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## COMMUNICATION

## Cooperative spin transition in a lipid layer like system†

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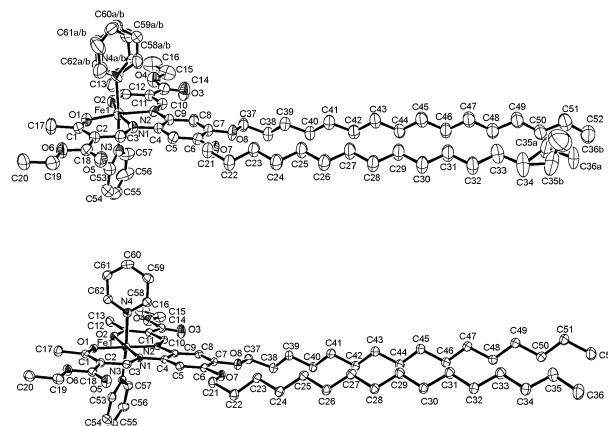
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A novel iron(II) mononuclear spin transition complex [FeL(py)<sub>2</sub>] displays an abrupt spin transition around 225 K accompanied by a very wide thermal hysteresis loop (~50 K) that spreads out over 100 K. Crystal structure analysis in both low-spin and high-spin states reveals a lipid layer-like arrangement of the complex molecules and provides insights into the spin switching mechanism.

Since the discovery of spin crossover (SCO) compounds in 1931 by Cambi *et al.*<sup>1</sup> the interest in this substance class never vanished,<sup>2</sup> as the thermochromism associated with the spin state change makes them potentially useful for various applications such as display and memory device units,<sup>3</sup> sensors<sup>4</sup> and cold channel control units in food and medical sectors.<sup>5</sup> In order to realise application it is important to explore different possibilities for the nanostructuration of SCO materials<sup>6</sup> and to investigate if additional properties can be combined (*e.g.* liquid crystal behaviour,<sup>7</sup> magnetic exchange interactions<sup>8</sup>) resulting in multifunctional SCO materials.<sup>9</sup> In this frame, we modified Schiff base like ligands used for the synthesis of SCO complexes<sup>10</sup> by adding long alkyl chains in the outer periphery by preparing [FeL(py)<sub>2</sub>] (**1**) with L = (*E,E*)-[diethyl-2,2'-[4,5-dihexadecyloxy-1,2-phenylenebis(iminomethylidene)]bis(3-oxobutanato)]. We aimed to study the influence of this modification on the crystal packing as well as on the SCO behaviour and investigate if additional features could be achieved for **1**.

A complete description of the synthesis of H<sub>2</sub>L and the iron complex (**1**) is given in the ESI.† Single crystals suitable for an X-ray analysis of 1·0.25H<sub>2</sub>O were obtained and the crystal structure was determined first at 250 K and then at 125 K (same crystal, in both cases space group *P1*), corresponding to the high-spin (HS) and low-spin (LS) states of the complex as seen in the magnetic measurement (Fig. 3). The crystallographic data are summarized in the ESI†, Table S1. Fig. 1 displays an ORTEP drawing of the asymmetric unit of **1** in the HS and the LS state. An excerpt of the coordination



**Fig. 1** ORTEP drawing of the asymmetric unit of **1** in the HS (top) and LS (bottom) states. Hydrogen atoms and the water molecule have been omitted for clarity. Displacement ellipsoids are shown at the 50% probability level.

**Table 1** Selected bond lengths [Å] and angles [°] within the inner coordination sphere of **1** in the HS and LS states

	Fe–N <sub>eq</sub>	Fe–O <sub>eq</sub>	Fe–N(L <sub>ax</sub> )	O–Fe–O	L <sub>ax</sub> –Fe–L <sub>ax</sub>	∠ L <sub>ax</sub> <sup>b</sup>
HS	2.059(3) 2.086(2)	2.001(2) 1.999(2)	2.284(3) 2.280(4) <sup>a</sup> 2.288(7) <sup>a</sup>	106.10(9)	173.73(12) <sup>a</sup> 176.95(23) <sup>a</sup>	21.5 <sup>a</sup> 47.5 <sup>a</sup>
LS	1.897(2) 1.907(2)	1.935(2) 1.947(2)	2.021(2) 2.014(2)	88.80(7)	175.06(8)	83.6

<sup>a</sup> Disorder. <sup>b</sup> Angle between the pyridine planes.

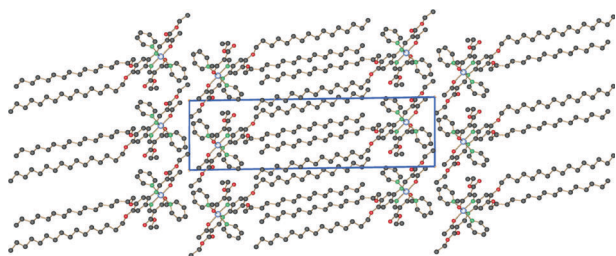
environment is given in Fig. S1 (ESI†). Selected bond lengths and angles around the inner coordination sphere of the iron centre are summarized in Table 1.

The average bond lengths within the first coordination sphere of the iron(II) centres in the HS structure are 2.07 Å (Fe–N<sub>eq</sub>), 2.00 Å (Fe–O<sub>eq</sub>) and 2.28 Å (Fe–L<sub>ax</sub>). Those and the observed O–Fe–O angle (106°) are in the region expected for HS complexes of this ligand type.<sup>10,11</sup> Upon the HS to LS transition a shortening of the bond lengths of about 10% is observed, as observed for other iron(II) SCO complexes.<sup>2,11</sup> The average bond lengths in the LS-structure are 1.90 Å (Fe–N<sub>eq</sub>), 1.94 Å (Fe–O<sub>eq</sub>) and 2.02 Å (Fe–L<sub>ax</sub>) with an O–Fe–O angle of 89°. In the HS state, a disorder is observed of one pyridine and at the end of one of the C<sup>16</sup> alkyl chains. The pyridine ring including N4 is contorted in two directions in a relative ratio of 60:40, the same ratio is observed for the ethyl endgroup (C35 and C36) in the alkyl chain bound by O7. The planes spanned by the two axial pyridine rings are

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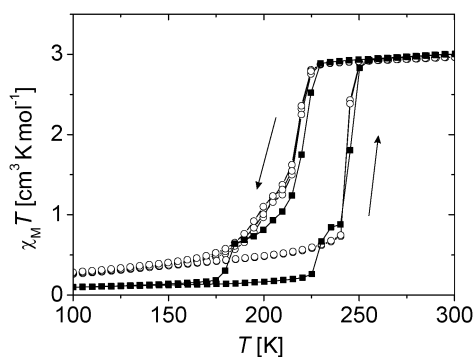
**Fig. 2** Packing of the molecules of **1**·0.25 H<sub>2</sub>O in the crystal projected along [1 0 0] in the HS state.

staggered in the HS state. Upon cooling the disordered pyridine ring changes its orientation resulting in a nearly perpendicular arrangement in the LS state. An additional water molecule is observed in the crystal packing with an approximate occupation number of 0.25. Several hydrogen bonds and short contacts are observed between the complex molecules and the water (Table S2, ESI†).

In the crystal the molecules are packed in a lipid layer like arrangement as illustrated in Fig. 2 with the layers running along the *a*–*b*-plane. Within one layer, the alkyl chains of the Schiff base like ligand are packed in the middle and the SCO centres are on the outer sides.

In the LS state a network of hydrogen bonds is formed between the SCO centres between the layers and within the layer, the additional water molecule being also involved, as illustrated in Fig. S2 (ESI†). In the HS state the rearrangement of the pyridine ring and the alkyl chain results in changes for the short contacts. Now, short contacts are only observed between the molecules of one lipid like layer but not between the layers as illustrated in Fig. S2 and S3 (ESI†).

Magnetic susceptibility data of **1**·0.25 H<sub>2</sub>O were recorded on cooling and warming over the temperature range 300–5 K (Fig. 3). At room temperature the  $\chi_M T$  product is 3.0 cm<sup>3</sup> K mol<sup>−1</sup> which is in the range expected for an iron(II) complex in the HS state. In the first cycle (squares) upon cooling the magnetic moment remains constant until 229 K where an abrupt spin transition (ST) takes place with about 60% of the molecules and  $T_c^{(60)\downarrow} = 222$  K. In the temperature range between 214 K and 185 K the magnetic moment decreases gradually from  $\chi_M T$  (214 K) = 1.1 cm<sup>3</sup> K mol<sup>−1</sup> to  $\chi_M T$  (185 K) = 0.6 cm<sup>3</sup> K mol<sup>−1</sup>. Below 185 K a second abrupt step is observed involving the remaining 20% of the molecules  $T_c^{(20)\downarrow} = 182$  K. At 100 K the remaining  $\chi_M T$  product is 0.1 cm<sup>3</sup> K mol<sup>−1</sup> which is in the range expected for

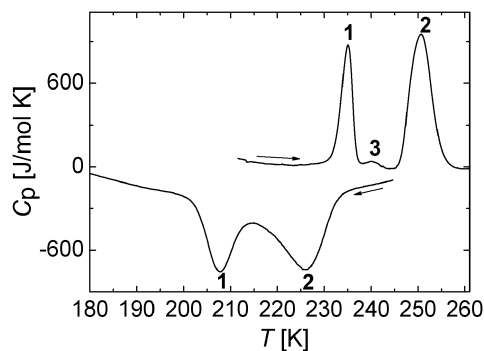


**Fig. 3**  $\chi_M T$  vs. *T* plot of **1** of the first cycle (squares) and all the following cycles (open circles). Displayed in the temperature range 300–100 K.

an iron(II) complex in the LS state. Upon heating the ST occurs in two steps with about 30% of the molecules involved in the first ( $T_c^{(30)\uparrow} = 229$  K) and about 70% of the molecules involved in the second step ( $T_c^{(70)\uparrow} = 245$  K). The width of the thermal hysteresis loop is 23 K for the first and 47 K for the second step. The second and third thermal cycles reveal the absence of the second gradual step and the remaining  $\chi_M T$  product is 0.3 cm<sup>3</sup> K mol<sup>−1</sup> at 100 K indicating an iron(II) complex almost completely in the LS state.

The sample has been studied by differential scanning calorimetry (DSC) over the temperature range 300–98 K in cooling and warming modes, at 10 K min<sup>−1</sup>, in order to extract thermodynamical parameters associated with the spin transition, and probe the order of the phase transitions. On warming from 98 K, two major endothermic peaks corresponding to a first order phase transition proceeding in two steps were detected (Fig. 4). The first one is observed at  $T_{\max}^{(1)\uparrow} = 235(1)$  K and a broader peak is found at  $T_{\max}^{(2)\uparrow} = 250(1)$  K. These data match the transition temperatures derived from SQUID measurements (see Fig. 3). A very less intense peak, whose shape indicates a continuous or weakly first-order phase transition, is observed at  $T^{(3)} = 240$  K corresponding to the plateau region between the two phase transitions. On cooling from room temperature, broader exothermic peaks arise at different temperatures, confirming the hysteretic character of the ST process. A peak is found at  $T_{\max}^{(2)\downarrow} = 226(1)$  K, in agreement with the first ST branch, which is followed by a second peak at  $T_{\max}^{(1)\downarrow} = 207(1)$  K, which only corresponds to the onset of the plateau region of the magnetic curve. The tiny second step in the cooling mode of the first cycle of the SQUID measurements around 182 K is not seen by DSC. Such a DSC profile was confirmed by two successive cooling and heating cycles. Interestingly, two peaks are still observed in the warming and cooling process after the first cycling. One is clearly related to the spin transition (2) whereas the other peak (1) relates to a thermal anomaly that plays a role in the spin state during the first cycle but not later, still existing but not affecting it. Enthalpy and entropy associated with these peaks are given in Table S3 (ESI†).

This unprecedented ST behaviour can be explained thanks to the results from X-ray structure analysis. Step-wise spin transitions are often related to the presence of two or more inequivalent iron centres.<sup>12</sup> In our case, the disorder of the pyridine ring in the HS state could be responsible for inequivalent iron centres and thus the steps in the transition curve. Such a



**Fig. 4** Heat capacity vs. *T* of **1** over the temperature range 261–180 K at a scanning rate of 10 K min<sup>−1</sup> in the cooling (←) and warming (→) modes.

situation was recently discussed by Matouzenko *et al.*<sup>13</sup> with the same symmetries in the HS and LS states and a symmetry breaking on the plateau.<sup>14</sup> As different intermolecular interactions are observed for the disordered parts of the HS state (Fig. S2, ESI†) the differences in the hysteresis width can be explained with differences in the H-bond network. It should be noted that the space group does not change upon ST and thus the observed hysteresis cannot be related to a structural phase transition but must be related to other cooperative effects. Due to the long alkyl chains in the outer periphery of the complex the observed change in the cell volume ( $\Delta V/V = 2.9\%$ ) is very small, especially when the contribution from the thermal contraction is considered. Thus the hysteresis cannot solely be related to elastic interactions. The changes in the hydrogen bond network are one possible explanation for the cooperative spin transition.<sup>15</sup> A very interesting feature concerns the second peak in the DSC measurements that is related to the gradual part in the ST curve. This thermal anomaly could result from an order–disorder transition of the pyridine ring,<sup>16</sup> which is disordered in the HS state and that orders in the LS, as detected by X-ray diffraction. Such types of transitions are known to be able to control the course of a ST leading to gradual regions.<sup>17</sup> The disappearance of the small step after the first cooling/heating cycle could be related to small changes in the molecule structure or the intermolecular interactions due to this order–disorder transition.

The introduction of long alkyl chains in the outer periphery of Schiff base like ligands did lead to head–tail molecules that crystallise in a lipid layer like structure and show a cooperative spin transition with a wide thermal hysteresis loop (up to  $\sim 50$  K). The X-ray structure solved in the HS and LS states allows a deeper insight into the ST mechanism of **1** which results from an interplay involving the H-bond network and order–disorder transition of the pyridine rings.

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