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Controlling Destructive Quantum Interference in Tunneling Junctions Comprising Self-assembled Monolayers via Bond Topology and Functional Groups

Yanxi Zhang,^{†,‡,¶} Gang Ye,^{†,‡,¶} Saurabh Soni,^{†,‡} Xinkai Qiu,^{†,‡} Theodorus L. Krijger,^{†,‡} Harry T. Jonkman,[‡] Marco Carlotti,^{†,‡} Eric Sauter,[§] Michael Zharnikov,[§] and Ryan C. Chiechi*,^{†,‡}

†Stratingh Institute for Chemistry, University of Groningen, Nijenborgh 4, 9747 AG Groningen, The Netherlands

‡Zernike Institute for Advanced Materials, Nijenborgh 4, 9747 AG Groningen, The

Netherlands

¶ These authors contributed equally to this work.

§Applied Physical Chemistry, Heidelberg University, Im Neuenheier Feld 253, Heidelberg 69120, Germany

E-mail: r.c.chiechi@rug.nl

ABSTRACT

Quantum interference effects (QI) are of interest in nano-scale devices based on molecular tunneling junctions because they can affect conductance exponentially through minor structural changes. However, their utilization requires the prediction and deterministic control over the position and magnitude of QI features, which remains a significant challenge. In this context, we designed and synthesized three benzodithiophenes based molecular wires; one linearly-conjugated, one cross-conjugated and one cross-conjugated quinone. Using eutectic Ga-In (EGaIn) and CP-AFM, we compared them to a well-known anthraquinone in molecular junctions comprising self-assembled monolayers. By combining density functional theory and transition voltage spectroscopy, we show that the presence of an interference feature and its position can be controlled independently by manipulating bond topology and electronegativity. This is the first study to separate these two parameters experimentally, demonstrating that the conductance of a tunneling junction depends on the position and depth of a QI feature, both of which can be controlled synthetically.

INTRODUCTION

Molecular electronics is concerned with the transport of charge through molecules spanning two electrodes, [1] the fabrication of which is a challenging area of nanotechnology. [2-4] In such junctions, π -conjugated molecules influence transport more than a simple, rectangular tunneling barrier; when a tunneling electron traverses the region of space occupied by orbitals localized on these molecules, its wave function can undergo constructive or destructive interference, enhancing or suppressing conductance. When the presence of different pathways in molecular system affects conductance, it is typically described as quantum interference (QI), [5] which was originally adapted from the Aharonov-Bohm effect [6] to substituted benzenes. [7,8] The concept "quantum interference effect transistor" was also proposed using meta-benzene structures for device application. [9] Solomon et al. further refined the concept in the context of Molecular Electronics where it is now well established that destructive QI leads to lower conductance in tunneling junctions. [10-15] We previously demonstrated QI in SAM-based junctions using a series of compounds based on an anthracene core; AC, which is linearly-conjugated; AQ, which is cross-conjugated via a quinone moiety; and AH, in which the conjugation is interrupted by saturated methylene bridges (Figure S1). [16] Subsequent studies verified these findings in a variety of experimental platforms and a consensus emerged that, provided the destructive QI feature (anti-resonances in transmission) is sufficiently close to the Fermi level, $E_{\rm F}$, cross-conjugation leads to QI. [17-24] However, experimental studies on conjugation patterns other than AC/AQ are currently limited to ring substitutions such as meta-substituted phenyl rings, [25–32] or varied connectivities in azulene, [33–35] which differ fundamentally [5,11,36-38] from cross-conjugated bond topologies [23,39,40] because they change tunneling pathways, molecular-lengths and bond topology simultaneously (Table S1). Isolating these variables is however important because the only primary observable is conductance, which varies exponentially with molecular length. More recent work has focused on "gating" OI effects by controlling the alignment of π -systems through-space [37,41,42] and affecting the orbital symmetry of aromatic rings with heteratoms. [43-45] These studies exclusively study the effects of the presence and absence of QI features; to date—and despite recent efforts ^[46]—the specific effects of bond topology and electronegativity on the depth and position of QI features have not been isolated experimentally.

To address this issue, we designed and synthesized the series of benzodithiophene derivatives (BDT-n); benzo[1,2-b:4,5-b']dithiophene (BDT-1, linearly conjugated), benzo[1,2b:4,5-b']dithiophene-4,8-dione (BDT-2, cross-conjugated with quinone), and benzo[1,2-b:5,4b'dithiophene (BDT-3, cross-conjugated and an isomer of BDT-1). These compounds separate the influence of cross-conjugation (bond topology) from that of the electron-withdrawing effects of the quinone functionality while controlling for molecular formula and length. We investigated the charge transport properties of these molecules in tunneling junctions comprising self-assembled monolayers (SAMs), which are relevant for solid-state molecular-electronic devices. [47-49] Through a combination of density functional theory (DFT) and transition voltage spectroscopy (TVS) we show that cross-conjugation produces QI features near occupied molecular states and that the position and depth of the QI feature is strongly influenced by the strongly electron-withdrawing quinone functionality, which places these features near unoccupied states while simultaneously bringing those states close to $E_{\rm F}$. Thus, by controlling bond topology and electronegativity separately, the conductance can be tuned independently of length and connectivity via the relative positions of the QI features and molecular states and not just the presence or absence of such features.

RESULTS AND DISCUSSION

To isolate molecular effects on transport, it is important to control for changes to the width of the tunneling barrier which, in SAMs, is typically defined by the end-to-end lengths of the molecules. Conductance G generally varies exponentially with the barrier-width d such that $G = G_0 \exp(-\beta d)$, where G_0 is the theoretical value of G when d = 0, and β is the tunneling decay coefficient. Since β depends on the positions of molecular states relative to

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 $E_{\rm F}$ and we are comparing compounds with very different redox potentials (orbital energies) we can only ascribe changes to G if d is invariant across the series. Furthermore, to isolate the variable of bond topology experimentally, the electronic properties of the linear- and cross-conjugated compounds must be nearly identical. Figure 1a shows the structures of the BDT-n series and AQ; the "arms" are linearly-conjugated phenylacetylenes (highlighted in the light blue background) and the cores (Ar, highlighted in the brown background) are substituted by the structures indicated. The variation in the end-to-end lengths of these compounds is within 1 Å and the linear- and cross-conjugated compounds BDT-1 and BDT-3 differ only by the relative position of sulfur atoms; they have the same molecular formula. The synthesis, full characterization and a detailed discussion of their properties are provided in the Supplementary Information. Note that we include AQ in the series as a benchmark for destructive QI effects.

We measured tunneling charge transport through metal-molecule-metal junctions comprising BDT-1, BDT-2, BDT-3 and AQ using conformal eutectic Ga-In (EGaIn) contacts as top electrodes.^[50] We utilized an established procedure of the *in situ* deprotection of thioacetates [41,51] to form well-defined SAMs on Au substrates; these substrates served then as bottom electrodes. We refer to the assembled junctions as Au/SAM//EGaIn where "/" and "//" denote a covalent and van der Waals interfaces, respectively. The geometry of the junctions is shown in Figure 1b. To verify that the structural similarities of the compounds carry over into the self-assembly process, we characterized the SAMs of BDT-nby several complementary techniques, including (high-resolution) X-ray photoelectron spectroscopy (HRXPS/XPS) and angle-resolved near-edge X-ray absorption fine structure spectroscopy (NEXAFS). These data are discussed in detail in the Supplementary Information and summarized in Table 1. The characterization of SAMs of AQ is reported elsewhere. [16,41] The XPS and NEXAFS data suggest that the molecules in the BDT-n SAMs are assembled upright with the tilt angle of approximately 35°. The molecules are packed densely on the order of 10¹⁴ molecules per cm² as are similar conjugated molecular-wire compounds. [41]

Figure 2a shows the current-density versus voltage (J/V) curves for the **BDT-**n series and **AQ** using EGaIn top contacts. **BDT-1** is the most conductive across the entire bias window. The conductance of linearly-conjugated BDT-1 and AC (Figure S1; a linearly-conjugated analog of AQ), are almost identical (Figure S21), meaning that the low-bias conductivity and/or values of J are directly comparable between the AC/AQ and BDT-n series. As expected, the cross-conjugated BDT-2, BDT-3 and AQ are all less conductive than BDT-1 (and AC). The low-bias conductivity (from the Ohmic region, $-0.1\,\mathrm{V}$ to $0.1\,\mathrm{V}$) of the quinones (BDT-2 and AQ), however, is even more suppressed than the cross-conjugated BDT-3, while the magnitudes of J for BDT-2, BDT-3 and AQ are similar beyond $-0.5 \,\mathrm{V}$. We observed similar behavior in QI mediated by through-space conjugation in which the compound with an interference feature very close to $E_{\rm F}$ exhibited a sharp rise in J, eventually crossing J/V curve of the compound with a feature further from $E_{\rm F}$. [41] This observation suggests that, as the junction is biased, the transmission probability "climbs" the interference feature rapidly, bringing highly transmissive conduction channels into the bias window at sufficiently low values of V to meet and exceed the total transmission of the compound for which the interference feature is far from $E_{\rm F}$ at zero bias. Further discussion on the asymmetry of J/V curves is included in the Supplementary Information.

To better compare the conductance of the molecules, we calculated the low-bias conductivities and normalized them to **BDT-1**. These values are plotted in Figure 2b, showing that cross-conjugation lowers the conductance of **BDT-3** by an order of magnitude compared to **BDT-1** and the quinone functionality of **BDT-2** and **AQ** lowers it by two orders of magnitude, in agreement with the analogous behavior of **AC** and **AQ**. [20] To control for large-area effects (e.g., if there are defects in the SAM), we measured **BDT-n** series by conducting-probe atomic force microscopy (CP-AFM) with Au electrodes and found the same trend: **BDT-1** > **BDT-3** > **BDT-2**, however, a direct comparison of low-bias conductivities was precluded by the extremely high resistance of **BDT-2** and **AQ** at low bias. These data are discussed in detail in the Supplementary Information. Thus, we conclude

that quinones suppress conductance more than cross-conjugation alone, irrespective of the measurement/device platform.

For insight into the shapes of the J/V curves and the conductance, we simulated the transmission spectra, T(E) vs. $E-E_{\rm F}$ ($E_{\rm F}$ value of $-4.3\,{\rm eV}$, see Experimental Section), of the **BDT-**n series using density functional theory (DFT) and compared the resulting curves with AQ (Figure 3). These calculations, which are discussed in more detail in the Computational Methodology section of the Supplementary Information, simulate the transmission spectra through isolated molecules in vacuum at zero bias and are useful for predicting trends in conductance. There are three important features of these curves: 1) Only the compounds with cross-conjugation (including quinones) show sharp dips (anti-resonances or QI features) [13,18] in the frontier orbital gap; 2) the dips occur near $E_{\rm F}$ only for the two quinones; and 3) the QI features are more pronounced for the molecules in which the cross-conjugation is caused by a quinone moiety as opposed to the carbon-carbon bond topology. When bias is applied to a junction, the x-axis of the transmission plot shifts and $E_{\rm F}$ broadens such that an integral starting at $E-E_{\rm F}=0\,{\rm eV}$ and widening to larger ranges of $E-E_{\rm F}$ is a rough approximation of how T(E) translates into current, I(V). This relationship is apparent in the slightly lower conductance of AQ compared to BDT-2 (Figure 2b) and the slightly lower values of T(E) for AQ compared to BDT-2 across the entire range of $E-E_{\rm F}$. The proximities of the QI features to $E_{\rm F}$ is also apparent in the J/V curves (Figure 2a). As the junction is biased, the minimum of the QI feature shifts such that, by 0.5 V, the transmission probabilities are roughly equal for BDT-n and AQ.

The shape of T(E) near $E - E_{\rm F} = 0\,{\rm eV}$ is roughly traced by differential conductance plots of $\log |\frac{dJ}{dV}|$ vs. V, allowing QI features near $E_{\rm F}$ to be resolved experimentally. [18,41,52] Figure 4 shows heatmap plots of differential conductance of Au/SAM//EGaIn constructed from histograms binned to $\log |\frac{dJ}{dV}|$ for each value of V. (Note that these are histograms of J/V curves with no data-selection, thus, brighter colors correspond to mean values of J and are not related to conductance histograms of single-molecule break-junctions; see

Supplementary Information for details.) Both BDT-1 and BDT-3 exhibit ordinary, Ushaped plots characteristic of non-resonant tunneling. By contrast, both AQ and BDT-2—the two compounds bearing quinone functionality—show V-shaped plots with negative curvature. These results are in agreement with Figure 3, which places the QI features for the quinone moieties, \mathbf{AQ} and \mathbf{BDT} -2, much closer to E_{F} than for \mathbf{BDT} -3. The positions of these features are related to the positions of highest-occupied and lowest-unoccupied π -states (HOPS and LUPS), which are in good agreement between DFT and experiment (Tables S2 and S3). Thus, the differential conductance heatmaps (experiment) and DFT (simulation) both indicate that cross-conjugation suppresses conductance because it creates a dip in T(E) in the frontier orbital gap, but that the electron-withdrawing nature of the quinone functionality simultaneously pulls the LUPS and the interference features close to $E_{\rm F}$ such that the J/V characteristics and transmission plots of ${\bf AQ}$ and ${\bf BDT}$ -2 are nearly indistinguishable despite the presence of two thienyl groups in BDT-2. These results also suggest that tunneling transport is mediated by the HOPS (hole-assisted tunneling) for BDT-1 and BDT-3 and by the LUPS (electron-assisted tunneling) for BDT-2 and AQ because tunneling current is dominated by the resonance(s) closest to $E_{\rm F}$.

To further investigate the mechanism of transport, we measured transition voltages, V_{trans} (Table S3, Figures S17 and S18), which provide information about the energy offset between E_{F} and the dominant frontier orbital. [53,54] Figure 5a shows the levels for the **BDT**-n series calculated by DFT with respect to E_{F} (-4.3 eV), clearly predicting LUPS-mediated tunneling for **BDT**-2 and **AQ**. Figure 5b compares the experimental values of V_{trans} to the energy differences between E_{F} and the frontier orbitals. The salient feature of Figure 5b is that the trend in $|E_{\text{HOPS}} - E_{\text{F}}|$ opposes the trend in V_{trans} such that the trend in experimental values of V_{trans} agrees with DFT only when we compare V_{trans} with $|E_{\text{HOPS}} - E_{\text{F}}|$ for **BDT**-1 and **BDT**-3, and with $|E_{\text{LUPS}} - E_{\text{F}}|$ for **BDT**-2 and **AQ**. Thus, DFT calculations combined with experimental values of V_{trans} predict electron-assisted tunneling for **BDT**-2 and **AQ**. This degree of internal consistency between the experiment and theory is important

because, ultimately, the only primary observable is conductance, which we plot as J/V curves, differential conductance heatmaps and Fowler-Nordheim plots (from which we extract V_{trans}). And we find remarkable agreement between these direct and indirect observations and DFT calculations on model junctions comprising single molecules.

CONCLUSION

The key question of this work is how cross-conjugation and electronegativity affect QI features. [11,20,52,55,56] Based on our experimental observations and calculations, we assert that destructive QI induced by cross-conjugation is highly sensitive to the functional groups that induce the cross-conjugation and that quinones are, therefore, a poor testbed for tuning QI effects (beyond switching them on and off^[57]) because their strong electron-withdrawing nature places a deep, destructive feature near $E_{\rm F}$ irrespective of other functional groups (in our case, two fused thiophene rings barely make a difference). Comparing a quinone to a hydrocarbon also compares HOPS-mediated tunneling to LUPS-mediated tunneling between molecules with significantly different band-gaps and absolute frontier orbital energies. In contrast, BDT-1 and BDT-3 are heterocyclic isomers with no functional groups, identical molecular formulas, nearly-identical HOPS, identical lengths that translate into SAMs of identical thicknesses, and transport is dominated by the HOPS. They isolate the single variable of conjugation patterns, allowing us to separate bond topology (cross-conjugation) from electronic properties (functional groups), giving experimental and theoretical insight into the relationship between bond topology and QI. Our results suggest that there is a lot of room to tune the conductance of moieties derived from BDT-3 by including pendant groups (e.g., halogens, CF₃ groups or acidic/basic sites) that shift the QI feature gradually towards $E_{\rm F}$ synthetically and/or in response to chemical signals.

Experimental

Synthesis

Reagents. All reagents and solvents were commercial and were used as received. Benzo[1,2-b;4,5-b']dithiophene was purchased from TCI. 2,6-dibromobenzo[1,2-b;4,5-b']dithiophene-4,8-dione^[58], 2,6-dibromobenzo[1,2-b;5,4-b']dithiophene^[59], 4-ethynyl-1-thioacetylbenzene^[60] and 1-tert-butylthio-4-ethynylbenzene^[61] were synthesized according to literature procedures.

NMR and Mass Spectra. ¹HNMR and ¹³CNMR were performed on a Varian Unity Plus (400 MHz) instrument at 25 °C, using tetramethylsilane (TMS) as an internal standard. NMR shifts are reported in ppm, relative to the residual protonated solvent signals of CDCl₃ ($\delta = 7.26$ ppm) or at the carbon absorption in CDCl₃ ($\delta = 77.0$ ppm). Multiplicities are denoted as: singlet (s), doublet (d), triplet (t) and multiplet (m). High Resolution Mass Spectroscopy (HRMS) was performed on a JEOL JMS 600 spectrometer.

UV-Vis and Cyclic Voltammetry UV-Vis measurements were carried out on a Jasco V-630 spectrometer. Cyclic voltammetry (CV) was carried out with a Autolab PGSTAT100 potentiostat in a three-electrode configuration.

General. Unless stated otherwise, all crude compounds were isolated by bringing the reaction to room temperature, extracting with CH₂Cl₂, washing with saturated NaHCO₃, water and then brine. The organic phase was then collected and dried over Na₂SO₄ and the solvents removed by rotary evaporation. Synthetic schemes and NMR spectra are provided in the Supplementary Information.

2,6-dibromobenzo[1,2-b:4,5-b']dithiophene (1). Benzo[1,2-b;4,5-b']dithiophene (540 mg, 2.84 mmol) were dissolved in 70 mL anhydrous THF under an atmosphere of N₂, cooled to -78 °C and n-butyllithium (8.5 mmol, 5.3 mL, 1.6 M in hexane) was added drop-wise. The solution was stirred for 30 min in the cold bath before being warmed to room temperature and stirred for and additional 20 min. The mixture was cooled to -78 °C again and a solution

of CBr₄ (2.8 g, 8.5 mmol) in 5 mL anhydrous THF was added. The solution was stirred for

30 min in the cold bath before being quenched with concentrated sodium bicarbonate solution

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2,6-Bis[(4-acetylthiophenyl)ethynyl]benzo[1,2-b:4,5-b']dithiophene-4,8-dione (BDT-

2). $^{[62]}$ TiCl₄ (0.04 mL, 0.364 mmol) was added drop-wise to a solution of compound 5 (100 mg,

 $0.167\,\mathrm{mmol})$ and $\mathrm{CH_3C(O)Cl}$ ($0.03\,\mathrm{mL},\,0.377\,\mathrm{mmol})$ in $\mathrm{CH_2Cl_2}$ at $0\,^\circ\mathrm{C}$. The resulting mix-

135.17, 131.73, 131.56, 130.31, 121.70, 98.14, 82.55, 46.81, 31.02.

ture was stirred at room temperature for 1 h and the conversion was monitored by TLC (hexanes/CH₂Cl₂, 2:1). Upon completion, the reaction was quenched with water (10 mL). The crude solid was purified by column chromatography to give **BDT-2** (50 mg, 53 %). 1 HNMR (400 MHz, CDCl₃) δ : 7.73 (s, 2H), 7.59 (d, J=8.2, 4H), 7.45 (d, J=8.2, 4H), 2.46 (s, 6H). 13 CNMR (100 MHz, CDCl₃) δ : 195.59, 175.96, 145.20, 136.95, 134.87, 134.20, 133.15, 132.57, 132.50, 125.24, 100.42, 85.49, 33.01. HRMS(ESI) calcd. for C₃₀H₁₇O₄S₄ [M+H]⁺: 569,00042, found: 568.99887.

2,6-Bis[(4-tert-butylthiophenyl)ethynyl]benzo[1,2-b:5,4-b']dithiophene (7). 2,6-dibromobenzo[1,2-b:5,4-b']dithiophene (6; 50 mg, 0.143 mmol) and 1-tert-butylthio-4-ethynylbenzene (4; 68 mg, 0.358 mmol) were dissolved in mixture of fresh distilled Et₃N (5 mL) and anhydrous THF (10 mL). After degassing, the catalysts Pd(PPh₃)₄ (16 mg, 0.014 mmol) and CuI (2.7 mg, 0.014 mmol) were added. The reaction mixture was refluxed overnight under N₂. The crude solid was purified by column chromatography to give 7 (40 mg, 49 %). ¹HNMR (400 MHz, CDCl₃) δ: 8.16 (s, 1H), 8.14 (s, 1H), 7.56 (s, 2H), 7.54 (d, J=4, 4H), 7.51 (d, J=4, 4H), 1.31 (s, 18H). ¹³CNMR (100 MHz, CDCl₃) δ: 141.35, 140.05, 139.89, 136.77, 134.10, 131.29, 125.64, 125.38, 120.85, 117.39, 97.51, 86.99, 49.30, 33.67.

2,6-Bis[(4-acetylthiophenyl)ethynyl]benzo[1,2-b:5,4-b'] dithiophene (BDT-3). [62] TiCl₄ (0.042 mL, 0.388 mmol) was added drop-wise to a solution of compound (7) (100 mg, 0.176 mmol) and CH₃C(O)Cl (0.03 mL, 0.397 mmol) in CH₂Cl₂ at 0 °C. The resulting mixture was stirred at room temperature for 10 min and the conversion was monitored by TLC (hexanes/CH₂Cl₂ 2:1). Upon completion the reaction was quenched with water (10 mL). The crude solid was purified by column chromatography to give BDT-3 (25 mg, 26 %). ¹HNMR (400 MHz, CDCl₃) δ: 8.17 (s, 1H), 8.15 (s, 1H), 7.59 (d, J=7.2, 4H), 7.58 (s, 2H), 7.43 (d, J=8.2, 4H), 2.45 (s, 3H). ¹³CNMR (100 MHz, CDCl₃) δ: 195.88, 141.43, 140.03, 136.90, 134.76, 131.51, 131.48, 126.27, 126.50, 120.94, 117.42, 97.22, 87.15, 32.97. HRMS(ESI) calcd. for C₃₀H₁₈O₂S₄ [M+H]⁺: 539,02624, found: 539.02476.

Self-assembled monolayers

The SAMs of BDT-n were formed via in situ deprotection [41,51] on template-stripped Au substrates. [63] Freshly template-stripped substrates were immersed into 3 mL of 50 μM solutions of the thioacetate precursors in freshly distilled toluene inside a nitrogen-filled glovebox and sealed under a nitrogen atmosphere. The sealed vessels were kept inside a nitrogen flow box [64] (O₂ below 3 %, RH below 15 %) overnight; all subsequent handling and EGaIn measurements were performed inside the flowbox. 1.5 h prior to measurement, 0.05 mL of 17 mM diazabicycloundec-7-ene(DBU) in toluene was added to the precursor/substrate solution. The substrates were then rinsed with toluene and allowed to dry for 30 min before performing the measurements.

Characterization

The SAMs of **BDT-**n were characterized by XPS (laboratory and synchrotron), NEXAFS spectroscopy, UPS and water contact angles. In some cases, SAMs of $CH_3(CH_2)_{15}SH$ or $CH_3(CH_2)_{17}SH$ on Au were used as a reference. See Supplementary Information for details.

Transport measurements

EGaIn. For each SAM, at least 10 junctions were measured on each of three different substrates by applying a bias from $0.00\,\mathrm{V} \to 1.00\,\mathrm{V} \to -1.00\,\mathrm{V} \to 0.00\,\mathrm{V}$ with steps of $0.05\,\mathrm{V}$. At least 20 trace/re-trace cycles were measured for each junction; only junctions that did not short over all 20 cycles were counted as "working junction" for computing yields.

CP-AFM. I-V measurements were performed on a Bruker AFM Multimode MMAFM-2 equipped with a Peak Force TUNA Application Module. The Au on mica substrates were removed from the flowbox immediately prior to measurement, which occurred under ambient conditions by contacting the SAM with a Au-coated Si₃N₄ tip with a nominal radius of 30 nm

(NPG-10, Bruker; resonant frequency: $65\,\mathrm{kHz}$, spring constant: $0.35\,\mathrm{N/m}$). The AFM tip was grounded and the samples were biased from $-1.0\,\mathrm{V} \to 1.0\,\mathrm{V} \to -1.0\,\mathrm{V}$ on $\mathrm{Au^{Mica}}$. 11 trace/re-trace cycles per junction were performed and the top electrode was removed from SAMs between junctions.

Processing. All raw data were processed algorithmically using Scientific Python to generate histograms, Gaussian fits, extract transition voltages and construct differential conductance heatmap plots.

DFT calculations

Calculation were performed using the ORCA 4 software package [65,66] and the ARTAIOS-030417 software package. [67,68] The molecules terminating with thiols were first minimized to find the gas-phase geometry and then attached to two 18-atom Au(111) clusters via the terminal sulfur atoms with a distance of 1.75 Å at hexagonal close-pack hollow sites (hydrogen atom from the thiol was deleted before attaching the electrodes). Single-point energy calculations were performed on this model junction using B3LYP/G and LANL2DZ basis sets according to literature procedures to compute the energy levels. [67] Transmission curves and isoplots of the central molecular orbitals, for isolated molecules without electrodes and terminal hydrogen atoms, were generated using the ARTAIOS-030417 software package and the energy axis was scaled using the $E_{\rm F}$ of $-4.3\,{\rm eV}$. The use of this $E_{\rm F}$ value for comparing transmission trends to the experimental tunneling conduction in Au/SAM//EGaIn junction is supported by the UPS measurements that also give a similar $E_{\rm F}$ value (see Supplementary Information section 1.3.3). It is has also been established experimentally that SAMs of aliphatic and conjugated molecules on Au shift the $E_{\rm F}$ values by 0.85 and 0.98 eV, respectively, (i.e., to $-4.2\,\mathrm{eV}$ to $-4.4\,\mathrm{eV}$) from $-5.2\,\mathrm{eV}$ for a clean gold surface. [69–71] This value of $E_{\rm F}$ was used for all DFT calculations. Further rational for choosing this value of $E_{\rm F}$ and the detailed step-wise procedure for all the calculations involved is further described in Supplementary Information.

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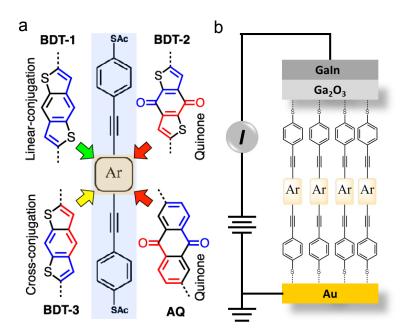


Figure 1: (a) Structures of **BDT-1**, **BDT-2** and **BDT-3** with linearly and cross-conjugated pathways of the cores drawn in blue and red, respectively. The phenylacetylene arms (highlighted in blue) are linearly conjugated. **BDT-1** is linearly-conjugated, **BDT-2** contains a cross-conjugation imposed by the central quinone ring analogous to **AQ** and **BDT-3** is similarly cross-conjugated, but the cross-conjugation separating the two linearly-conjugated pathways arises from the positions of the sulfur atoms relative to the central phenyl ring (there are no exocyclic bonds). (b) Schematic of Au/SAM//EGaIn junction ("/" and "//" denote a covalent and van der Waals interfaces, respectively).

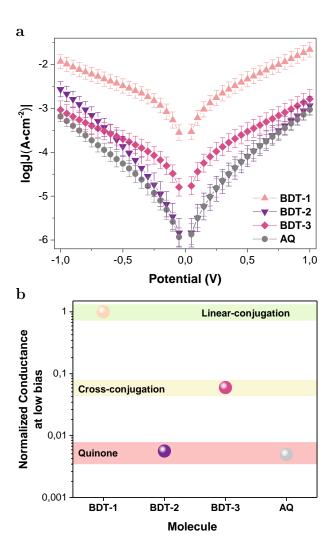


Figure 2: (a) Plots of $\log |J(\text{Acm}^{-2})|$ versus V of Au/SAM//EGaIn junctions comprising SAMs of **BDT-1** (salmon up-triangles), **BDT-2** (purple down-triangles), **BDT-3** (pink diamonds) and \mathbf{AQ} (grey circles). Each datum is the peak position of a Gaussian fit of $\log |J|$ for that voltage. The error bars are 95% confidence intervals taking each junction as a degree of freedom. (b) Normalized low bias conductance, linearly conjugated **BDT-1** (salmon ball) features the highest values, the quinone **BDT-2** (purple ball) and \mathbf{AQ} (grey ball) the lowest and cross conjugated **BDT-3** (pink ball) is in between.

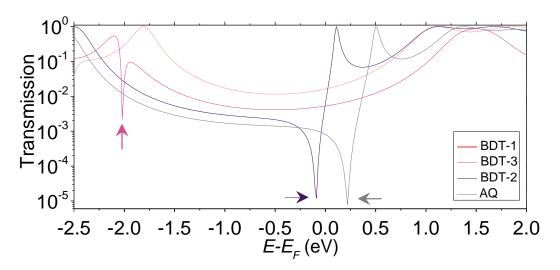


Figure 3: Transmission spectra for isolated molecules of **BDT**-n and **AQ**. The spectrum of **BDT**-1 (salmon) is featureless between the resonances $(T(E) \to 1)$ near the frontier orbitals. The sharp dips in the spectra of **BDT**-2 (purple), **BDT**-3 (pink) and **AQ** (grey) indicated with arrows are destructive QI features. The energies on the bottom axis are with respect to the $E_{\rm F}$ value of $-4.3\,{\rm eV}$.

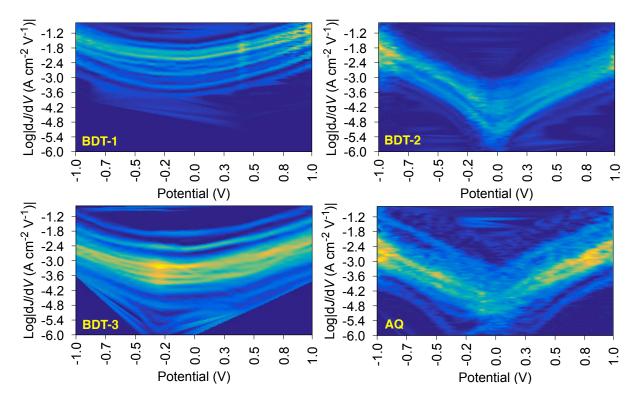


Figure 4: Differential conductance heatmap plots of Au/SAM//EGaIn junctions comprising **BDT-1** (top-left), **BDT-2** (top-right), **BDT-3** (bottom-left) and **AQ** (bottom-left) showing histograms binned to $\log \left| \frac{dJ}{dV} \right|$ (differential conductance, Y-axis) versus potential (V, X-axis). The colors correspond to the frequencies of the histograms and lighter (more yellow) colors indicate higher frequencies. The bright spots near $\pm 1 \, \text{V}$ are due to the doubling of data that occurs in the forward/return J/V traces. The plots for both **BDT-2** and **AQ**, which contain quinones, are V-shaped at low bias and exhibit negative curvature, indicating a destructive QI feature near E_{F} , while the plots of **BDT-1** and **BDT-3** are U-shaped.

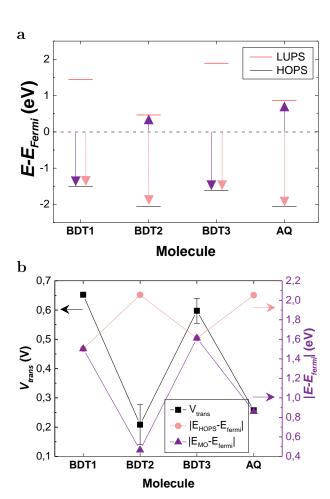


Figure 5: (a) Energy offsets of the frontier orbitals calculated using DFT with respect to $E_{\rm F}$ value of $-4.3\,{\rm eV}$. (b) The energy offsets (salmon and purple lines, right axis) plotted with the measured values of $V_{\rm trans}$ (black line, left axis). The salmon line plots the energy offsets of the HOPS. The purple line plots the smallest energy difference (purple arrows in Figure 5a); $|E_{\rm HOPS}-E_{\rm Fermi}|$ for BDT-1 and BDT-3, $|E_{\rm LUPS}-E_{\rm Fermi}|$ for BDT-2 and AQ. The exact values of $V_{\rm trans}$ and the orbital energies are shown in Table S3 and S5.

Table 1: Summary of the properties of SAMs of \mathbf{BDT} -n and \mathbf{Au}/\mathbf{BDT} - $n//\mathbf{EGaIn}$ junctions.

Compound	BDT-1	BDT-2	BDT-3	C18 reference
XPS thickness (Å)	17 ± 3	18 ± 4	19 ± 4	n.d.
HRXPS thickness (Å)	19.81 ± 0.40	22.30 ± 0.45	17.17 ± 0.34	20.9
Averaged XPS thickness (Å)	18.4	20.2	18.4	n.d.
Water contact angle (°)	68.3 ± 4.8	65.8 ± 4.0	62.8 ± 4.6	104.2 ± 2.2
Density (10^{14} molecules per cm ²)	2.05	3.30	2.33	4.63
Area molecules per $Å^2$)	48.8 ± 2	30.3 ± 2	43.0 ± 2	21.6
$\log J \ @0.5\mathrm{V}\ (\mathrm{Acm}^{-2})$	-2.34 ± 0.17	-4.09 ± 0.23	-3.53 ± 0.20	$-4.96 \pm 0.87^{[41]}$
Yield of working junctions $(\%)$	88.9	93.8	84.2	$79^{[41]}$
Num. working EGaIn junctions	32	30	32	$28^{[41]}$
Total J/V traces	643	626	666	$280^{[41]}$

Graphical TOC Entry

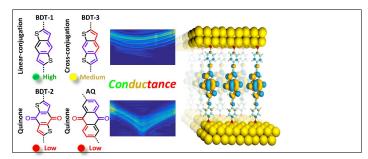


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Three different benzodithiophene derivatives were designed to isolate the effects of bond topology from that of functional groups in quantum interference to examine the role of the quinone functionality separate from cross-conjugation in SAMs based large area molecular junctions. Quinones are found to suppress tunneling transport further than cross-conjugation alone and to switch the mechanism from tunneling mediated by occupied states to tunneling mediated by unoccupied states.