# **On Thermodynamic Characteristics of In–I System Compounds**

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**Abstract.** The temperature dependences of the total vapour pressure for solid and liquid  $InI_2$  and liquid  $InI_3$  were measured by the static method with a membran-gauge manometer. The heat capacities above phases and solid  $InI_3$  were also obtained. The partial pressures  $In_2I_6$ ,  $InI_3$ , InI,  $In_2I_4$ ,  $I_2$ , I were calculated. Absolute entropies and enthalpies of formation  $InI_2$  (crII),  $InI_2$  (l),  $In_2I_4$  (g),  $InI_3$  (l),  $InI_3$  (g),  $In_2I_6$  (g) were obtained. We used the least square method to ob-

tain the mutual consistent sets of data on thermodynamic characteristics of indium iodides in condensed and gaseous phases. The result of this work is the set of standard formation enthalpies and absolute entropies of the In–I system compounds.

Keywords: Indium iodides; Thermodynamic data

# Thermodynamische Daten von Indiumiodid

**Inhaltsübersicht.** Die Messungen des Gesamtdruckes über InI<sub>3</sub> (fl), InI<sub>2</sub> (f) und InI<sub>2</sub> (fl) werden mit der statischen Methode mit einem Quarz-Membran-Manometer ausgeführt. Die Temperaturabhängigkeiten der spezifischen Wärmen werden für festes und flüssiges InI<sub>2</sub> und InI<sub>3</sub> bestimmt. Die Partialdrücke für In<sub>2</sub>I<sub>6</sub>, InI<sub>3</sub>, InI, In<sub>2</sub>I<sub>4</sub>, I<sub>2</sub>, I werden berechnet. Daraus werden die Bildungenthalpien und Standardentropien von  $InI_2$  (crII),  $InI_2$  (fl),  $In_2I_4$  (g),  $InI_3$  (fl),  $InI_3$  (g),  $In_2I_6$  (g) abgeleitet. Abschließend wird ein kompletter Satz von thermodynamischen Daten für das System Indium–Iod vorgeschlagen.

#### Introduction

The measurement of the vapour total pressure as a function of temperature and vapour density in a wide concentration range permits us to design various thermodynamic cycles. Such an approach allows one to determine the thermodynamic characteristics of examined compounds in the investigated system, to avoid possible systematic errors and to get the mutual consistent thermodynamic data table. Such an investigation and its advantages have been demonstrated by the author of [1] for the Ge–I system. We have tried to apply a similar approach for study of the properties of the In–I system compounds.

According to the literature data three individual compounds InI,  $InI_2$  and  $InI_3$  in the In–I system are known to exist in crystalline state. Polymorphic transitions for indium diiodide and indium triiodide have been determined reliably. However, different information on temperatures of these transitions has been presented in literature [2–6]: from 426 to 435 K for InI<sub>2</sub> and from 200 to 342 K for InI<sub>3</sub>. According to [7–8],  $In_2I_6$ ,  $InI_3$ ,  $InI_2$ , InI, and  $I_2$  are the main species in the gas phase. But, electronographic vapour investiga-

Institute of Inorganic Chemistry Russian Academy of Sciences, Sibirian Branch 3 Lavrentiev av. Novosibirsk 90/Russia tions performed in the In–I system [9] has some doubt whether the amount of  $InI_2$  being considerable or not. Besides, the presence of  $In_2I_4$  dimers is revealed. The latter conclusion is indirectly confirmed by the data of [6] on the structure of solid indium diiodide.

#### **Experimental**

Indium iodides were synthesized from the elements in sealed quartz ampoules under vacuum. Indium, grade HP 50-4 (High Pure 50-4) (impurities sum of 50 controlled elements is not more than  $10^{-4}$  at.%), and iodine, grade HP 16-4 (impurities sum of 16 controlled elements is not more than  $10^{-4}$ at.%) were used. Iodine was resublimated three times in vacuum for additional purification from other halogens and organic impurities; its vapour being overheated up to 1300 K. Monoiodide was synthesized in the excess of indium and triiodide - in excess of iodine. The synthesis was performed in a three section ampoule, initial reagents were loaded into the side sections. Iodine temperature was held at 400 K. The indium temperature was chosen in such a way, that the reaction product pressure did not exceed the iodine pressure. The synthesis being over the product was purified in the same ampoule by its distillation into the central section. The impurities being more volatile and the iodine excess were collected in the section previously occupied by iodine. For this purpose the section was cooled by liquid nitrogen. Indium diiodide was synthesized via melting stoichiometric amounts of InI and InI<sub>3</sub>. The purification was performed by freezing out light volatile impurities in liquid nitrogen. The InI<sub>3</sub> synthesis from elements yields the high temperature

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	Sample No	Obtained, mass%		Calculated, mass.%	
		In	Ι	In	Ι
InI	N 3	$47.57 \pm 0.11$	$52.49 \pm 0.17$	47.50	52.50
InI <sub>2</sub>	N 5	$31.06\pm0.15$	$68.79 \pm 0.16$	31.15	68.85
InI <sub>3</sub>	N 1	$23.14\pm0.18$	$76.67 \pm 0.16$	23.17	76.83

**Table 1**Chemical analyses of samples

phase – "yellow" triiodide one (InI<sub>3</sub>(cr. I)). Its transition to the low temperature "red" phase (InI<sub>3</sub> (cr. II)) is kinetically restricted. The "red" phase was prepared via the "yellow" phase recrystallization from its solution in purified n-butyliodide according to the procedure described elsewhere [10]. All the compounds obtained were characterized by chemical, spectral and X-ray powder diffraction analysis. The indium content was controlled by the weight method (In<sub>2</sub>O<sub>3</sub>). The iodine content was measured by potentiometric titration using AgNO<sub>3</sub>. The results of three typical samples (three parallel runs were carried for every sample) are presented in Table 1, errors being given for 95% reliable intervales. Spectral analysis showed the sum of all impurities controlled to be not more than 0.015 at.%. According to the X-ray power diffraction results all the samples were found to be single phases. Unfortunately, in the reference books there are no X-ray power diffraction data for InI2 and InI3. Hence, we have recorded their Raman scattering spectra which agree well with previously obtained data [6, 11]. All the operations with the samples were carried out in such a way that contact with moisture and atmosphere gases was excluded.

#### **Experimental Device and Measurements**

To determine the heat capacities of condensed phases we have used the drop calorimetry allowing us to measure  $H^0(T)-H^0(298)$  (P<sup>0</sup> = 101,325 kPa) temperature dependence [12]. The experiments were performed on samples placed into 25–30 mm long quartz ampoules with inner diameter of 7–8 mm and 1 mm wall thickness. Thus, efficiency heat is 30–

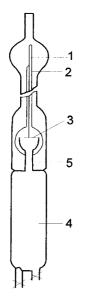


Fig. 1 Quartz membrane zero-manometer.

1 – the moving stock; 2 – the basic stock; 3 – the sensor element; 4 – the working cell; 5 – the compensative part.

50% of the total heat evolving in the calorimeter. The efficiency heat can be determined as the difference of experiments performed for ampoules with samples and without. The ampoule is installed in the gradient-less zone of a furnace (0.2 K per 50 cm), the accuracy of thermodynamic temperature measurements being 0.3-0.5 K. The duration of the ampoule exposure in the furnace depends on how fast the heat equilibrium was attained. In our case it was 20-25 min. The heat losses upon the sample transfer from the furnace to the calorimeter did not exceed 0.5%. The calorimetric cell (massive copper block thermostated at 298.15 K) was screened from the furnace by special movable shutters. Temperature fluctuations in the thermostat did not exceed 0.002 K (the thermometer circuit sensitivity was  $4 \cdot 10^{-5}$  K). The average duration of main period was 15-20 min. The Pfaundler correction for the heat exchange did not exceed 5% of the overall temperature raise in the cell. Since the spread of the calorimeter cooling constants was about 7% and the error of the heat equivalent determination was 0.15%, the overall error of heat measurement was less than 0.7%. The calorimeter calibration with respect to standards ( $\alpha$ -Al<sub>2</sub>O<sub>3</sub> and mercury) gave the random spread of 0.5%.

When the measurements being completed, the ampoule was transferred into the furnace heated to a new temperature with a special device. This equipment essentially improved the efficiency of a calorimeter.

The pressure of saturated and unsaturated vapours was measured by the static method using quartz membrane zero manometers of the Novikov-Suvorov construction [13]. These manometers have the same or even higher sensitivity as those used by *Oppermann* [14], their sensor area being much smaller. This factor decreases elastic and inelastic deformation always arising in a glass at the high temperature change. Fig. 1 presents the draft of the device used.

First, this device with both sides opened to atmosphere, was installed in the furnace and the calibration experiment was carried out to determine the corrections for the temperature drift of zero position of the moving stock. Such corrections did not exceed  $1.5 \pm 0.1$  Torr (1 Torr = 133,32 Pa). The loading of the device with the substance under study was performed as described elsewhere [14].

Before loading the device was heated for 4 hours under high vacuum at 1000–1100  $^{\circ}$ C to remove moisture and other substances. During the measurements the zero manometer was thermostated. The accuracy of the temperature maintenance was 0.2 K.

The maximum temperature drop in the device did not exceed 2 K at 1200 K. The error of the temperature measurement estimated from experiments with nitrogen and mercury was  $\pm 1$  K at 600 K and increased to 2 K at 1200 K. The pressure measurement error including compensation errors and the zero position determination errors was less than 0.3 Torr.

#### **Experimental Results**

In this work we used following expression for temperature function of heat capacity:

$$Cp = A \cdot T^{-2} + B \cdot T^{-1} + C + D \cdot T + E \cdot T^2.$$

For the substance under consideration some terms were statistical negligible.

The enthalpy temperature dependencies for the  $InI_2$  low temperature modification ( $InI_2$ , crII) were measured in 24 calorimetric experiments at the temperature range of 346–421 K.

Results were approximated by the following expression:

$$\begin{split} H^0(T) - H^0 & (298.15) = -18100 + 37.35 \cdot T \\ &+ 0.1016 \cdot T^2 - 0.7803 \cdot 10^{-4} \ T^{-3} \ J/mol. \end{split}$$

Thus,  $C_p^0(T) = 37.35 + 0.2032 \cdot T - 0.2341 \cdot 10^{-3} \cdot T^2 \text{ J/mol/K}$ 

17 measurements of the high temperature  $InI_2,cr$  enthalpy at 434–488 K give:

 $H^0(T)$  –  $H^0$  (298.15) = 24520 – 134.6  $\cdot$  T + 0.2453  $\cdot$  T  $^2$  J/mol and

 $C_p^0(T) = 0.4906 \cdot T - 134.6 \text{ J/mol/K}$ 

For the liquid  $InI_2$  15 measurements at 499–684 K give:

 $H^{0}(T) - H^{0}(298.15) = 96.11 \cdot T + 0.01577 \cdot T^{2} - 21820 \text{ J/mol};$ 

 $C_p^0(T) = 96.11 + 0.03154 \cdot T J/mol/K;$ 

 $\Delta_{tr}^{T} H^{0} (InI_{2}, 428 \text{ K}) = 1.5 \pm 0.2 \text{ kJ/mol}$ 

 $\Delta_{\rm m} {\rm H}^0 ({\rm InI}_2, 497 {\rm K}) = 11.6 \pm 0.3 {\rm kJ/mol}$ 

It should be noted that we have met some kinetic restrictions when studying thermodynamic characteristics of the diiodide transition  $crII \rightarrow crI$ .

We also studied the InI<sub>3</sub>, crI-"yellow" phase (28 measurements at 335–480 K) and liquid triiodide (34 measurements, 483–723 K). For InI<sub>3</sub>,crI we obtained:  $H^{0}(T) - H^{0}(298.15) = -32720 + 123.4 \cdot T - 0.0701 \cdot T^{2} + 0.8206 \cdot 10^{-4} T^{3} \cdot J/mol;$  $C_{0}^{n}(T) = 123.4 - 0.1402 \cdot T + 2.462 \cdot 10^{-4} T^{2} \cdot J/mol \cdot K$ 

For InI<sub>3</sub>, liq we obtained:

$$\begin{split} H^0(T) - H^0 & (298.15) = -224900 + 0.05884 \cdot T^2 + 4.035 \cdot \ln T \text{ J/mol}; \\ C^0_n(T) = 0.1177 \cdot T + 40350/T \text{ J/mol} \end{split}$$

The equilibrium of gas phase with liquid  $InI_3$  was studied by the static method. Four series of experiments were performed, varying the ratio sample weight to volume from 0.2 to 40 g/l, the temperature from 477 K to 674 K. Data analysis on saturated vapor shows that under the examined conditions the pressure does not depend on initial concentrations. Thus, at  $InI_3$  evaporation the equilibrium is monovariant. The data of all four series (81 experimental points) can be approximated by the following expression:

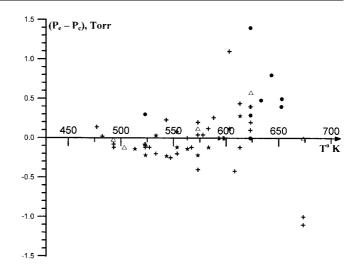
 $\ln(P/P^{0}) = 13.723 - 10130/T \pm 2\sigma,$ where  $\sigma^{2} = 0.000700 - 0.872/T + 272/T^{2}$  (477–674 K)

For this and following expressions  $P^0 = 1$  atm = 760 Torr = 101.325 kPa.

The approximation was connected with minimization of the following function [23]:

 $Y = \Sigma (P_e - P_c)^2 / W$ 

where  $P_e$  – experimental value of the pressure,  $P_c$  –



**Fig. 2** The difference between the experimental and calculated from the regression expression values of total pressure for the saturated vapour of the liquid InI<sub>3</sub> as the temperature function at the different InI<sub>3</sub> concentration, C mol  $\cdot$  cm<sup>-3</sup> + -0.80  $\cdot$  10<sup>-4</sup>;  $\Delta$  - 0.13  $\cdot$  10<sup>-4</sup>;  $\bullet$  - 0.40  $\cdot$  10<sup>-5</sup>;  $\star$  - 0.64  $\cdot$  10<sup>-5</sup>.

calculated value, W – statistical weight of the experimental point. The maximal difference between the total pressure values was -1.4 Torr at P(tot) = 200 Torr and T = 673 K (Fig. 2).

As for unsaturated vapor the average gas molecular mass greatly decreases if the temperature is increased. For example, for the initial concentration of  $0.64 \cdot 10^{-3}$  mol/l the molecular mass was changed from 600 (680 K) to 433 (1173 K) (M<sub>InI3</sub> = 495.53). So, the polymerization and dissociation processes occur in the equilibrium gas phase.

At studing the evaporation processes of  $InI_2$  it was found that the saturated vapor pressure was not constant at constant temperature but depended on the initial  $InI_2$  concentration. Increasing the ratio of the sample weight to volume, we could find the conditions at which the solution composition changed insignificantly. 4 measurement series (75 points, concentra-

**Table 2** Pressure concentration dependence over  $InI-InI_3$ melts as a function of temperature at 520–830 K. Approximation accuracy ( $\Delta P/P$ ) was 5% for 520 K and 0.5% for 800 K.

Composition	Number of	$\ln P/P^0 = A - B/T^a)$		
InI <sub>x</sub>	points	А	В	
2.103	12	11.26	9526	
2.088	13	11.53	9841	
2.082	15	11.55	9857	
2.077	28	11.74	10027	
2.073	17	11.72	10037	
2.061	10	11.91	10244	
2.038	10	12.09	10496	
2.026	7	12.36	10870	
1.816	8	12.37	11690	
1.683	4	11.53	11270	

<sup>a</sup>)  $P^0 = 1$  atm = 101,325 kPa

tions ranged from 50 to 100 g/l) can be approximated by the following expression:

 $\ln(P/P^{0}) = 11.562 - 10524/T \pm 2\sigma,$ where  $\sigma^{2} = 0.00266 - 4.101/T + 1590/T^{2}$  (523–843 K)

In special experimental series we also studied how the pressure over  $InI-InI_3$  solutions depends on the concentration.

Table 2 presents the results obtained.

Interpolation the data of Table 2 to x = 2 gives equation:  $ln(P/P^0) = 12.34 - 10980/T$ 

This equation agrees well with the measurements over liquid  $InI_2$ .

#### Discussion

Earlier the same samples were studied by electronographic method [9] and low temperature calorimetry [15–17]. The structure and the atom–atom distances for the gas indium iodides were determined in [9]. Besides set of frequencies of vibration was recommended in this publication. The low temperature calorimeter measurements were performed in range from 7 to 300 K [15–17]. The data [9, 15–17] obtained gave information on the heat capacities and absolute entropies of the gas phase (model: harmonic oscillator – rigid rotator) and condensed indium iodides:

$$\begin{split} C^0_p(\mathrm{InI},g) &= 37.40 - 0.418 \cdot 10^5 \ \mathrm{T}^{-2}, \\ S^0_{298}(\mathrm{InI},g) &= 267.2 \pm 1.0 \ \mathrm{J/mol/K} \\ C^0_p(\mathrm{InI}_{3,g}) &= 83.10 + 2.978 \cdot 10^{-5} \ \mathrm{T} - 17710 \cdot \mathrm{T}^{-2}, \\ S^0_{298}(\mathrm{InI}_{3,g}) &= 401 \pm 4 \ \mathrm{J/mol/K} \\ C^0_p(\mathrm{In}_2\mathrm{I}_{6,g}) &= 182.9 - 3.594 \cdot 10^{-5} \ \mathrm{T} - 306500 \cdot \mathrm{T}^{-2}, \\ S^0_{298}(\mathrm{In}_2\mathrm{I}_{6,g}) &= 655.6 \pm 8 \ \mathrm{J/mol/K} \\ C^0_p(\mathrm{In}_2\mathrm{I}_{4,g}) &= 131.3 + 1.216 \cdot 10^{-4} \ \mathrm{T} - 383600 \ \mathrm{T}^{-2}, \\ S^0_{298}(\mathrm{In}_2\mathrm{I}_{4,g}) &= 512.5 \pm 10 \ \mathrm{J/mol/K} \\ C^0_p(\mathrm{In},\mathrm{cr}) &= 58.10 - 0.03708 \cdot \mathrm{T} + 70.59 \ 10^{-6} \ \mathrm{T}^2 \ (298\text{-}637 \ \mathrm{K}) \\ S^0_{298}(\mathrm{InI},\mathrm{cr}) &= 119.9 \pm 0.2 \ \mathrm{J/mol/K} \ [14] \\ S^0_{298}(\mathrm{InI}_2,\mathrm{crII}) &= 173.6 \pm 0.4 \ \mathrm{J/mol/K} \ [16] \end{split}$$

 $S_{298}^{0}(InI_{2},crII) = 1/3.6 \pm 0.4 J/mol/K [16]$  $S_{298}^{0}/InI_{3},crI) = 223.0 \pm 0.4 J/mol/K [17]$ 

The study of the unsaturated vapor pressures within wide temperature and concentration ranges in the In–I system [18] (these data are consistent with results [9]) allowed us to obtain thermodynamic characteristics of processes occurring in the gas phase:

 $In_2I_6,g = 2 InI_3,g \tag{1}$ 

 $InI_{3,g} = InI_{,g} + I_{2,g}$ (2)

 $In_2I_4,g = InI,g + InI_3,g$ 

$$\begin{split} \Delta_{r1} H^0 & (298.15) = 102.1 \pm 0.9 \text{ kJ} \\ \Delta_{r1} S^0 & (298.15) = 147.2 \pm 1.2 \text{ J/K} \end{split}$$

$$\begin{split} \Delta_{r2} H^0 & (298.15) = 193.6 \pm 3.8 \text{ kJ} \\ \Delta_{r2} S^0 & (298.15) = 127.9 \pm 3.3 \text{ J/K} \end{split}$$

$$\begin{split} \Delta_{r3} H^0 & (298.15) = 137.4 \pm 1.8 \text{ kJ} \\ \Delta_{r3} S^0 & (298.15) = 154.6 \pm 1.9 \text{ J/K} \end{split}$$

Data processing [18] was performed using the above heat capacities and including the dissociation reaction of molecular iodine [19]. Taking into account the reaction characteristics (1) and (2) we can write the following expression for the partial pressures of  $InI_3$  and  $In_2I_6$  in equilibrium with liquid triiodide indium:

InI<sub>3</sub>,liq = InI<sub>3</sub>,g

(5)

(6)

(7)

(8)

 $\ln(P/P^{0}) = -59.415 - 33373/T - 10650/T^{2} + 9.9946 \cdot \ln T - 0.007076 \cdot T + 4852.9 \cdot \ln T/T, (477-674 \text{ K})$ 

Extrapolation gives:

$$\begin{split} \Delta_{r4} H^0 & (298) = 108.07 \pm 0.13 \text{ kJ} \\ \Delta_{r4} S^0 & (298.15) = 163.8 \pm 0.9 \text{ J/K} \end{split}$$

For the process:

 $2 \operatorname{In}I_3, \operatorname{liq} = \operatorname{In}_2 I_6, g$ 

 $\ln(P/P^{0}) = -149.93 - 53890/T - 18431/T^{2} + 21.993$  $\cdot \ln T - 0.014153 \cdot T + 9705.8 \cdot \ln T/T, (477-674 \text{ K})$ 

$$\begin{split} &\Delta_{r5}H^0(298) = 114.05 \pm 0.95 \text{ kJ} \\ &\Delta_{r5}S^0 \ (298.15) = 180.4 \pm 1.5 \text{ J/K} \end{split}$$

In a similar way we obtained for the congruent transition process of liquid diiodide into vapor:

## $2 \operatorname{In} I_2, \operatorname{liq} = \operatorname{In}_2 I_4, g$

 $\Delta_{r6} H^0 (497) = 104.5 \pm 7.5 \text{ kJ}$  $\Delta_{r6} S^0 (497) = 92.8 \pm 2.8 \text{ J/K}$ 

Taking into consideration the temperature dependencies of the phase heat capacities and the enthalpies of the phase transitions for the diiodide sublimation process we can write:

## $2InI_2,crII = In_2I_4,g$

$$\begin{split} \Delta_{r7} H^0 & (298.15) = 137.9 \pm 8.0 \text{ kJ} \\ \Delta_{r7} S^0 & (298.15) = 164.0 \pm 3.0 \text{ J/K} \end{split}$$

Earlier [24] the thermodynamic characteristics of the vapourisation process of indium monoiodide were investigated. It is monovariant equilibrium (M(experimental) =  $239.5 \pm 2.7$ ; M(theory) = 241.72):

According [24]  $C_p^0(InI, liq) = 69 \pm 2 \text{ J/mol/K} (640-940 \text{ K})$ 

For reaction (8):

(3)

 $\Delta_{\rm r8} {\rm H}^0$  (637.5 K) = 106.2 ± 0.2 kJ/mol  $\Delta_{\rm r8} {\rm S}^0$  (637.5 K) = 107.05 ± 0.25 J/mol/K

The above experimental data allow us to determine thermodynamic characteristics for condensed phases and gas species by means of several thermodynamic cycles. For example, we can calculate the absolute entropy of gaseous InI<sub>3</sub> as:

$$\begin{split} S^0(InI_{3,g}) &= S^0(InI_{3,c}rII) + \Delta_{tr}S^0(InI_{3,c}rII) \\ &+ \Delta_m S^0(InI_{3,c}rI) + \Delta_{r4}S^0 \\ &= S^0(InI_{3,c}rI) + \Delta_m S^0(InI_{3,c}rI) + \Delta_{r4}S^0 \\ &= (S^0(In_2I_{6,g}) - \Delta_{r1}S^0)/2 \\ &= S^0(I_{2,g}) + S^0(InI,g) - \Delta_{r2}S^0 \\ &= S^0(In_2I_{4,g}) - S^0(InI,g) + \Delta_{r3}S^0 \end{split}$$

Moreover, we can compare this value with that calculated from molecular constants according to the model "harmonic oscillator – rigid rotator" and can determine the absolute entropies of gaseous InI,  $In_2I_4$ ,  $In_2I_6$ by means of various thermodynamic cycles. We can design the similar cycles to determine the formation enthalpies of indium iodides using the calorimetry data on the direct synthesis of these compounds [20, 21]:

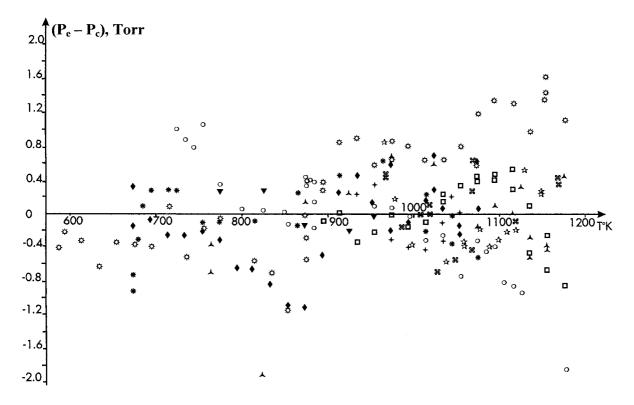
$In(cr) + 1.5 I_2(cr) = InI_3(crI)$	$\Delta H_{298}^0 = -223.9 \pm 0.6 \text{ kJ}$
$InI(cr) + I_2(cr) = InI_3(crI)$	$\Delta H_{298}^0 = -121.9 \pm 0.3 \text{ kJ}$
$InI_2(crII) + 0.5 I_2(cr) = InI_3(crI)$	$\Delta H_{298}^0 = -47.7 \pm 1.3 \text{ kJ}$
$InI_3(crII) = InI_3(crI)$	$\Delta H^0_{298} = 5.0 \pm 0.2 \text{ kJ}$
$In(cr) + 0.5 I_2(cr) = InI(cr)$	$\Delta H^0_{298} = -102.9 \pm 1.3 \text{ kJ}$
$In(cr) + 1.5 I_2(cr) = InI_3(crI)$	$\Delta H^0_{298} = -224.3 \pm 0.8 \text{ kJ}$
$In(cr) + 1.5 I_2(cr) = InI_3(crII)$	$\Delta H^0_{298} = -219.5 \pm 1.4 \text{ kJ}$

 Table 3
 Standard formation enthalpies and absolute entropies of indium iodides

Composition	Phase	$\Delta_{\rm f} {\rm H}^0$ (298.15) kJ/mol	S <sup>0</sup> (298.15) J/mol/K
InI	cr	$-102.5 \pm 0.2$	$120.1 \pm 0.1$
InI	g	$26.4 \pm 0.5$	$268.1 \pm 0.5$
InI <sub>2</sub>	crII	$-176.0 \pm 0.4$	$173.9 \pm 0.4$
$In_2I_4$	g	$-215 \pm 9.0$	$511.8 \pm 8.4$
InI <sub>3</sub>	crII	$-229.2 \pm 0.6$	$205.4 \pm 0.5$
InI <sub>3</sub>	crI	$-224.0 \pm 0.2$	$223.0 \pm 0.4$
InI <sub>3</sub>	g	$-105.4 \pm 1.5$	$402.0 \pm 2.0$
$In_2I_6$	g	$-311.9 \pm 7.9$	$650.1 \pm 8.5$

**Table 4**Characteristics of phase transitions of the In–I system compounds

Compositi	on Type of transition	Т, К	P(atm)	$\Delta H^0$ , kJ
InI	$cr \rightarrow l$	637.5	$9.08 \times 10^{-4}$	$17.26 \pm 0.25$
InI	$l \rightarrow InI(g)$	637.5	$9.08 \times 10^{-4}$	105.12
InI <sub>2</sub>	crII → crI	428.1	$1.59 \times 10^{-8}$	$1.29 \pm 0.21$
InI <sub>2</sub>	$crI \rightarrow In_2I_4$	428.1	$0.59 \times 10^{-8}$	130.62
$InI_2$	$crI \rightarrow InI_3(g) + InI(g)$	428.1	$1.00 \times 10^{-8}$	265.02
$InI_2$	$crI \rightarrow l$	497.2	$2.60 \times 10^{-6}$	$11.85 \pm 0.35$
$InI_2$	$21 \rightarrow In_2I_4(g)$	497.2	$0.90 \cdot 10^{-6}$	103.64
$InI_2$	$21 \rightarrow InI_3(g) + InI(g)$	) 497.2	$1.70 \times 10^{-6}$	237.19
InI <sub>3</sub>	$crII \rightarrow crI$	291.5	$1.27 \times 10^{-12}$	$5.16 \pm 0.05$
InI <sub>3</sub>	$crI \rightarrow InI_3(g)$	291.5	$1.26 \pm 10^{-12}$	118.71
InI <sub>3</sub>	$2crI \rightarrow In_2I_6(g)$	291.5	$0.01 \times 10^{-12}$	136.31
InI <sub>3</sub>	$crI \rightarrow l$	480.2	$2.60 \times 10^{-4}$	$18.48 \pm 0.24$
InI <sub>3</sub>	$l \rightarrow InI_3(g)$	480.2	$2.10 \times 10^{-4}$	95.55
InI <sub>3</sub>	$2l \rightarrow In_2I_6(g)$	480.2	$0.50 \times 10^{-4}$	93.13



**Fig. 3** The difference between the experimental and calculated according to the mutual consistent sets of the thermodynamic data values of total pressure for the unsaturated vapour for the In–I system as the temperature function at the different concentration (C mol · cm<sup>-3</sup>) of In and I, correspondently. O  $- 0.5917 \cdot 10^{-5}$ ,  $0.9187 \cdot 10^{-5}$ ;  $\bigcirc - 0.5956 \cdot 10^{-5}$ ,  $0.9267 \cdot 10^{-5}$ ;  $\bigstar - 0.4378 \cdot 10^{-5}$ ,  $0.5846 \cdot 10^{-4}$ ;  $\Re - 0.3643 \cdot 10^{-5}$ ;  $0.7031 \cdot 10^{-5}$ ;  $\blacktriangledown - 0.4088 \cdot 10^{-5}$ ,  $1.2264 \cdot 10^{-5}$ ;  $\blacktriangle - 0.6403 \cdot 10^{-5}$ ,  $1.9209 \cdot 10^{-5}$ ;  $\blacklozenge - 0.1267 \cdot 10^{-4}$ ,  $3.8010 \cdot 10^{-4}$ ;  $\Re - 0.6245 \cdot 10^{-5}$ ,  $2.0625 \cdot 10^{-5}$ ;  $\square - 0.1412 \cdot 10^{-5}$ ,  $0.8358 \cdot 10^{-5}$ ;  $\Re - 0.6083 \cdot 10^{-5}$ .

It should be noted that all the calorimetric data of papers [20, 21] were obtained on samples identical to those used in this study. It allow us to avoid possible systematic errors related to procedures of the sample preparation and pretreating. The calculation of the same characteristics (e.g. absolute entropy and enthalpy of formation for InI<sub>3</sub>) via various thermodynamic cycles showed that, within the accuracy of measurements, the values obtained were in a good agreement. So, we can use the least square method to obtaining the recommended characteristics of indium iodides. (Under recommended characteristics we mean the formation enthalpies of gas species and condensed phases, that are most stable at P = 1 atm and T = 298.15 K). Moreover, the data on polymorph transitions and melting of substances are included here as well. Using these values we can calculate all the other thermodynamic characteristics of physico-chemical processes.

Thermodynamic values measured experimentally cannot always be considered as independent. For example, the experimental values of the evaporation entropy and enthalpy are correlated to each other. So, to obtain the consistent set of recommended characteristics we used the estimation procedure of parameters of models with several dependent variables [22]. According to this procedure one should minimize the following function:

$$\psi$$
 y<sub>Et</sub> y<sub>Ct</sub>  $\sigma$ <sup>ts</sup> y<sub>Es</sub> y<sub>Cs</sub>  $t$  1 s 1

are  $y_E$  experimental values of thermodynamic characteristics of physico-chemical processes,  $y_c$  – the values calculated over a set of recommended values, **t** or **s** is the item No of the value observed in the experiment (t, s = 1, ..., ),  $\sigma^{ts}$  – the element of matrix reverse to covariant matrix:

$$\begin{array}{ccc} \mathbf{A} & \begin{matrix} \sigma_{11} & \sigma_1 \\ \\ \sigma_{1} & \sigma \end{matrix}$$

Diagonal coefficients of the A-matrix are determined by the estimates of experimental value dispersions, nondiagonal coefficients equal zero for independent observations and can be calculated from correlation coefficients in other cases [22]. Table 3 and 4 present the mutual consistent sets of data on thermodynamic characteristics of indium iodides. The maximal difference between experimental and calculated values of total pressure of unsaturated vapour is 1.8 Torr at P(tot) = 170 Torr, vapour density = 0,0036 mol/l for mixture InI<sub>3</sub> + 2.5 I<sub>2</sub> and T = 1180 K (Fig. 3).

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