1-diphenylethylene, the radical formed by hydrogen abstraction from t-butylbenzene is trapped by the olefin to give an adduct whose e.s.r. spectrum can be interpreted in terms of the following hyperfine splitting constants: a_o 3.08, a_m 1.28, a_p 3.32, a_{H_B} 9.74 G. A further splitting of 0.15 G from an even number of protons (6 or 8) can also be resolved. Since the latter coupling is practically the same as those from the δ protons in the adducts of Bu^t or CMe₂CN, the trapped radical can be recognized as the rearranged 'CMe₂CH₂Ph. Further, in the hypothetical adduct of neophyl a coupling of ca. 0.5 G from the two γ-protons, as in the Ph₂ČCH₂Buⁿ radical, ought to be expected.

Therefore, if considering that the neophyl rearrangement is a rather slow process $(k_{370 \text{ K}}^{\text{R}} 5.8 \times 10^3 \text{ s}^{-1})^{48}$ the rate constant K^{A} for the addition of this carbon-centred radical to Ph₂C=CH₂ has to be relatively small for a radical reaction. A rough estimate of its higher limit can be obtained by using equation (18).⁴⁹ Although equation (18) is strictly valid only if

$$k^{\mathbf{A}} = \frac{k^{\mathbf{R}}[\mathrm{Ph}_{2}\dot{\mathrm{C}}\mathrm{CH}_{2}\mathrm{CH}_{2}\mathrm{CMe}_{2}\mathrm{Ph}]}{[\mathrm{Ph}_{2}\mathrm{C}=\mathrm{CH}_{2}][\mathrm{Ph}_{2}\dot{\mathrm{C}}\mathrm{CH}_{2}\mathrm{CMe}_{2}\mathrm{CH}_{2}\mathrm{Ph}]} \quad (18)$$

the spin trapping is not reversible and when the resulting adducts do not decay with time, the limit of the trapping rate determined as $k_{370} \kappa^{A} < 400 \ l \ mol^{-1} \ s^{-1}$ compares favourably with that estimated for the addition of the hex-5-enyl radical to di-t-butylethylene ($k_{298} \kappa^{A} < 850 \ l \ mol^{-1} \ s^{-1}$).⁴⁹

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