

## Vibrational spectra of hexabromobenzene and hexaiodobenzene

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**Abstract**—Infrared and Raman spectra of polycrystalline  $C_6Br_6$  and  $C_6I_6$  are measured at both room temperature and liquid nitrogen temperatures. Vibrational assignment of fundamental modes as well as combination bands are established on the basis of normal coordinate analysis.

### INTRODUCTION

Vibrational spectra of hexabromobenzene  $C_6Br_6$  and hexaiodobenzene  $C_6I_6$  have been studied by a few investigators [1–3]. DELORME *et al* [1] reported i.r. spectra of these molecules in the region from 4000–70  $cm^{-1}$ , and MARSAULT *et al* [2] measured far i.r. spectra in the region from 200 to 15  $cm^{-1}$ . Raman spectra of  $C_6Br_6$  and  $C_6I_6$  were studied by ABRAMOWITZ *et al* [3]. These authors were interested primarily in the planarity of these hexahalobenzenes. The assignments of intramolecular and lattice vibrations are yet to be settled. To establish the vibrational assignments consistent with those in other hexahalobenzenes, i.r. and Raman spectra of these molecules were remeasured. Normal coordinate analysis has been carried out, and the complete assignments are presented for  $C_6Br_6$  and  $C_6I_6$ .

### EXPERIMENTAL

The commercially available sample of  $C_6Br_6$  was purified by repeated recrystallization from saturated benzene solution.  $C_6I_6$  was prepared from benzene, fuming sulfuric acid and solid iodine and recrystallized from boiling nitrobenzene solution as described by RUPP [4] and DURAND *et al* [5]. Result of elementary analysis of  $C_6I_6$  is as follows: expected C 8.65%, found C 8.76%.

Raman spectra of powder  $C_6Br_6$  were taken on a JEOL JRS-400D double monochromator using 5145 and 4880 Å excitation lines from an argon ion gas laser with an output of 700 mw. For orange polycrystalline  $C_6I_6$ , a Spex Model 1401 double monochromator with 6328 Å excitation line from a He–Ne gas laser was used with an output of 50 mw. Raman lines scattered from the surface of  $C_6I_6$  disc were detected.

Infrared spectra of  $C_6Br_6$  and  $C_6I_6$  were recorded on a Hitachi G2 (4000–400  $cm^{-1}$ ), EPI-L (700–200  $cm^{-1}$ ) and FIS-1 (500–60  $cm^{-1}$ ) spectrometers. Far i.r. spectra were also taken on a Hitachi 070 interferometer to check the reproducibility in the region from 200–60  $cm^{-1}$ . Both spectra obtained by using different spectrometers match completely.

Low temperature measurements of i.r. and Raman spectra were made with a conventional cryostat holding liquid nitrogen. Infrared and Raman spectra of  $C_6Br_6$  and  $C_6I_6$  are given in Fig. 1 and 2, respectively.

### NORMAL COORDINATE TREATMENT

The planarity of the  $C_6X_6$  type molecules ( $X = F, Cl, Br, \text{ and } I$ ) has been the subject of a number of investigations. X-ray diffraction study (6, 7) indicates the puckered  $D_{3d}$  structure for the  $C_6I_6$  crystal. However, six iodine atoms are located only 0.04 Å off the plane of the benzene ring. An electron diffraction investigation of gaseous  $C_6Br_6$  indicates a distorted configuration [8], while X-ray analysis [9] supported planar configuration in the crystalline state which is isomorphous to crystalline  $C_6Cl_6$  [10, 11]. Since vibrational frequencies are insensitive to slight geometrical changes of configuration, normal coordinate treatments were undertaken on a planar  $D_{6h}$  model. The calculated frequencies may be compared with those for planar  $C_6Cl_6$  [12].

According to the group theory, the irreducible representations for the normal modes of  $C_6X_6$  molecule assuming  $D_{6h}$  symmetry are:  $\Gamma(D_{6h}) = (2a_{1g} + a_{2g} + 4e_{2g} + 2b_{1u} + 2b_{2u} + 3e_{1u}) + (2b_{2g} + e_{1g} + a_{2u} + 2e_{2u})$ , where the symmetry species enclosed in the first set of braces correspond to the

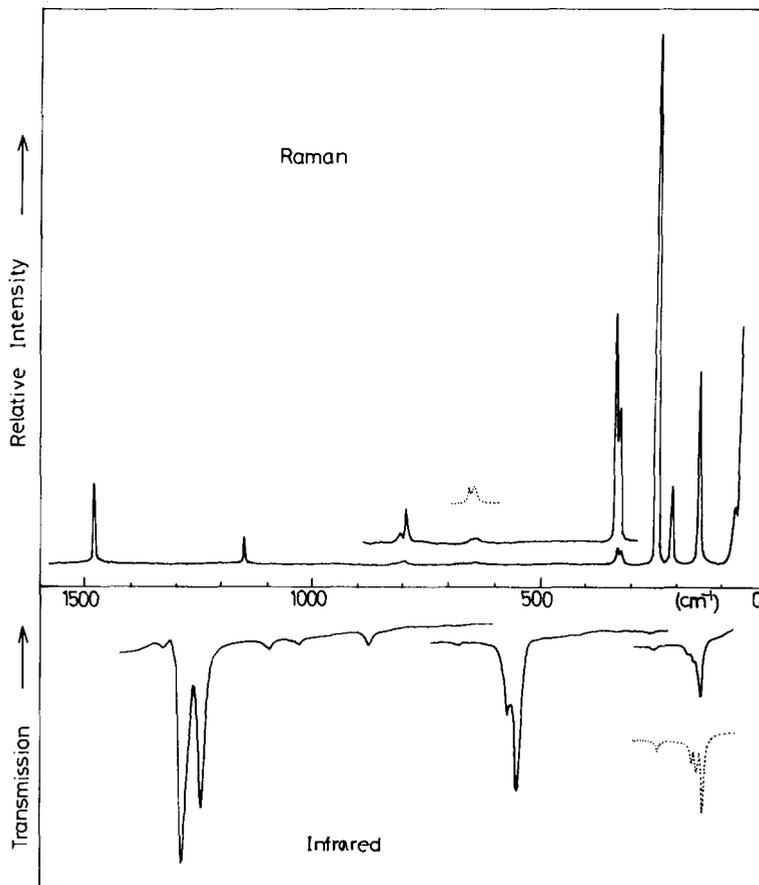


Fig. 1. Raman and i.r. spectra of solid  $C_6Br_6$  at room temperature. (Dotted curves indicate the spectra obtained at liquid nitrogen temperature.)

in-plane modes, and the species in the second braces correspond to the out-of-plane modes. Of these twenty normal modes, seven modes ( $2a_{1g} + 4e_{2g} + e_{1g}$ ) are Raman active, and four ( $3e_{1u} + a_{2u}$ ) are i.r. active. The correlation between  $D_{6h}$  and  $D_{3d}$  indicates that two  $b_{2g}$  modes are Raman active and two  $b_{1u}$  and two  $e_{2u}$  modes are i.r. active in a puckered  $D_{3d}$  structure. Hereafter the notation of symmetry species in  $D_{6h}$  symmetry will be used.

Internal symmetry coordinates chosen are the same as those of HARADA *et al.* in the calculation of the benzene molecule [13]. The potential energy is given as [13]:

$$2V_{in} = 2V_{UBFF} + 2 \left( o_{cx} \sum_{i=1}^6 s_i s_{i+1} + m_{cx} \sum_{i=1}^6 s_i s_{i+2} + p_{cx} \sum_{i=1}^3 s_i s_{i+3} \right) + 2 \left( o_{ccx} \sum_{i=1}^6 \psi_i \psi_{i+1} + m_{ccx} \sum_{i=1}^6 \psi_i \psi_{i+2} + p_{ccx} \sum_{i=1}^3 \psi_i \psi_{i+3} \right)$$

$$+ 2 \left( o_{cc} \sum_{i=1}^6 t_i t_{i+1} + m_{cc} \sum_{i=1}^6 t_i t_{i+2} + p_{cc} \sum_{i=1}^3 t_i t_{i+3} \right) + 2 \left( o_{ccc} \sum_{i=1}^6 \phi_i \phi_{i+1} + m_{ccc} \sum_{i=1}^6 \phi_i \phi_{i+2} + p_{ccc} \sum_{i=1}^3 \phi_i \phi_{i+3} \right)$$

$$2V_{out} = A \sum_{i=1}^6 \gamma_i^2 + B \sum_{i=1}^6 \delta_i^2 + 2C \left( \sum_{i=1}^6 \gamma_i \delta_i - \sum_{i=1}^6 \gamma_i \delta_{i-1} \right) + 2D \left( \sum_{i=1}^6 \gamma_i \delta_{i+1} - \sum_{i=1}^6 \gamma_i \delta_{i-2} \right) + 2 \left( a_0 \sum_{i=1}^6 \gamma_i \gamma_{i+1} + a_m \sum_{i=1}^6 \gamma_i \gamma_{i+2} + a_p \sum_{i=1}^3 \gamma_i \gamma_{i+3} \right) + 2 \left( b_0 \sum_{i=1}^6 \delta_i \delta_{i+1} + b_m \sum_{i=1}^6 \delta_i \delta_{i+2} + b_p \sum_{i=1}^3 \delta_i \delta_{i+3} \right).$$

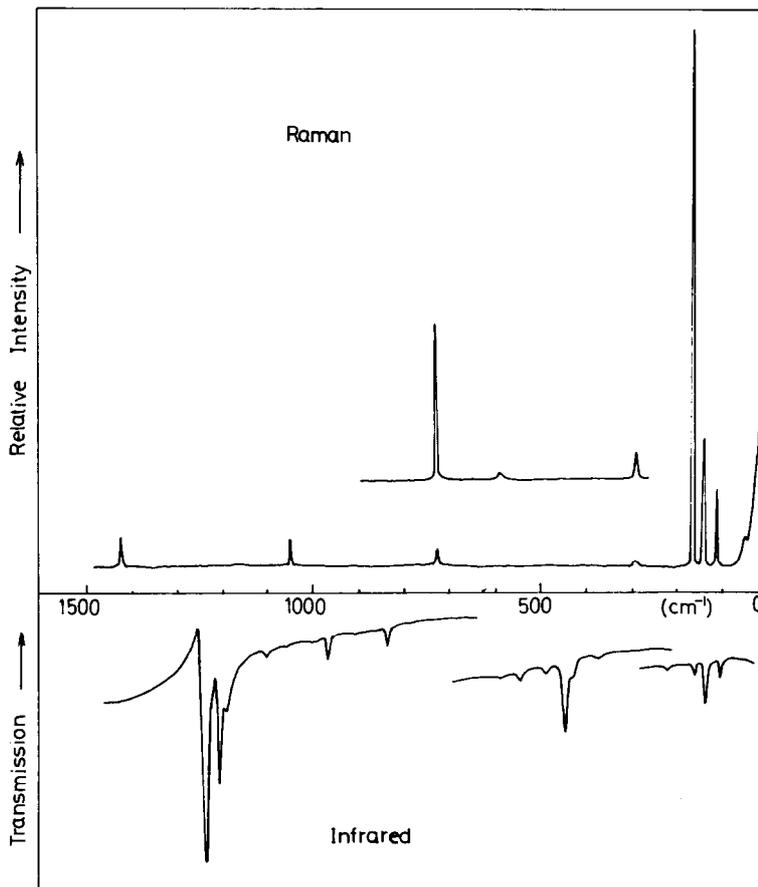


Fig. 2. Raman and i.r. spectra of solid  $C_6I_6$  at room temperature.

In the preliminary calculation, the values of force constants adopted are the same as those used for the benzene molecule [13] except  $K_{cx}$ ,  $H_{cx}$ , and  $F_{cx}$ , which are listed in Table 1. These values are transferred from  $CH_3Br$ ,  $C_2Br_4$  and  $CH_3I$  [14]. Calculated frequencies are also shown in Table 1.

#### VIBRATIONAL ASSIGNMENTS

##### Raman spectra

Although general features of the Raman spectra of  $C_6Br_6$  and  $C_6I_6$  are similar to those reported by ABRAMOWITZ *et al.* [3], we observe a few additional bands at 648, 637, 68, 47, 41, 34, 22, and  $20\text{ cm}^{-1}$  for the  $C_6Br_6$  molecule at room temperature. Also, the bands at  $489\text{ cm}^{-1}$  and  $125\text{ cm}^{-1}$  they reported are completely missing. For  $C_6I_6$ , a weak band at  $592\text{ cm}^{-1}$  was observed, while a sharp band at  $383\text{ cm}^{-1}$  [3] was not recorded even at low tempera-

Table 1. Vibrational frequencies and force constants of  $C_6Br_6$  and  $C_6I_6$  in preliminary calculation

	Calculated frequencies ( $\text{cm}^{-1}$ )	
	$C_6Br_6$	$C_6I_6$
$a_{1g}$	1104, 229	1090, 173
$b_{2g}$	748, 61	746, 45
$e_{1g}$	300	290
$e_{2g}$	1559, 797, 197, 96	1547, 769, 152, 69
$a_{2u}$	117	107
$b_{1u}$	1142, 230	1122, 173
$e_{1u}$	1300, 539, 118	1284, 487, 86
$e_{2u}$	561, 46	555, 34
	Force constants (md/A)	
	$C_6Br_6$	$C_6I_6$
$K_{cx}$	1.757	1.547
$H_{ccx}$	0.08	0.07
$F_{ccx}$	0.68	0.60

Values of other force constants are the same as used for benzene molecule [13].

ture. No splitting was observed around  $592\text{ cm}^{-1}$  at low temperature.

BATES *et al.* [15] determined the assignments of molecular and lattice vibrations of the  $\text{C}_6\text{Cl}_6$  crystal from the polarized Raman spectra. They assigned the six bands at 56, 54, 45, 38, 25, and  $21\text{ cm}^{-1}$  at 300 K to the lattice vibrations about the three principal axes of a molecule. X-ray diffraction studies [6, 7, 9, 10, 11] indicate that  $\text{C}_6\text{Cl}_6$ ,  $\text{C}_6\text{Br}_6$ , and  $\text{C}_6\text{I}_6$  crystals have isomorphic structures with space group  $P21/c$  and that each contains two molecules in a unit cell with the same packing. Therefore, the patterns of lattice vibrations in  $\text{C}_6\text{Br}_6$  and  $\text{C}_6\text{I}_6$  are expected to be parallel to that of  $\text{C}_6\text{Cl}_6$ . As intermolecular forces differ hardly in  $\text{C}_6\text{Cl}_6$  and in  $\text{C}_6\text{Br}_6$ , the librational frequencies are determined mainly by the principal moments of inertia. The lower librational frequencies, therefore, are expected for  $\text{C}_6\text{Br}_6$ . This consideration leads us to the assignment that four sharp bands at 47, 41, 22, and  $20\text{ cm}^{-1}$  as well as a broad band at  $34\text{ cm}^{-1}$  are due to the librational modes. The band at  $68\text{ cm}^{-1}$  is assigned to an intramolecular vibration, namely the  $b_{2g}$  CX out-of-plane bending mode which corresponds to the  $\text{C}_6\text{Cl}_6$  band at  $103\text{ cm}^{-1}$  [15].

The corresponding  $b_{2g}$  band for  $\text{C}_6\text{I}_6$  as well as the librational bands located at lower frequencies

were not observed. This is probably due to the different experimental conditions used for  $\text{C}_6\text{Br}_6$  and  $\text{C}_6\text{I}_6$ .

Another  $b_{2g}$  fundamental is harder to assign. Normal coordinate analysis indicates that this is a chair type deformation and halogen atoms scarcely move in this mode. The frequency shifts from  $\text{C}_6\text{Cl}_6$  to  $\text{C}_6\text{I}_6$  are expected to be small. BATES [15] assigned a band at  $678\text{ cm}^{-1}$  of  $\text{C}_6\text{Cl}_6$  to this mode. Therefore, the doublet located at 648 and  $637\text{ cm}^{-1}$  of  $\text{C}_6\text{Br}_6$  and the  $\text{C}_6\text{I}_6$  band at  $592\text{ cm}^{-1}$  may be assigned to this fundamental. This assignment is less unambiguous, since overtones of the  $e_{1g}$  doublet at 326 and  $319\text{ cm}^{-1}$  for  $\text{C}_6\text{Br}_6$  and/or the  $\text{C}_6\text{I}_6$  band at  $295\text{ cm}^{-1}$  may also appear in the same frequency range.

In other spectral regions, the present assignment is identical to that given in [3].

#### Infrared spectra

The i.r. spectra in Figs. 1 and 2 are essentially identical to those obtained by DELORME [1] above  $600\text{ cm}^{-1}$ . The spectra also match with the far i.r. spectra of MARSAULT [2] in the region of  $200\text{--}15\text{ cm}^{-1}$ . There are, however, considerable discrepancies between our spectra and those of Delorme below  $500\text{ cm}^{-1}$ . We observed a band at  $246\text{ cm}^{-1}$ ,

Table 2. Assignment of fundamental frequencies of  $\text{C}_6\text{X}_6$  molecules (in  $\text{cm}^{-1}$ )

	Sym. species $D_{6h}$	$\text{C}_6\text{Cl}_6$		$\text{C}_6\text{Br}_6$		$\text{C}_6\text{I}_6$		Sym. species $D_{3d}$	
		Ref. [15] obs.	Ref. [12] calc.	This work obs.	This work calc.	This work obs.	This work calc.		
$a_{1g}$	$\left\{ \begin{array}{l} \nu_1 \\ \nu_2 \end{array} \right.$	1226	1165	1159	1130	1057	1039	$\left. \begin{array}{l} \nu_1 \\ \nu_2 \end{array} \right\}$	$a_{1g}$
		372	383	232	243	166	177		
$b_{2g}$	$\left\{ \begin{array}{l} \nu_7 \\ \nu_8 \end{array} \right.$	678	677	$\left\{ \begin{array}{l} 648 \\ 637 \end{array} \right.$	648	592	601	$\left. \begin{array}{l} \nu_3 \\ \nu_4 \end{array} \right\}$	$a_{1g}$
		103	98	68	68	59	48		
$e_{1g}$	$\left\{ \begin{array}{l} \nu_{11} \\ \nu_{15} \end{array} \right.$	345	389	$\left\{ \begin{array}{l} 326 \\ 319 \end{array} \right.$	300	295	290	$\left. \begin{array}{l} \nu_{13} \\ \nu_{11} \end{array} \right\}$	$e_g$
		1523	1585	1485	1485	$\left\{ \begin{array}{l} 1430 \\ 1423 \end{array} \right.$	1430		
$e_{2g}$	$\left\{ \begin{array}{l} \nu_{16} \\ \nu_{17} \\ \nu_{18} \end{array} \right.$	883	837	795	785	731	706	$\left. \begin{array}{l} \nu_{12} \\ \nu_{14} \\ \nu_{15} \end{array} \right\}$	$e_g$
		323	303	208	178	148	134		
		218	204	148	130	119	99		
$a_{2u}$	$\left\{ \begin{array}{l} \nu_4 \\ \nu_5 \end{array} \right.$	170	197	145	145	144	144	$\left. \begin{array}{l} \nu_8 \\ \nu_9 \end{array} \right\}$	$a_{2u}$
		1080	1067	1031	1036	968	978		
$b_{1u}$	$\left\{ \begin{array}{l} \nu_6 \\ \nu_{12} \end{array} \right.$	395	368	246	216	167	154	$\left. \begin{array}{l} \nu_{10} \\ \nu_{16} \end{array} \right\}$	$a_{2u}$
		1344†	1356	1292	1310	1238	1248		
$e_{1u}$	$\left\{ \begin{array}{l} \nu_{13} \\ \nu_{14} \end{array} \right.$	694†	661	558	580	455	485	$\left. \begin{array}{l} \nu_{16} \\ \nu_{17} \end{array} \right\}$	$e_u$
		215	215	$\left\{ \begin{array}{l} 162 \\ 155 \end{array} \right.$	141	113	108		
$e_{2u}$	$\left\{ \begin{array}{l} \nu_{19} \\ \nu_{20} \end{array} \right.$	521	581	567	560	543	555	$\left. \begin{array}{l} \nu_{19} \\ \nu_{20} \end{array} \right\}$	$e_u$
		79	79	$\left\{ \begin{array}{l} 48.8^* \\ 46.4^* \end{array} \right.$	46	$\left\{ \begin{array}{l} 45.8^* \\ 41.6^* \end{array} \right.$	34		

† Ref. [18]. \* Ref. [3].

a weak doublet at 162 and 155  $\text{cm}^{-1}$ , and a strong band at 145  $\text{cm}^{-1}$  in the spectrum of  $\text{C}_6\text{Br}_6$ , while DELORME *et al.* [1] reported bands at 433, 246, 145, and 135  $\text{cm}^{-1}$ . The last band is missing from the spectrum of Marsault. Reversal of intensity is also observed for the bands at 558 and 526  $\text{cm}^{-1}$ .

Three bands are observed at 167, 144, and 113  $\text{cm}^{-1}$  for  $\text{C}_6\text{I}_6$ . This agrees with the observation by Marsault, while only a band at 113  $\text{cm}^{-1}$  was recorded by Delorme. Below 600  $\text{cm}^{-1}$  we used the combined data of Marsault's and our own.

From the normal coordinate analysis of  $\text{C}_6\text{Br}_6$  two bands at 1292 and 558  $\text{cm}^{-1}$  and a weak doublet at 162 and 155  $\text{cm}^{-1}$  are assigned, respectively, to the ring distortion, the CBr stretching, and the CBr deformation vibrations in the  $e_{1u}$  species. The strong band at 145  $\text{cm}^{-1}$  is assigned to the  $a_{2u}$  CBr bending vibration. Preliminary calculation indicates that two  $e_{2u}$  and two  $b_{1u}$  vibrations have frequencies around 561, 46, 1142, and 230  $\text{cm}^{-1}$ , respectively. These modes are i.r. active in  $D_{3d}$  symmetry, and the observed band at 565  $\text{cm}^{-1}$  is assigned to the  $e_{2u}$  boat type deformation, while the bands at 1031 and 246  $\text{cm}^{-1}$  are assigned to the  $b_{1u}$  triangular deformation and CBr stretching modes respectively. The doublet observed at 48.8 and 46.4  $\text{cm}^{-1}$  by Marsault is assigned to the  $e_{2u}$  CBr bending mode.

Normal coordinate treatment also indicates that the displacements of halogen atoms in the  $a_{2u}$  vibration are very small,\* while those in the  $e_{1u}$  CX deformation vibrations are comparable to the displacements of carbon atoms. The latter mode is expected to be much more sensitive to the halogen substitutions. The band at 144  $\text{cm}^{-1}$  is assigned to the  $a_{2u}$  vibration for  $\text{C}_6\text{I}_6$  (which corresponds to the 145  $\text{cm}^{-1}$  band in  $\text{C}_6\text{Br}_6$ ) and the band at 113  $\text{cm}^{-1}$  to the  $e_{1u}$  mode. The bands at 1238 and 455  $\text{cm}^{-1}$  are assigned to the other  $e_{1u}$  modes, i.e. the ring distortion and the CI stretching vibrations. The weak band at 229  $\text{cm}^{-1}$  may arise from a combination of 113  $\text{cm}^{-1}$  ( $e_{1u}$ ) and 119  $\text{cm}^{-1}$  ( $e_{2g}$ ) bands. The above assignments differ considerably from earlier workers.

The preliminary calculation also indicates that the observed bands at 968, 167, and 543  $\text{cm}^{-1}$  are assigned, respectively, to the triangular deformation and the CI stretching vibrations in the  $b_{1u}$  species and to the boat type deformation in the  $e_{2u}$  species. It is very hard to assign the 167  $\text{cm}^{-1}$  band to the  $e_{1u}$  mode as indicated by Marsault. The calculation shows that the doublet at 45.8 and 41.6  $\text{cm}^{-1}$  observed by MARSAULT [2] is due to the  $e_{2u}$  CX

bending mode. The final assignments are listed in Table 2 along with the assignments for  $\text{C}_6\text{Cl}_6$  [15].

## DISCUSSION

Since the assignments for the fundamental frequencies were settled, the least-squares refinement of force constants has been attempted to reproduce the frequencies in the  $a_{1g}$ ,  $e_{2g}$ ,  $e_{1u}$ ,  $a_{2u}$ , and  $b_{2g}$  species as nearly as possible. The force constants  $K_{cx}$ ,  $K_{cc}$ ,  $H_{cx}$ ,  $H_{cc}$  as well as  $c_{xx}$ , a cis interaction constant introduced by MANN *et al.* [16, 17] in the normal coordinate analysis in  $\text{C}_2\text{X}_4$ , are treated as adjustable parameters in the in-plane modes. As for out-of-plane force constants,  $A$ ,  $B$ ,  $a_0$ ,  $a_m$ ,  $a_p$ ,  $b_0$ ,  $b_m$ , and  $b_p$  are taken as variables with the constraints of  $\Delta A = \Delta a_m$ ,  $\Delta a_0 = \Delta a_p$ ,  $\Delta B = \Delta b_m = -\Delta b_0 = -\Delta b_p$ . The constraints leave the  $e_{2u}$  and  $e_{1g}$  frequencies unchanged: the frequency fits in these species are good in the preliminary calculation. The final results are given in Table 2 and the final set of force constants is given in Table 3.

The overall agreements of the calculated and observed frequencies are quite satisfactory. It is encouraging that the frequency fits in the  $b_{1u}$  species are improved considerably in the refinement process which supports the present revised assignments. Some frequencies are interpreted as combination bands, and they are listed in Table 4.

The obtained value of  $K_{\text{CBr}}$  for  $\text{C}_6\text{Br}_6$  is reasonable considering the corresponding value for fluorobromobenzenes is reported to be 2.9314  $\text{md}/\text{A}$  [19]. It is in-

Table 3. Values of force constants (final set)

	In-plane f. c. (md/A)		Out-of-plane f. c. (md·A)		
	$\text{C}_6\text{Br}_6$	$\text{C}_6\text{I}_6$	$\text{C}_6\text{Br}_6$	$\text{C}_6\text{I}_6$	
$K_{cc}$	4.699	4.313	$A$	0.4082	0.4129
$K_{cx}$	2.578	1.656	$B$	0.4064	0.4409
$H_{cx}$	0.260	0.239	$C$	0.0398	0.0398
$F_{cx}$	0.387	0.387	$D$	-0.0027	-0.0027
$H_{ccc}$	0.098	0.115	$a_0$	0.0696	0.1016
$F_{ccc}$	0.241	0.241	$a_m$	-0.0018	0.0029
$o_{cc}$	0.481	0.481	$a_p$	0.0591	0.0912
$m_{cc}$	-0.204	-0.204	$b_0$	-0.0258	-0.0603
$p_{cc}$	0.377	0.377	$b_m$	0.0680	0.1025
$o_{cx}$	0.024	0.024	$b_p$	-0.0396	-0.0741
$m_{cx}$	0.0	0.0			
$p_{cx}$	-0.006	-0.006			
$o_{cex}$	0.005	0.005			
$m_{cex}$	-0.014	-0.014			
$p_{cex}$	0.007	0.007			
$o_{ccc}$	-0.018	-0.018			
$m_{ccc}$	0.0	0.0			
$p_{ccc}$	0.0	0.0			
$c_{xx}$	0.373	0.423			

\* 0.14 for  $\text{C}_6\text{Br}_6$  and 0.09 for  $\text{C}_6\text{I}_6$  relative to the displacements of carbon atoms.

Table 4. Assignment of combination bands of  $C_6Br_6$  and  $C_6I_6$  (in  $cm^{-1}$ )

$C_6Br_6$				
	obs.	assignment	obs.	assignment
i.r.	447 (2)	208 ( $e_{2g}$ ) + 246 ( $b_{1u}$ )	1249 (8)	1485 ( $e_{2g}$ ) - 246 ( $b_{1u}$ )
	630 (3)	68 ( $b_{2g}$ ) + 567 ( $e_{2u}$ )	1336 (4)	1485 ( $e_{2g}$ ) - 145 ( $a_{2u}$ )
	677 (3)	637 ( $b_{2g}$ ) + 48.8 ( $e_{2u}$ )	1520 (2)	1292 ( $e_{1u}$ ) + 232 ( $a_{1g}$ )
	865 (2)	319 ( $e_{1g}$ ) + 558 ( $e_{1u}$ )	1952 (2)	1292 ( $e_{1u}$ ) + 648 ( $b_{2g}$ )
	880 (5)	1031 ( $b_{1u}$ ) - 148 ( $e_{2g}$ )		
	965 (2)	1031 ( $b_{1u}$ ) - 68 ( $b_{2g}$ )	R	
	1057 (2)	1292 ( $e_{1u}$ ) - 232 ( $a_{1g}$ )	819 (1)	246 ( $b_{1u}$ ) + 567 ( $e_{2u}$ )
	1098 (2)	68 ( $b_{2g}$ ) + 1031 ( $b_{1u}$ )	1215 (1)	1159 ( $a_{1g}$ ) + 68 ( $b_{2g}$ )
$C_6I_6$				
	obs.	assignment	obs.	assignment
i.r.	229 (4)	119 ( $e_{2g}$ ) + 113 ( $e_{1u}$ )	702 (4)	166 ( $a_{1g}$ ) + 543 ( $e_{2u}$ )
	249 (3)	295 ( $e_{1g}$ ) - 45.8 ( $e_{2u}$ )	773 (4)	731 ( $e_{2g}$ ) + 41.6 ( $e_{2u}$ )
	267 (3)	119 ( $e_{2g}$ ) + 144 ( $a_{2u}$ )	834 (5)	295 ( $e_{1g}$ ) + 543 ( $e_{2u}$ )
	280 (2)	{ 166 ( $a_{1g}$ ) + 113 ( $e_{1u}$ ) 119 ( $e_{2g}$ ) + 167 ( $b_{1u}$ )	862 (4)	968 ( $b_{1u}$ ) - 119 ( $e_{2g}$ )
	315 (4)	{ 166 ( $a_{1g}$ ) + 144 ( $a_{2u}$ ) 148 ( $e_{2g}$ ) + 167 ( $b_{1u}$ )	897 (2)	731 ( $e_{2g}$ ) + 167 ( $b_{1u}$ )
	339 (2)	455 ( $e_{1u}$ ) - 119 ( $e_{2g}$ )	921 (2)	1057 ( $a_{1g}$ ) - 144 ( $a_{2u}$ )
	385 (4)	543 ( $e_{2u}$ ) - 166 ( $a_{1g}$ )	980 (3)	1430 ( $e_{2g}$ ) - 455 ( $e_{1u}$ )
	440 (4)	295 ( $e_{1g}$ ) + 144 ( $a_{2u}$ )	1009 (3)	1057 ( $a_{1g}$ ) - 45.8 ( $e_{2u}$ )
	491 (4)	543 ( $e_{2u}$ ) - 59 ( $b_{2g}$ )	1074 (3)	1238 ( $e_{1u}$ ) - 166 ( $a_{1g}$ )
	564 (4)	731 ( $e_{2g}$ ) - 167 ( $b_{1u}$ )	1089 (3)	968 ( $b_{1u}$ ) + 119 ( $e_{2g}$ )
	574 (3)	119 ( $e_{2g}$ ) + 455 ( $e_{1u}$ )	1119 (4)	1238 ( $e_{1u}$ ) - 119 ( $e_{2g}$ )
	609 (2)	148 ( $e_{2g}$ ) + 455 ( $e_{1u}$ )	1171 (2)	1057 ( $a_{1g}$ ) + 113 ( $e_{1u}$ )
	663 (2)	119 ( $e_{2g}$ ) + 543 ( $e_{2u}$ )	1205 (8)	1057 ( $a_{1g}$ ) + 144 ( $a_{2u}$ )
	676 (3)	968 ( $b_{1u}$ ) - 295 ( $e_{1g}$ )	1316 (3)	1430 ( $e_{2g}$ ) - 113 ( $e_{1u}$ )
	692 (2)	148 ( $e_{2g}$ ) + 543 ( $e_{2u}$ )	R	
			1349 (1)	1057 ( $a_{1g}$ ) + 295 ( $e_{1g}$ )

Numbers in the parentheses indicate the relative intensities.

interesting to note that the values of  $K_{CC}$  for  $C_6Br_6$  and  $C_6I_6$  are smaller than the value for  $C_6H_6$  (5.237 md/A) [13], and that the force constant  $c_{xx}$ , which was completely ignored in  $C_6H_6$ , has rather large values for  $C_6Br_6$  and  $C_6I_6$ . As for out-of-plane force constants  $A$  and  $B$ , they are slightly larger than those found for  $C_6H_6$  ( $A = 0.3970$  mdA,  $B = 0.3668$  mdA) [13]. In order to make more quantitative discussion, it is necessary to study a complete set of brominated and iodinated benzenes, as were attempted for that of fluorinated benzenes [20, 21] and that of chlorinated benzenes [22].

In conclusion, all the observed i.r. and Raman bands are assigned to fundamentals as well as combination bands on the basis of normal coordinate analysis assuming a planar  $D_{6h}$  configuration. The present assignments for the  $b_{2g}$ ,  $a_{2u}$ ,  $b_{1u}$ ,  $e_{1u}$ , and  $e_{2u}$  modes differ from those given earlier [1-3]. A few bands can be interpreted as due to i.r. inactive ( $b_{1u}$ ,  $e_{2u}$ ) or Raman inactive ( $b_{2g}$ ) bands in  $D_{6h}$  symmetry. This supports a slightly puckered molecular model of  $D_{3d}$  symmetry for  $C_6Br_6$  and  $C_6I_6$  crystals as were proposed by previous investigators [1-3].

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