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## The Iminyl Radical O<sub>2</sub>SN\*\*

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Dedicated to Prof. Werner Uhl on the occasion of his 60th birthday

Sulfonylazides, RSO<sub>2</sub>N<sub>3</sub>, are widely used diazo<sup>[1]</sup> and azide<sup>[2]</sup> transfer reagents, and in particular the photolytic or thermal elimination of N<sub>2</sub> from these azides have been well-studied.<sup>[3]</sup> Although the latter routes are generally used for the in situ generation of sulfonyl nitrenes, RSO<sub>2</sub>N, both were found to be complicated owing to secondary product formation, such as SO<sub>2</sub> + RN, as well as the pseudo Curtius rearrangement product RN=SO<sub>2</sub>.<sup>[3,4]</sup>

Recently it became apparent that initially formed reactive singlet sulfonyl nitrenes might be rapidly quenched by efficient intersystem crossing (ISC) to yield the lowerenergy triplet nitrenes of sluggish reactivity.<sup>[4]</sup> Thermally persistent triplet FSO<sub>2</sub>N was produced by flash pyrolysis (ca. 800 °C) of FSO<sub>2</sub>N<sub>3</sub> in yields up to 66 %.<sup>[5]</sup> Only traces of SO<sub>2</sub> and FSO<sub>2</sub> were found as byproducts, and the lifetime of triplet FSO<sub>2</sub>N in the gas phase was found to be dominated by precursor concentration-dependent triplet nitrene dimerization.<sup>[5]</sup> Triplet nitrenes, produced by flash-pyrolysis of selected azides, have also offered unique access to synthetically challenging molecules, such as OCN–NCO,<sup>[6]</sup> *cyclo*-N<sub>2</sub>CO,<sup>[7]</sup> OPN/ONP,<sup>[8]</sup> SPN/SNP/*cyclo*-PSN,<sup>[9]</sup> and the FSO<sub>2</sub> radical.<sup>[5]</sup>

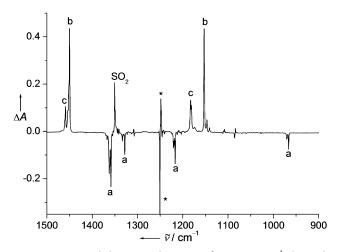
Herein, we extend our flash vacuum pyrolysis studies to the simple alkyl sulfonyl azides  $CF_3SO_2N_3$  and  $CH_3SO_2N_3$ , and present the detection and the photochemistry of the novel  $O_2SN$  (sulfonyliminyl) radical.

Flash vacuum pyrolysis of  $CF_3SO_2N_3$  highly diluted in argon (1:500) was performed by passing the gas mixture through a hot quartz furnace (ca. 800 °C, inside diameter 1.0 mm, length 30 mm). The reaction mixture was immediately deposited onto the cold matrix support (16 K) in a high vacuum (see the Supporting Information for details). The IR spectrum of the deposit reveals almost complete decomposition of the azide (Supporting Information, Figure S1). Along with the known IR bands of SO<sub>2</sub>, SO<sub>3</sub>, CF<sub>3</sub>, and C<sub>2</sub>F<sub>6</sub>, there are a number of new IR bands showing distinct <sup>15</sup>N isotopic shifts in experiments using a mixture of  $CF_3SO_2^{15}N_{\alpha}NN$  and  $CF_3SO_2NN^{15}N_{\gamma}$  (1:1). To distinguish between different possi-

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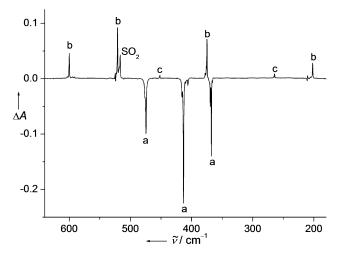
ble carriers of the new IR bands, the deposit was subjected to light with wavelength  $\lambda > 360$  nm and radiation from a 365 nm LED source, which both were found to efficiently deplete the same set of the new bands. The mid-IR difference spectrum obtained from the first experiment is shown in Figure 1.



**Figure 1.** IR spectral changes in the region of 1500–900 cm<sup>-1</sup> obtained after photolysis ( $\lambda > 360$  nm) of Ar matrix-isolated (16 K) pyrolysis products of CF<sub>3</sub>SO<sub>2</sub>N<sub>3</sub>. Bands assigned to O<sub>2</sub>SN (a) decreased, and bands that are due to SO<sub>2</sub>, *syn* (b) and *anti* OSNO (c) appeared. Residual bands of the CF<sub>3</sub> radical (marked by asterisks) occurred in the difference spectrum owing to matrix site changes upon photolysis.

Based on their same photolytic behavior, four IR bands at 1358.6, 1328.3, 1216.9, and 966.9 cm<sup>-1</sup> in the mid-IR spectrum (Figure 1, a) are assigned to species a. Its IR spectrum was completed by recording three additional bands in the far-IR region (Figure 2) at 474.5, 413.2, and  $367.6 \text{ cm}^{-1}$ . Each band exhibits a satellite that is due to matrix sites, and along with the complete set of <sup>14/15</sup>N isotopic shifts, bands associated with the naturally abundant <sup>34</sup>S isotopologue were also observed. The experimental vibrational data are compared in Table 1 with predicted values for the planar  $O_2SN$  ( $C_{2\nu}$ ) radical obtained at the DFT B3LYP/6-311 + G(3df) level. With the exception of the band at 1328.3 cm<sup>-1</sup>, the agreement is very good. This band is attributed to the combination 966.9  $(a_1)$  + 367.6 (b<sub>2</sub>) = 1334.5 cm<sup>-1</sup> (b<sub>2</sub>), which takes intensity through a Fermi resonance from the strong fundamental at 1358.6 cm  $^{-1}$  (b<sub>2</sub>). Apart from the six fundamentals, two combination and two overtone bands of O2SN were observed and assigned.

The assignment of the IR spectrum of O<sub>2</sub>SN is given in Table 1. The two strongest bands at 1358.6 and 1216.9 cm<sup>-1</sup> are readily attributed to antisymmetric ( $\nu_5$ ) and symmetric



**Figure 2.** IR spectral changes in the far-IR region of 640–180 cm<sup>-1</sup> obtained after irradiation ( $\lambda > 360$  nm) of the Ar matrix-isolated (16 K) pyrolysis products of CF<sub>3</sub>SO<sub>2</sub>N<sub>3</sub>. Bands assigned to O<sub>2</sub>SN (a) decreased, and bands due to SO<sub>2</sub>, *syn* (b) and *anti* OSNO (c) appeared.

**Table 1:** Observed and calculated IR frequencies  $\nu$ , isotope shifts  $\Delta \nu$  (cm<sup>-1</sup>) and band assignments for the O<sub>2</sub>SN (<sup>2</sup>B<sub>2</sub>, C<sub>2 $\nu$ </sub>) radical.

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ν		$\Delta \nu$	( <sup>14/15</sup> N)	$\Delta \nu$	( <sup>32/34</sup> S)	Mode
expt (Ar) <sup>[a]</sup>	calcd <sup>[b]</sup>	expt	calcd	expt	calcd	$v_i^{[c]}$
2700.8 (1)		< 0.5				2v <sub>5</sub>
2561.0 (1)		8.6				$\nu_1 + \nu_5$
2424.5 (2)		13.0				$2\nu_1$
1358.6 (100)	1336 (156)	1.7	0.1	16.5	17.4	$\nu_5$ , $\nu_{as}(SO_2)$
1328.3 (18)		21.9				$\nu_3 + \nu_6$
1216.9 (37)	1201 (46)	6.5	6.4	13.5	14.0	$\nu_1, \nu_s(SO_2)$
966.9 (15)	954 (9)	17.8	19.2	1.2	1.1	$\nu_2$ , $\nu$ (SN)
474.5 (26)	462 (19)	1.4	1.5	2.4	2.5	$\nu_3, \delta(SO_2)$
413.2 (38)	414 (20)	1.8	1.9	7.1	7.5	$\nu_4$ , $\delta_{\text{o.o.p.}}$
367.6 (22)	358 (20)	6.0	6.0	1.8	1.8	$\nu_{\rm 6},  \delta({\rm OSN})$

[a] Band positions of the most intense matrix sites in solid argon at 16 K, and relative intensities (in parenthesis) based on integrated areas of all matrix sites. [b] Harmonic frequencies calculated at the DFT B3LYP/6-311 + G(3df) level of theory and scaled by a factor of 0.9679. Calculated IR intensities (km mol<sup>-1</sup>) are given in parentheses. [c] Numbering of fundamentals  $v_i$  (i=1–6) according to their symmetry ( $a_1$ : i=1–3,  $b_1$ : 4,  $b_2$ : 5 and 6) and frequency.

( $\nu_1$ ) stretches of a SO<sub>2</sub> moiety, respectively, and the assignment of the SN stretch ( $\nu_2$ ) to the band at 966.9 cm<sup>-1</sup> is supported by its large <sup>15</sup>N isotope shift of 17.8 cm<sup>-1</sup>. A significant <sup>34</sup>S isotope shift for the band at 413.2 cm<sup>-1</sup> is consistent with its assignment to an out-of-plane mode  $\nu_4$  of the planar radical. Thus, the band at 474.5 cm<sup>-1</sup> should be attributed to the SO<sub>2</sub> bending mode ( $\nu_3$ ), which is lower in frequency than that of free SO<sub>2</sub> (516.7 cm<sup>-1</sup>, Ar matrix) owing to vibrational coupling with the SN stretch in O<sub>2</sub>SN, and the lowest energy vibration at 367.6 cm<sup>-1</sup> to the OSN bending mode  $\nu_6$ . We note that the O<sub>2</sub>SN radical was also detected among the flash pyrolysis products of CH<sub>3</sub>SO<sub>2</sub>N<sub>3</sub> (Supporting Information, Figure S2). However, the main products were SO<sub>2</sub> and CH<sub>2</sub>NH, and the yield of O<sub>2</sub>SN was significantly lower compared to that found by pyrolysis of CF<sub>3</sub>SO<sub>2</sub>N<sub>3</sub>.

The identification of the photolysis products of O<sub>2</sub>SN is straightforward. Traces of SO<sub>2</sub> formed during UV irradiation  $(\lambda > 360 \text{ nm})$  of Ar matrix-isolated O<sub>2</sub>SN (Figure 1) indicate its photodecomposition into SO<sub>2</sub> and nitrogen atoms. Apart from the known SO<sub>2</sub> bands, strong bands appeared at 1450.3 and 1152.9 cm<sup>-1</sup> (Figure 1, b) and weaker bands at 1459.2 and 1182.9 cm<sup>-1</sup> (Figure 1, c). These band positions are close to reported NO and SO stretching frequencies for the open-chain molecules *syn* OSNO (1450.8 and 1154.9 cm<sup>-1</sup>) and *anti* OSNO (1456.0 and 1181.2 cm<sup>-1</sup>) isolated in argon matrices.<sup>[10]</sup> Large <sup>15</sup>N isotope shifts obtained for the NO stretches (*syn* 24.9 cm<sup>-1</sup>, *anti* 24.8 cm<sup>-1</sup>) support their assignment, whereas the SO stretches are much less affected by <sup>15</sup>N isotopic substitution (Supporting Information, Figure S3).

In previous work, the bent OSNO isomers were obtained in low yield by laser photolysis (248 nm) of matrices containing OCS and NO<sub>2</sub>.<sup>[10]</sup> Only their two strongest IR bands were observed, and the presence of a CO molecules along with the OSNO radicals in the matrix cage are responsible for the differences between reported and our band positions. IR radiation from the globar source was found to convert *anti* OSNO to the more stable *syn* isomer.<sup>[10]</sup> We also observed this conversion using 365 nm LED radiation.

The efficient photoisomerization of  $O_2SN(C_{2\nu})$  into the chain-like OSNO ( $C_3$ ) radicals allows for a more complete analysis of their IR spectra (Figure 2). Their vibrational data are summarized in Table 2 and Table 3, respectively. For the *syn* isomer the weak SN stretching band ( $\nu_3$ ) appeared at

**Table 2:** Observed and calculated IR frequencies  $\nu$  and isotope shifts  $\Delta \nu$  (cm<sup>-1</sup>) for syn OSNO (<sup>2</sup>A', C<sub>s</sub>).

ν		$\Delta \nu$ ( <sup>14/15</sup> N)		$\Delta \nu$ ( <sup>32/34</sup> S)		Mode
expt (Ar) <sup>[a]</sup>	calcd <sup>[b]</sup>	expt	calcd	expt	calcd	
2866.7 (6)		48.3		< 0.5		2v <sub>1</sub>
1450.3 (100)	1461 (188)	24.9	25.9	< 0.5	0.2	ν <sub>1</sub> , ν(NO)
1152.9 (87)	1144 (183)	2.3	0.3	11.8	11.8	ν <sub>2</sub> , ν(SO)
600.2 (4)	614 (3)	5.6	7.0	8.4	8.9	ν <sub>3</sub> , ν(SN)
521.2 (15)	520 (18)	10.9	11.2	2.3	2.6	ν <sub>4</sub> , δ(OSN)
375.0 (11)	393 (9)	7.5	8.1	1.8	2.0	ν₅, δ(ONS)
202.2 (4)	200 (3)	< 0.5	0.4	< 0.5	0.3	ν <sub>6</sub> , τ

[a] Band position of the most intense matrix site in solid argon at 16 K, relative intensities (in parenthesis) based on the integrated areas of all matrix sites. [b] Harmonic frequencies calculated at the B3LYP/6-311 + G(3df) level of theory and scaled by a factor of 0.9679; calculated IR intensities [km mol<sup>-1</sup>] are given in parentheses.

600.2 cm<sup>-1</sup> (Figure 2, b) along with three additional deformations in the far-IR region. Owing to a lower abundance of the less-stable *anti* conformer, the assignment of its weak SN stretch is difficult. A weak overtone band  $2\nu_1$  was however observed for both isomers.

Upon 266 nm laser irradiation, *syn* OSNO rearranged into the less stable *anti* isomer. Using mercury arc radiation ( $\lambda >$  255 nm), traces of SNO radicals were also detected by means of two weak IR bands at 1593.3 ( $2\nu_2$ ,  $\Delta\nu$  (<sup>14/15</sup>N) = 28.3 cm<sup>-1</sup>) and 1522.6 cm<sup>-1</sup> ( $\nu_1$ ,  $\Delta\nu$  (<sup>14/15</sup>N) = 24.7 cm<sup>-1</sup>)<sup>[11]</sup> (Supporting Information, Figure S4).

**Table 3:** Observed and calculated IR frequencies  $\nu$  and isotope shifts  $\Delta \nu$  (cm<sup>-1</sup>) for *anti* OSNO (<sup>2</sup>A', C<sub>s</sub>).

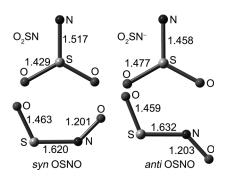
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ν		Δ	ν ( <sup>14/15</sup> N)	Mode
expt (Ar) <sup>[a]</sup>	calcd <sup>[b]</sup>	expt	calcd	
2887.3 (2)		49.8		2v <sub>1</sub>
1459.2 (67)	1471 (253)	24.8	26.7	ν <sub>1</sub> , ν(NO)
1182.9 (100)	1162 (206)	1.0	0.1	ν <sub>2</sub> , ν(SO)
	582 (2)		10.9	ν <sub>3</sub> , ν(SN)
452.0 (2)	451 (5)	3.8	4.0	ν <sub>4</sub> , δ(OSN)
272.6 (<1)	279 (0.3)	4.7	5.3	ν <sub>6</sub> , τ
264.4 (3)	261 (8)	2.1	2.1	ν₅, δ(ONS)

[a] Band position of the most intense matrix site in solid argon at 16 K, relative intensities (in parenthesis) based on the integrated areas of all matrix sites. [b] Harmonic frequencies calculated at the B3LYP/6-311 + G(3df) level of theory and scaled by a factor of 0.9679; the calculated IR intensities [km mol<sup>-1</sup>] are given in parentheses.

In the initially obtained IR spectrum of the pyrolysis products of  $CF_3SO_2N_3$ , some bands remained unassigned. They also revealed distinct <sup>15</sup>N isotope shifts (Supporting Information, Figure S5), but remained almost unaffected by any of the aforementioned irradiations. These bands appeared at 1413.1, 1313.6, 1243.8, 1213.1, 1196.3, and 1045.4 cm<sup>-1</sup>, indicating the presence of SO<sub>2</sub> (1413.1, 1045.4 cm<sup>-1</sup>) and CF<sub>3</sub> (1243.8, 1213.1, 1196.3 cm<sup>-1</sup>) group vibrations. A large <sup>15</sup>N isotope shifts (19.3 cm<sup>-1</sup>) for the strong band at 1313.6 cm<sup>-1</sup> suggests an S=N stretching band. IR frequencies calculated for the most likely carrier, *N*-sulfonyl imine CF<sub>3</sub>N=SO<sub>2</sub>, fit very well with the observations (Supporting Information, Table S1). Three weaker far-IR bands at 585.3, 471.9, and 406.0 cm<sup>-1</sup> were also assigned to this species.

The identification of intermediates from the flash pyrolysis of sulfonyl azides  $RSO_2N_3$  (R = CH<sub>3</sub>, CF<sub>3</sub>) sheds light on their thermal decomposition mechanism and the formation of O<sub>2</sub>SN. Upon pyrolysis, both azides are expected to eliminate N<sub>2</sub> furnishing the short-lived singlet nitrene RSO<sub>2</sub>N. Initially formed singlet RSO<sub>2</sub>N may undergo: 1) ISC to the triplet ground state; 2) Curtius-type rearrangement to  $RN = SO_2$ ; or 3) R-S bond fission to yield the radical pair O<sub>2</sub>SN and R. The favorable formation of O<sub>2</sub>SN by flash pyrolysis of CF<sub>3</sub>SO<sub>2</sub>N<sub>3</sub> according to (3) is consistent with the weak  $F_3C^{\delta+}$ -S<sup> $\delta+$ </sup> bond in the nitrene, CF<sub>3</sub>SO<sub>2</sub>N, for which a dissociation energy of  $66 \text{ kJmol}^{-1}$  (singlet CF<sub>3</sub>SO<sub>2</sub>N) is predicted by preliminary calculations using the complete basis set CBS-QB3 method.<sup>[12]</sup> For the methyl analogue, the C-S bond energies of singlet  $CH_3SO_2N$  is considerably higher (117 kJ mol<sup>-1</sup>), suggesting (2) and (3) are competing processes in the flash pyrolysis of CH<sub>3</sub>SO<sub>2</sub>N<sub>3</sub>. The Curtius-type rearrangement (2) eventually vields fragments (SO<sub>2</sub> and CH<sub>2</sub>NH) of CH<sub>3</sub>NSO<sub>2</sub>, although their formation by radical reactions between CH<sub>3</sub> and NSO<sub>2</sub>, formed under the pyrolysis conditions, cannot be excluded.

In recent studies on NO<sub>2</sub>S species,<sup>[10,13]</sup> the planar O<sub>2</sub>SN radical was not considered at all. We have calculated the molecular structures and energies of various planar NO<sub>2</sub>S isomers on the doublet potential energy surface using different theoretical methods (Supporting Information, Table S2). At all of the levels applied, *syn* OSNO (<sup>2</sup>A', Figure 3) was found to be the lowest-energy isomer, O<sub>2</sub>SN (<sup>2</sup>B<sub>2</sub>) and *anti* OSNO (<sup>2</sup>A', Figure 3) were found to be slightly higher in



*Figure 3.* Calculated planar molecular structures (bond lengths in Å) of  $O_2SN$  ( $^2B_2$ ,  $\angle OSO = 121.7^\circ$ ),  $O_2SN^-$  ( $^1A_1$ ,  $\angle OSO = 110.1^\circ$ ), syn OSNO ( $^2A'$ ,  $\angle OSN = 115.3^\circ$ ,  $\angle SNO = 134.9^\circ$ ), and anti OSNO ( $^2A'$ ,  $\angle OSN = 111.9^\circ$ ,  $\angle SNO = 132.5^\circ$ ) at the B3LYP/6-311 + G(3df) level of theory.

energy by, respectively, 17.4 (13.6) kJ mol<sup>-1</sup> and 15.2 (17.3) kJ mol<sup>-1</sup> at the CCSD(T)//MP2 (CBS-QB3) levels of theory. We note the existence of another weakly interacting *syn* OS…NO minimum (SN 2.047 Å; Supporting Information, Table S2), having the unpaired electron in a bonding a" orbital; however, in the following we will focus on the experimentally confirmed isomers only.

DFT calculations predict a planar molecular structure of  $C_{2\nu}$  symmetry for the O<sub>2</sub>SN radical (Figure 3). A comparison with the related structure of the *N*-sulfonyl imine CF<sub>3</sub>N=SO<sub>2</sub>, calculated at the same level of theory (see Supporting Information), revealed similar S=O bond lengths and OSO angles (O<sub>2</sub>SN: 1.429 Å (S=O),  $\angle$ OSO = 121.7°; CF<sub>3</sub>N = SO<sub>2</sub>: 1.427 and 1.422 Å (S=O),  $\angle$ OSO = 121.7°) for these two related molecules. The S=N bond in O<sub>2</sub>SN (1.517 Å) is slightly longer than that of CF<sub>3</sub>NSO<sub>2</sub> (1.507 Å).

Although O<sub>2</sub>SN (<sup>2</sup>B<sub>2</sub>) is isostructural and isoelectronic to the known FCO<sub>2</sub> ( $^{2}B_{2}$ ) radical,<sup>[14]</sup> the electronic structures of these two radicals are very different. The electron spin density of FCO<sub>2</sub> is equally shared between the two oxygen atoms,<sup>[14]</sup> but localized at the nitrogen atom in O<sub>2</sub>SN. Calculations for the closed-shell O<sub>2</sub>SN<sup>-</sup> anion (Figure 3)<sup>[15]</sup> reveal a significantly shortened S=N (1.458 Å) associated with elongated S= O bond lengths (1.477 Å), indicating that the unpaired electron in  $O_2SN$  resides in a  $\pi(NS)$  bonding (b<sub>2</sub>) molecular orbital. The O<sub>2</sub>SN<sup>-</sup> anion is isoelectronic to recently detected F<sub>2</sub>PN.<sup>[16]</sup> According to natural bond orbital (NBO) analysis, two doubly occupied atomic p-type orbitals of nitrogen in these molecules are engaged in  $\pi$  bonding to the central atom, whereas in the  $O_2$ SN radical the corresponding in-plane (b<sub>2</sub>) HOMO is only half-filled (Supporting Information, Figure S6, Table S3). The lowest-energy transition between the highest doubly occupied b<sub>1</sub> and the singly occupied b<sub>2</sub> orbitals in O<sub>2</sub>SN is forbidden by symmetry (A<sub>2</sub>). However, timedependent (TD) DFT calculations predict reasonably strong UV bands at 322 (A<sub>1</sub>, f=0.0094), 224 (B<sub>1</sub>, 0.0025) and 217 nm (B<sub>2</sub>, 0.0074). These predictions are consistent with the observed photosensitivity of O<sub>2</sub>SN and the presence of two structured UV absorptions centered at  $\lambda_{max} \approx 350 \text{ nm}$  and below 250 nm (Supporting Information, Figure S7). The latter absorption shows a regular vibrational spacing of about 490 cm<sup>-1</sup>. These absorptions are associated with O<sub>2</sub>SN, as they

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vanished completely after irradiation with light of  $\lambda > 360$  nm (Supporting Information, Figures S8 and S9).

The SN bond lengths predicted for the planar bent OSNO isomers of 1.62–1.63 Å (Figure 3) are between the expected values for single (1.74 Å) and double SN bonds (1.54 Å).<sup>[17]</sup> In a MO scheme, bonding interactions between the diatomic fragments OS and NO may arise from a combinations of their partially filled  $\pi^*$  fragment orbitals (for the frontier molecular orbitals, see the Supporting Information, Figure S6). For the bent OSNO<sup>-</sup> anion (24 valence electrons), the corresponding in-plane (a') and out-of-plane (a'') ( $\pi^*\pi^*$ ) bonding interactions affords a formal S=N double bond (Supporting Information, Table S4). For OSNO (23 valence electrons), the unpaired electron resides in the a' bonding combination (Supporting Information, Figure S6). However, the spin densities for *syn* and *anti* OSNO were calculated to be localized mainly at the NO moiety.

A strong UV absorption at  $\lambda_{max} \approx 285$  nm with a regular vibrational spacing of about 1000 cm<sup>-1</sup> occurred in the experimental spectrum after UV photolysis of O<sub>2</sub>SN (Supporting Information, Figure S7), which is most likely associated with the bent OSNO  $(^{2}A')$  radicals. This assignment is supported by TD DFT calculations, which predict strong UV transitions (f > 0.09) for syn and anti OSNO at 266 and 269 nm, respectively (Supporting Information, Figure S7). The anti to syn interconversion was predicted to proceed by a low-energy N-inversion process rather than by S-N bond rotation, and for both isomers low barriers of less than 40 kJ mol<sup>-1</sup> were reported for their dissociation into the diatomic fragments OS + NO.<sup>[13]</sup> In contrast, the thermally robust nature of the O<sub>2</sub>SN radical is confirmed by the absence of SO and NO in the flash pyrolysis products. On the other hand, photoisomerization from O2SN to bent OSNO radicals may proceed either via a cyclic transition state or by photodissociation and recombination of the photofragments OSN+O in a matrix cage. The dissociative mechanism has been suggested for the rearrangement of the related molecules NO<sub>3</sub>  $(\rightarrow ON + O_2)^{[18]}$  and SO<sub>3</sub>  $(\rightarrow OSOO)$ .<sup>[19]</sup>

In summary, flash vacuum pyrolysis of sulfonyl azides,  $CF_3SO_2N_3$  and  $CH_3SO_2N_3$ , yields the planar  $O_2SN$  (<sup>2</sup>B<sub>2</sub>) radical. IR and UV spectra of the novel radical have been assigned and its photorearrangement into the planar bent isomers, *syn* and *anti* OSNO (<sup>2</sup>A'), has been observed in solid noble-gas matrices. The facile generation of  $O_2SN$  in the gas phase from two easily accessible azides will stimulate further studies on its spectroscopy, structure, and chemistry.

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