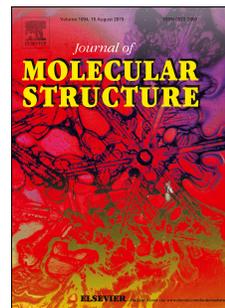


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Synthesis, characterization and theoretical study in gaseous and solid phases of the imine 4-Acetyl-N-(4-methoxybenzylidene)aniline

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Abstract

In the present paper we describe the synthesis and characterization of the Schiff's base or imine 4-Acetyl-N-(4-methoxybenzylidene)aniline (**1**), which provided experimental support for the theoretical calculations. The imine was characterized by infrared spectroscopy and single crystal XRD techniques. The computational studies were performed using the density functional theory (DFT) for the gaseous and solid phases. As similar compounds already shown biological activity, the pharmacokinetic properties of (**1**) were evaluated. Our results shown that (**1**), in its gaseous form, it is electronically stable and has pharmacological drug like properties. Due to its structural similarity with commercial drugs, it is a promise candidate to act as a nonsteroidal anti-inflammatory and to treat dementia, sleep disorders, alcohol dependence, and psychosis. From the solid state calculations we obtain that (**1**) is a low gap semiconductor and can act as an absorber for electromagnetic radiations with energy greater than ~ 0.9 eV.

Keywords: Schiff's base, ab initio calculations, electronic structure, ADME properties

1. Introduction

1 Nitrogenated compounds obtained by condensation reaction of primary amines and
2 aldehydes are called Schiff's base or imines and have a characteristic C=N double bond,
3 where the nitrogen atom is bonded to an aryl or alkyl group. Imines play an important
4 role in Coordination Chemistry due to the formation of stable complexes with transition
5 metals [1]. In our recent papers, several nitrogenated ligands, mainly imines, are employed
6 coordinating to the metal center, such as palladium(II), in order to obtain new complex
7 species with antitumour [2], tuberculostatic [3; 4] and leishmanicidal activities [5].

8 Computational methods are widely used in material science [6] and in drug devel-
9 opment [7] to simulate and predict physical and chemical relevant properties. In material
10 science, properties like electron density, electrostatic potential, electronic orbitals and lower
11 energy conformer (for isolated molecules) and electronic bands and crystal energy (for pe-
12 riodical systems) are among the successfully predicted properties [6]. In drug development
13 process, the quantitative structure activity relationship and the absorption, distribution,
14 metabolism and excretion (ADME) properties are some predicted properties [7].

15 Due to its wide variety of applications [8; 9], Schiff's base form an important class of
16 organic compounds. In the literature, several computational studies of imines that could aid
17 the design and development of new drugs, the understanding of the coordination geometry
18 and electronic properties of isolated and complexes Schiff's base have been performed [10-
19 14].

20 The aim of this work is to present an experimental and theoretical study of the title
21 compound (see Fig. 1 for a 2D diagram). The paper is organized as follows. The ex-
22 perimental details are presented in section 2. The theoretical methodology and tools are
23 described in section 3. In section 4, the discussion of the results for the gaseous and solid
24 phase of (**1**) are presented. Our conclusions are summarized in section 5.

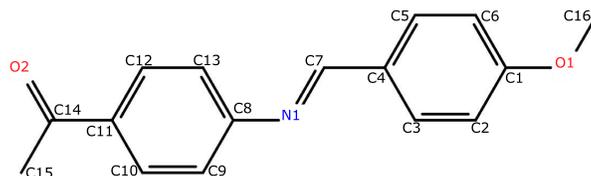


Figure 1: 2D structure diagram. Chemical Abstracts Service (CAS) registry number: 23596-02-3.

2. Experimental methodology

2.1. Synthesis

The materials commercially available were used without purification. The synthesis of the title imine (**1**) was carried out at room temperature. This compound was prepared as described: to a solution containing p-acetophenoneimine 1.2014 g (0.01 mol) and 30.0 mL of ethanol, 1.3614 g (0.01 mol) of p-anisaldehyde was added. The resulting solution was stirred for 8 h and then concentrated under reduced pressure. Distilled water was added and a white solid was formed. The solid was filtered off, washed thoroughly with diethyl ether and water and dried in vacuum. Yield: 90%, mp: 125 °C.

2.2. Infrared Spectroscopy

Infrared (IR) spectrum was recorded on a SHIMADZU Prestige 21 FTIR spectrophotometer. The spectrum of (**1**), in the form of KBr pellets, was recorded in the region 4000 – 400 cm^{-1} .

2.3. Single crystal X-ray diffraction

Suitable crystals of (**1**) were grown by slow lower the temperature in the ethanol. Intensity data were measured with the crystal at room temperature (298 K) using the Enraf-Nonius Kappa-CCD diffractometer monochromated by graphite, with MoK_{α} radiation ($\lambda = 0.71073 \text{ \AA}$). The cell refinements were performed using the software Collect [15] and Scalepack [16] and the final cell parameters were obtained on all reflections. Data reduction was carried out using the software Denzo-SMN and Scalepack [16]. The structure

45 was solved and refined using the software SHELXS97 and SHELXL97, respectively [17].
46 Non-hydrogen atoms of the molecules were unambiguously solved and full-matrix least-
47 squares refinement of these atoms with anisotropic thermal parameters was carried on.
48 The H atoms were located from the difference electron density synthesis and allowed to
49 ride on their parent atoms, with CH(aromatic)= 0.95 Å and CH(aliphatic)= 0.97 Å and
50 $U_{iso}(H) = 1.5U_{eq}$ for methyl H atoms or $1.2U_{eq}$ for the remaining H atoms. Tables were
51 generated by WinGX [18] and the structure representations by ORTEP-3 [18] and Mer-
52 cury [19]. The intramolecular parameters were analyzed using Mogul [20], a knowledge
53 database of molecular geometry derived from Cambridge Structural Database CSD [21],
54 which provides access to information on the ideal values of bond lengths, valence angles
55 and acyclic torsion angles of similar molecular fragments.

56 3. Theoretical methodology

57 To study the structural and electronic properties of the isolated molecule (gaseous
58 phase) we performed quantum mechanics calculations. These calculations were carried out
59 with the *ab initio* density functional theory (DFT) methodology [22–24] together with the
60 6-311++G** basis-set considering both polarization (**) and diffuse (++) functions and
61 the B3LYP exchange-correlation functional [25] as implemented in the Jaguar package [26].
62 As input, the molecular geometry determined by the X-ray diffraction analysis (see section
63 2.3 for details) was used and then optimized. As this conformation is not necessary in the
64 lower energy conformation, all properties presented here were calculated over the optimized
65 structure.

66 Using the Janak's theorem [27] it is possible to estimate the chemical potential (μ),
67 the molecular hardness (η) [28] and the electrophilicity index (ω) [29] from the HOMO and
68 LUMO energies ϵ_H and ϵ_L , respectively as following:

$$\mu \cong \frac{\epsilon_L + \epsilon_H}{2} \quad (1)$$

$$\eta \cong \frac{\epsilon_L - \epsilon_H}{2} \quad (2)$$

$$\omega = \frac{\mu^2}{2\eta} \quad (3)$$

69 To determine the ADME properties the simulations were done with the QikProp soft-
 70 ware [30]. When performing an evaluation, QikProp quickly analyzes the type of atoms
 71 and its charges, and angles, volume and surface area of the molecule. QikProp then uses
 72 these information together with the calculated physical descriptors, using Monte Carlo sta-
 73 tistical mechanics in the regression equations and semi-empirical calculations. The result
 74 is a precise prediction of the pharmacologically relevant properties of a molecule.

75 To study the structural and electronic properties of the solid phase, the calculations
 76 were done considering the DFT methodology for periodical systems [24; 31]. In the case
 77 of organic molecular crystals, one of the main interaction that maintain the cohesion of
 78 the crystal is the hydrogen interaction [32; 33]. This type of non-chemically bonded,
 79 nonlocal and long-ranged interaction is often named as “van der Waals” (vdW) interaction.
 80 Until a few years ago, the mostly used DFT approximations implemented in the solid
 81 phase simulations software were the local density approximation (LDA) and the generalized
 82 gradient approximation (GGA). Since the beginning of the last decade, the first works
 83 appear with new approximations including van der Waals dispersion forces [34–39].

84 We used the software SIESTA [40] together with the van der Waals density functional
 85 as proposed by Klimeš *et al.* [36]. Norm conserving Troullier–Martins pseudopotentials [41]
 86 were generated with the ATOM software (part of the SIESTA package) to avoid the explicit
 87 treatment of the electrons. For the valence electrons, we used a split-valence double-zeta
 88 basis set with polarization functions (DZP) [42]. The structure optimizations were done
 89 until the Hellman-Feynman forces were below $0.04 \text{ eV}\text{\AA}^{-1}$. To obtain results sufficiently
 90 accurate, convergence studies were done for the mesh cutoff energy and the number of
 91 \mathbf{k} -points. The total energy convergence for the system was obtained for a mesh cutoff of
 92 300 Ry and for a \mathbf{k} -point set of $8 \times 8 \times 8$. The Brillouin zone was sampled following the

93 Monkhorst and Pack scheme [43].

94 To better understand the intermolecular interactions, the use of tables with the inter-
95 atomic distances is widely used. This methodology is difficult to interpret, suffering of a
96 natural lack of information. Here we decided to use the Hirshfeld surface and fingerprint
97 tools [44–48] to easily visualize and analyze the intermolecular interactions and the packing
98 behavior.

99 The purpose of the Hirshfeld surface is to define the molecule occupied space in a crystal
100 from the partition of electron density into molecular fragments. Using this idea, a weight
101 function ($w_A(\mathbf{r})$) for a molecule A in a crystal is defined as:

$$w_A(\mathbf{r}) = \rho^{\text{promolecule}}(\mathbf{r}) / \rho^{\text{procrystal}}(\mathbf{r}), \quad (4)$$

102 where $\rho^{\text{promolecule}}(\mathbf{r})$ is the sum of the electron density over the atoms in the molecule A
103 (promolecule) and $\rho^{\text{procrystal}}(\mathbf{r})$ is the sum of the electron density over the crystal (pro-
104 crystal). The Hirshfeld surface is obtained setting $w_A(\mathbf{r}) = 0.5$.

105 4. Results and discussion

106 4.1. Structural properties

107 The crystal data and refinement parameters are presented in Table 1. Crystallographic
108 data for **(1)** has been deposited in the Cambridge Crystallographic Data Centre as a
109 supplementary publication, CCDC789276.

110 Figure 2 shows an ORTEP representation of the molecular structure of **(1)**. An overlay
111 of the molecular backbones in both the experimental and the theoretical intramolecular
112 structures of **(1)** clearly shows the conformational similarity between them. The r.m.s.
113 deviation between analogous non-H atoms is 0.015 Å.

114 The Mogul [20] analysis revealed that all bond lengths correspond with the expected
115 values for a good X-ray diffraction structure refinement. There are two molecular moieties,
116 which are individually almost flat: the moiety 1 containing the atoms O2, C15 and N1

Table 1: Crystal data and structure refinement for (1).

Empirical formula	$C_{16}H_{15}N_1O_2$
Formula weight ($gmol^{-1}$)	253.29
Temperature (K)	298(2)
Wavelength (\AA)	0.71069
Crystal system	monoclinic
Space group	$P2_1/c$
Unit cell dimensions:	
a(\AA)	16.553(5)
b(\AA)	7.928(5)
c(\AA)	10.581(5)
β ($^\circ$)	107.727(5)
Volume (\AA^3)	1322.6(11)
Z (dimer/cell)	4
$D_{(calc)}$ (Mg/m^3)	1.272
Absorption coefficient (mm^{-1})	0.084
F(000)	536
Crystal size	$0.04 \times 0.21 \times 0.24 \text{ mm}^3$
Theta range for data collection ($^\circ$)	3.27 to 25.00
Index ranges	$-19 \leq h \leq 18, -9 \leq k \leq 9, -12 \leq l \leq 11$
Reflections collected	7904
Independent reflections	2325 [R(int) = 0.0443]
Completeness to theta = 0.50° (%)	0.0
Refinement method	Full-matrix least-squares on F^2
Data / restraints / parameters	2325 / 0 / 172
Goodness-of-fit on F^2	1.024
Final R indices [I > 2sigma(I)]	R1 = 0.0485, wR2 = 0.1246
R indices (all data)	R1 = 0.0732, wR2 = 0.1419
Largest diff. peak and hole (\AA^3)	0.212 ⁷ and -0.159

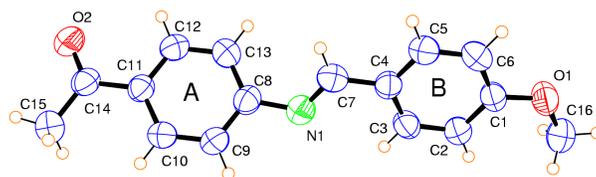


Figure 2: ORTEP representation of (1).

117 and those ones present in rings A and the moiety 2 containing ring B and N1, O1, C7
 118 and C16. The largest deviations from the least squares plane through the moieties 1 and
 119 2 are $-0.124(2)$ and $-0.054(2)$ Å for N1 and C7 atoms, respectively. The least squares
 120 planes of the moieties 1 and 2 form an angle of $50.52(6)^\circ$. This geometric feature is also
 121 highlighted by the torsional angle considering C7-N1-C8-C13 (torsion= $45.08(1)^\circ$). The N1
 122 atom, which belongs to both planar moieties, is slightly closer to the moiety 2 ($0.039(1)$ Å)
 123 than the moiety 1 ($-0.124(2)$ Å).

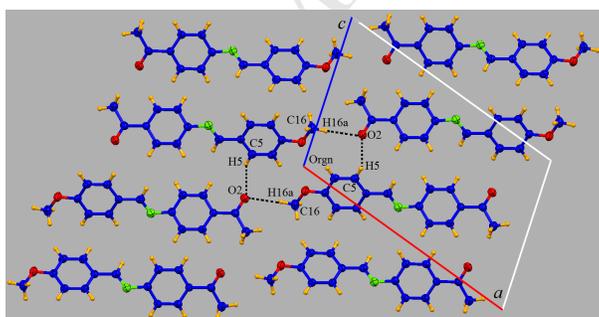


Figure 3: Non-classical hydrogen bonds of (1) forming chains along [101] direction.

124 The supramolecular analysis of (1) shows that there are two non-classical hydrogen
 125 bonds contributing to the stabilization of the crystal packing (figure 3). Either the C5H5
 126 or C16H16A groups act as intermolecular hydrogen bond donors to the O2, forming a
 127 bifurcated hydrogen bond that gives rise to chains along the [101] direction. Parallel
 128 chains are linked together by van der Waals interactions forming an infinite two-dimensional
 129 network parallel to the plane (010). No significant $\pi - \pi$ interactions are observed and the
 130 parallel planar 2D network on the (010) planes are linked along the [010] direction by van

131 der Waals interactions completing the 3D packing.

132 Starting from the X-ray structure, it were performed two types of structure optimiza-
 133 tions. The first one was done setting the cell parameters fixed and letting the atoms free to
 134 move (hereafter called Type-I structure). The second optimization was considering the cell
 135 parameters and the atoms free to move (hereafter called Type-II structure). The crystal
 136 structure parameters are summarized in Table 2.

Table 2: Crystal parameters as obtained from X-ray analysis and for the optimized crystal.

Parameter	Type-I	Type-II
a (Å)	16.553(5)	15.9987
b (Å)	7.928(5)	7.5556
c (Å)	10.581(5)	9.6178
α	90.0	90.0
β	107.727(5)	103.7432
γ	90.0	90.0
$Volume$ (Å ³)	1322.6(11)	1129.314

137 From Table 2 we can see that Type-II (full optimized) structure is a more dense struc-
 138 ture. The Type-II structure has lower values for the cell parameters and the angles did
 139 not present appreciable variations. As a result, the full optimized structure maintain its
 140 monoclinic symmetry and the decreased in volume after the optimization gave place to a
 141 more compact structure.

142 The calculated and experimental infrared spectra are shown in Fig. 4. All calculated
 143 oscillation frequencies were positive (red ticks in Fig. 4), indicating that a minimum energy
 144 optimized structure was reached.

145 The solid state experimental IR spectrum (upper solid blue line in Fig. 4), shows high
 146 intensity bands at 1670 cm^{-1} , assigned to C=N stretch characteristics of the imines or
 147 Schiff's base and the bands at 1585 cm^{-1} and 1510 cm^{-1} are assigned to C=C stretches of

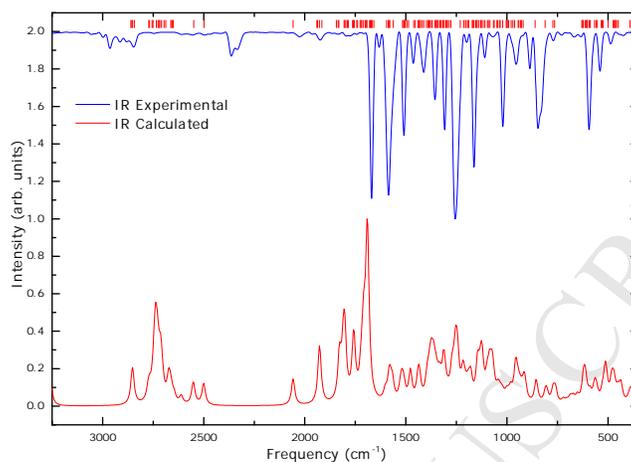


Figure 4: Infrared spectra. Experimental spectrum (upper solid blue line), theoretical spectrum (down solid red line) and calculated frequency positions (red ticks).

148 the aromatic rings. Two strong bands appears in 831 cm^{-1} and 846 cm^{-1} . These bands
 149 are seen in the typical spectrum of para-disubstituted rings present in **(1)**. In the 2800 to
 150 3500 cm^{-1} region, several bands assigned to C-H stretching of the methyl groups appear.

151 The bands observed at 2330 cm^{-1} and 1700 cm^{-1} are assigned to $\nu(\text{CO})$ from CO_2
 152 and ketonic group, respectively. Also the larger band at 3520 cm^{-1} was assigned to $\nu(\text{OH})$
 153 of water, contaminants hard to eliminate by the KBr pellets technique.

154 The theoretical spectrum was generated from the calculated peak positions/intensities
 155 using Gaussian functions with 100 cm^{-1} width. It is well known that *ab initio* vibrational
 156 frequencies are typically larger than the experimental ones [49; 50]. Taking this shift into
 157 account, Fig. 4 shows a good correspondence between experimental and calculated IR
 158 spectra in the peaks positions. This means that the theoretical methodology used here is
 159 suitable to predict the molecular properties.

160 4.2. ADME properties

161 *In silico* predictions of the ADME properties had become a very useful tool decreasing
 162 costs in the discovery and development of new drugs. Consequently, eliminating compounds

163 with high-risk or bad ADME properties earlier in the drug production pathway has become
 164 an attractive approach [51]. Alternatively, the ADME properties could guide scientists to
 165 suggest modifications based on the structure-activity relationship.

166 To know the drugability of (1), the ADME properties were simulated using the QikProp
 167 software [30]. Some of the most used descriptors are shown in Table 3. As it can be seen,
 168 all the parameters are in agreement with the statistically favorable values. It is noteworthy
 169 that no violation occurs to the Lipinski's rule of five [52] and Jorgensen's rule of three [53],
 170 widely used descriptors in drugability predictions.

Table 3: Some of the calculated ADME properties of compound (1).

Property	Value	Recommended
MW ¹	253.300	130.0 – 725.0
QPlogPoct ²	11.553	8.0 – 35.0
QPlogPw ³	6.020	4.0 – 45.0
QPlogPo/w ⁴	3.259	2.0 – 6.5
QPlogS ⁵	-3.993	-6.5 – 0.5
RuleOfFive ⁶	0	maximum is 4
RuleOfThree ⁷	0	maximum is 3

¹ Molecular weight of the molecule. ² Predicted octanol/gas partition coefficient.

³ Predicted water/gas partition coefficient. ⁴ Predicted octanol/water partition coefficient. ⁵ Predicted aqueous solubility, $\log S$. S in $\text{mol} \cdot \text{dm}^3$ is the concentration of the solute in a saturated solution that is in equilibrium with the crystalline solid. ⁶ Number of violations of Lipinski's rule of five. ⁷ Number of violations of Jorgensen's rule of three.

171 Together with the calculation of ADME descriptors, the QikProp package is able to
 172 identify similar drug molecules. In this case, the top 4 compounds with similarity greater
 173 that 90% were Melperone[®] (used to treat dementia, sleep disorders, alcohol dependence

174 and psychosis), Medifoxamine[®] (an antidepressant which acts as a dopamine reuptake
 175 inhibitor), Fenbufen[®] and Ketoprofen[®] both used as nonsteroidal anti-inflammatory drug
 176 (NSAID). These results point out **(1)** as a promise candidate for treating these diseases.

177 4.3. Electronic properties

178 The highest occupied molecular orbital (HOMO) and the lowest unoccupied molecular
 179 orbital (LUMO) are shown in Fig. 5. From the diagram we can see that these orbitals are
 180 asymmetrically spread over the entire molecular structure. This is a desirable characteristic
 181 for molecular electronic since the device can behave differently with different bias.

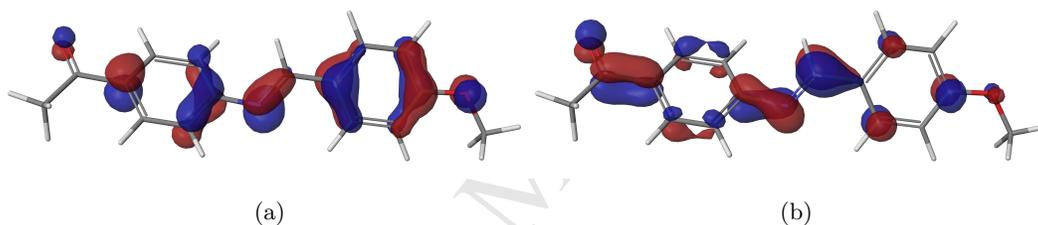


Figure 5: Molecular orbitals (a) HOMO and (b) LUMO (diagrams rendered with the Maestro software [54]).

182 The calculated values of chemical potential (μ), the molecular hardness (η) and the
 183 electrophilicity index are shown in Table 4. From them, we can infer that as η has a posi-
 184 tive value, the redistribution of the electrons in the molecule is energetically unfavorable.
 185 Also, the higher η value is, the more stable the molecule is and, therefore, the harder the
 186 rearrangement of its electrons [55]. The electrophilicity index, ω , can be used as a measure
 187 of the molecule energy lowering due to the maximal electron flow between the environ-
 188 ment and the molecule [29]. This is a very useful tool when studying the ligand-binding
 189 phenomena in drug design: higher values of ω implies lower energies for the ligand (more
 190 stability).

191 In Fig. 6 the Hirshfeld surface together with the hydrogen interaction for Type-II struc-
 192 ture is shown. The figure stands for a mapping of d_{norm} over the Hirshfeld surface. The

Table 4: Calculated parameters ϵ_L , ϵ_H , μ , η and ω .

Property	Value(eV)
ϵ_L	-2.183
ϵ_H	-6.221
μ	-4.202
η	2.019
ω	4.372

193 function d_{norm} represents a symmetric function of distances from the inside and outside
 194 of the Hirshfeld surface nuclei (d_i and d_e , respectively), relative to their respective van
 195 der Waals radii. The red regions represent distances shorter than vdW separations (high-
 196 lighting both donor and acceptor equally) whereas white regions represent distance equal
 197 to vdW and blue regions for distances greater than vdW. The red spot in Fig. 6 indicate
 198 that the van der Waals interactions are the main responsible for the crystal stability and
 199 cohesion, therefore a more compact structure.

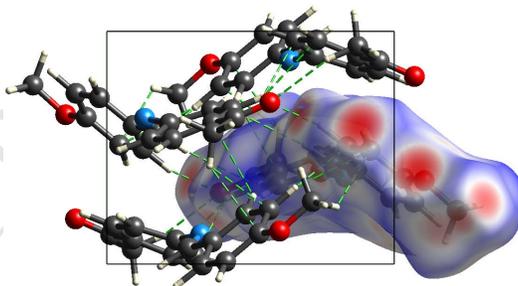


Figure 6: Hirshfeld surface and hydrogen interactions (dashed green lines). Diagram rendered with the *CrystalExplorer* software [56].

200 Another useful tool to visualize the distribution and type of interactions inside the
 201 crystal is the fingerprint plot. The fingerprint plots correspond to 2D histograms for the
 202 distribution of (d_i, d_e) pairs from the Hirshfeld surface. In Fig. 7 the calculated fingerprint

203 plots for Type-II structure are shown.

204 In Fig. 7a the contributions of all atoms from the molecule are represented. At first
 205 sight, it can be seen that the intermolecular contact distribution is below 2.2\AA . This is
 206 in accordance to the fact that Type-II structure is more compact than Type-I were the
 207 contact distances extend to more than 2.5\AA . The sharp spikes in the plots characterize the
 208 hydrogen interactions: longer spikes represent stronger hydrogen bonds, the upper spike
 209 corresponds to the hydrogen interaction donor ($d_e > d_i$) and the other one to the hydrogen
 210 interaction acceptor ($d_i > d_e$).

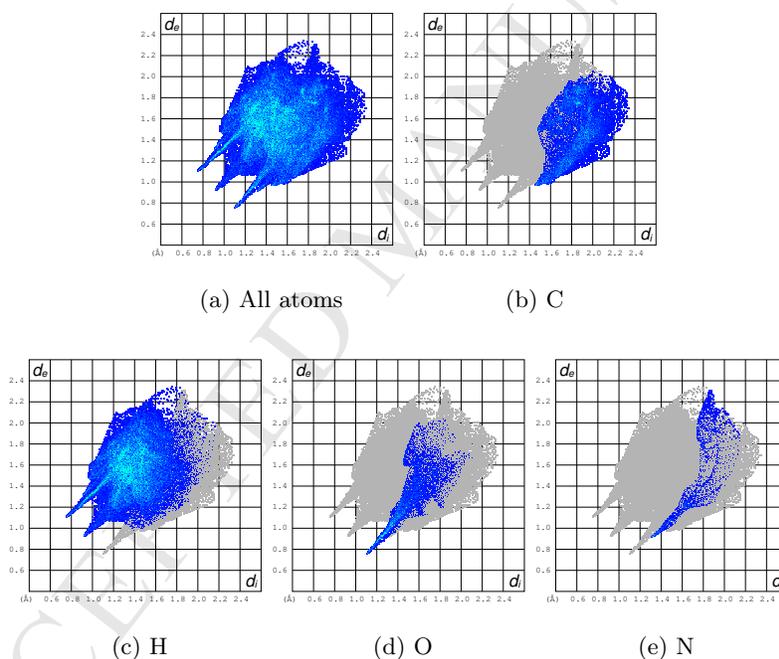


Figure 7: Fingerprint plots for Type-II structure. Maps for atoms in the promolecule (a) all atoms; (b) carbon atoms; (c) hydrogen atoms; (d) oxygen atoms and (e) nitrogen atoms. Diagrams rendered with the *CrystalExplorer* software [56].

211 To get a better understanding about what atoms contribute more to the crystal cohesion
 212 and how they do that, it is possible to plot the fingerprints for a given type of atom from
 213 the promolecule, representing in this way, the interaction of it with all other atoms from

214 the procrystal. In both structures, the carbon and hydrogen atoms are responsible for
 215 more than 80% of the interactions. Their contributions are shown in figures 7b (carbon
 216 atoms) and figures 7c (hydrogen atoms). Figure 7d and Fig. 7e shown the contributions
 217 of oxygen and nitrogen atoms, respectively. Both atoms behave like hydrogen interaction
 218 acceptors.

219 One of the properties that can be calculated within the quantum mechanical ap-
 220 proach is the partial atomic charge distribution. To accomplish this task, there are several
 221 methodologies like (i) Mulliken population analysis [57–60], (ii) natural population analysis
 222 (NPA) [61], (iii) the Breneman-Wiberg (BW) model [62], (iv) Merz-Kollman-Singh (MKS)
 223 electrostatic potential derived charges [63; 64], (v) the Bader partition scheme [65–68],
 224 (vi) the Voronoi partition scheme [69; 70] and (vii) the Hirshfeld partition scheme [71].
 225 However, none of them is universally accepted as the “best” for computing partial atomic
 226 charges. Here, we decided to use the Hirshfeld partition scheme as it proves to be highly
 227 insensitive to the choice of the basis set [71; 72]. In Fig. 8 the calculated charges for Type-II
 228 structure are shown.

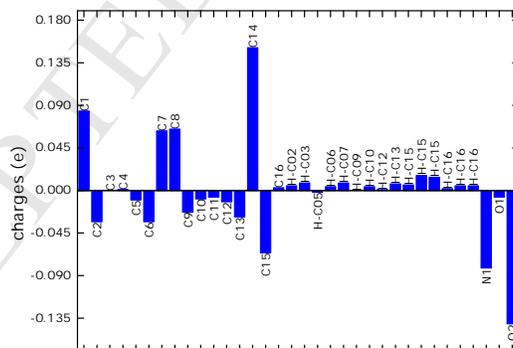


Figure 8: Charge distribution using the Hirshfeld [71] partition scheme for Type-II structure.

229 In Fig. 9 the calculated total density of states (DOS) together with each atom contri-
 230 bution are shown. The systems did not show spin polarization. The states over the Fermi
 231 energy (E_F , dashed green line) are the states available in the conduction band and the

232 states below the Fermi energy represents the states from the valence bands. As it can be
 233 seen, the conduction states are formed mainly with carbon orbitals whereas the valence
 234 band states are formed with the contribution of carbon and oxygen orbitals.

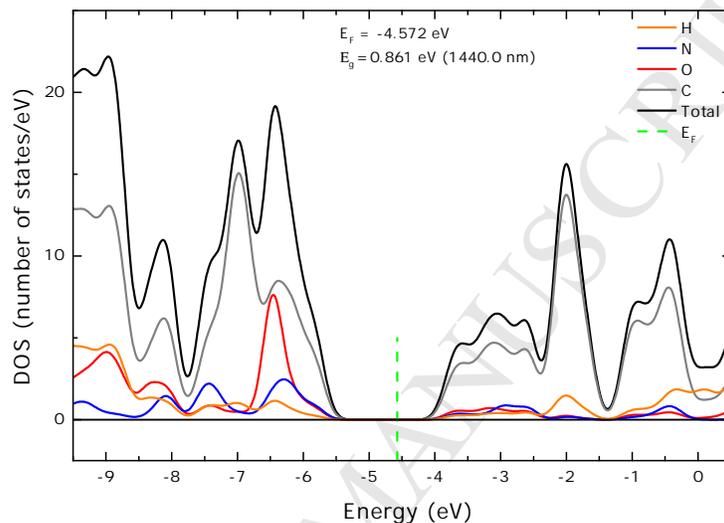


Figure 9: Total density of state (DOS, solid black line) for Type-II structure and individual atoms contributions: carbon (solid gray line), oxygen (solid red line), nitrogen (solid blue line) and hydrogen (solid orange line).

235 A visible energy gap (E_g) equal to equal to 0.861 eV appears. From an optical point
 236 of view, this energy gap implies that **(1)** is transparent for electromagnetic radiation with
 237 energies below 0.861 eV and have available states for the absorption of radiation with ener-
 238 gies greater that 0.861 eV including part of the infrared region, the visible and ultraviolet
 239 regions. These properties indicate that **(1)** can be used as a low gap organic semiconductor
 240 very useful in optoelectronic applications.

241 4.4. Gaseous vs solid

242 The structural and electronic properties of an isolated molecule (gaseous form) are
 243 very different when the same molecule is in its crystal form. This is due to the collective
 244 contribution of each molecule and the geometrical restriction imposed by crystal periodicity.

245 The title compound (**1**), in its gaseous phase, has the lower possible symmetry (C_1) for
 246 a molecule. In its crystal form, the system belongs to the $P2_1/c$ space group, which is also
 247 a low-symmetry space group. A low-symmetry system has low geometrical restrictions
 248 which are translated in fewer limitations for the physical properties [73].

249 In table 5, the calculated root-mean-square deviation (RMSD) and maximum distance
 250 between two equivalent atoms (MaxD) using the molecule overlay feature of Mercury [19]
 software is shown for the three structures.

Table 5: Calculated RMSD/MaxD (in Å)

	Gaseous	Type-I	Type-II
Gaseous	0	0.8645/2.2184	0.2919/0.5968
Type-I	0.8645/2.2184	0	0.8460/2.4884
Type-II	0.2919/0.5968	0.8460/2.4884	0

251

252 The lower values for the RMSD and MaxD is for the pair gaseous/Type-II structures.
 253 This is expected as the Type-II was optimized without any constrains, permitting the full
 254 variations of atoms position together with cell parameters.

255 In Fig. 10 the calculated infrared spectra for the gaseous and solid phases (Type-II
 256 structure) are shown.

257 The infrared spectra of molecules in the gaseous phase are characteristic of a free
 258 molecule without any kind of interaction. In the solid phase, the crystalline order has
 259 two effects. First, the symmetry of each molecule can be changed due to distortions in
 260 its structure owing to the interactions with its nearest neighbors. The second effect is
 261 the appearance of new oscillation modes due to the coupling (in-phase or out-of-phase)
 262 between neighboring molecules in the crystal unit cell [74].

263 The oscillation modes of (**1**) in its gaseous phase are concentrated in the regions
 264 1750 cm^{-1} to 250 cm^{-1} and over 3000 cm^{-1} . The first region of oscillations is also visi-
 265 ble in the spectrum for the solid phase whereas the oscillations in the second region were

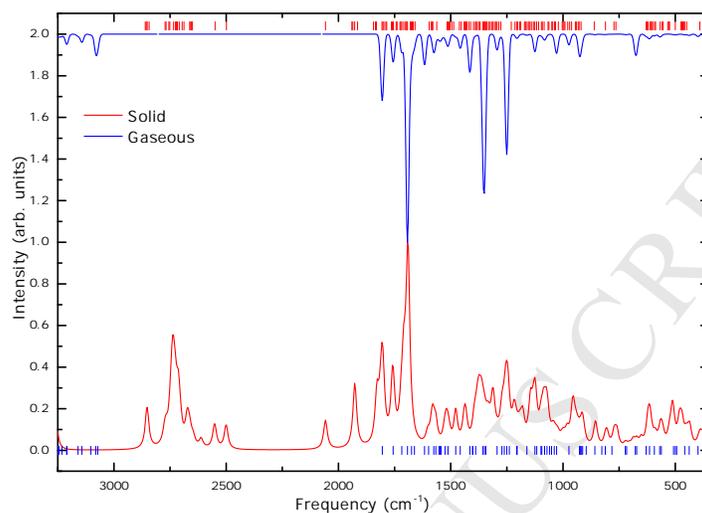


Figure 10: Infrared spectra. Gaseous phase (upper blue lines), solid phase (Type-II structure, down red lines) and calculated frequency positions (ticks).

not reproduced at the solid phase. The appearance of the other bands in the solid phase spectrum were associated with the coupling of the 4 molecules present in the unit cell.

The molecular orbitals of a single molecule are formed from the superposition of atomic orbitals and can be calculated using the linear combinations of atomic orbitals (LCAO) approximation [6]. The results are discrete energy levels. In the solid phase (organic crystals), the overlap of atomic and molecular orbitals produce bands instead discrete energy levels. In this case, the translational symmetry of the crystal gives rise to another quantum number \mathbf{k} that represents the momentum of the electrons.

The results from the Natural Bond Orbital (NBO) analyzes [75; 76] (using the Canonical Molecular Orbital Composition) for the HOMO and LUMO orbitals are summarized in table 6. The HOMO has significant contributions from at least 9 NBOs whereas the LUMO has a contribution from 6 NBOs. All the participant NBOs has a p character. The results obtained from NBO analysis are in agreement with the orbital distribution (see inset of Fig. 11).

Table 6: Natural bond orbital analysis for HOMO and LUMO

HOMO	LUMO
0.395*[33]: BD(2) C3-C7	0.536*[72]: BD*(2) N2-C5*
0.356*[29]: BD(2) N2-C5	0.373*[90]: BD*(2) C8-O19*
0.337*[42]: BD(2) C6-C11	0.316*[80]: BD*(2) C4-C13*
0.298*[21]: LP(2) O1(lp)	-0.311*[76]: BD*(2) C3-C7*
-0.290*[37]: BD(2) C4-C13	0.273*[68]: LV(1) C14(lv)
0.272*[23]: LP (1) C10(lp)	-0.247*[50]: BD(2) C9-C16
-0.269*[100]: BD*(2) C12-C15*	
0.248*[68]: LV(1) C14(lv)	
0.236*[93]: BD*(2) C9-C16*	

BD: bond, BD*: antibond, LP: lone pair (unfilled valence-shell), LV: lone vacancy (unfilled valence nonbonding).

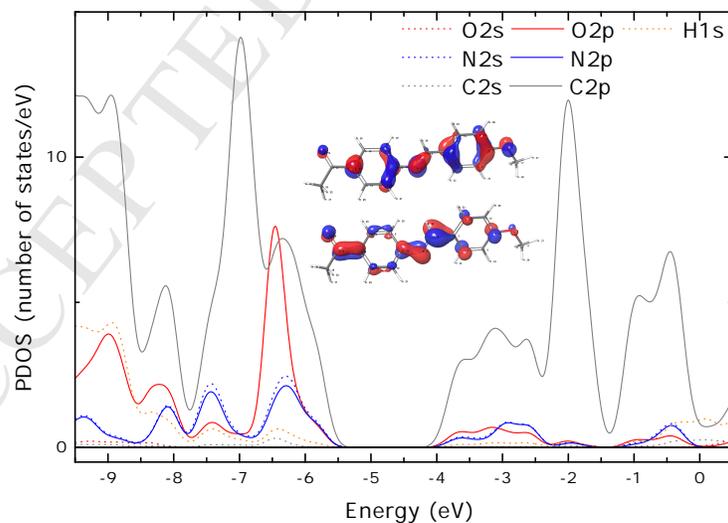


Figure 11: Partial density of states for Type-II system and HOMO/LUMO molecular orbitals.

280 For the solid state phase, the partial density of states is shown in Fig. 11. As can be
281 seen, the main contributions to the conduction and valence bands are from the carbon
282 atoms. For the valence band, there is also a contribution from the $2p$ orbitals from oxygen
283 and nitrogen atoms.

284 Comparing both results (table 6 and Fig. 11) it is possible to affirm that the same type
285 of atoms contributes for the HOMO/LUMO and valence/conduction bands, respectively.

286 5. Conclusions

287 The structural and electronic properties of compound (**1**) were determined in gaseous
288 and solid phases from experimental and theoretical approaches.

289 In this case, we obtain a good correspondence between experimental and calculated IR
290 spectra indicating that the theoretical methodology is adequate to simulate the properties
291 of these compound. The calculated molecular hardness (η) and the electrophilicity index
292 (ω) show an electronically stable molecule which is suitable to bind or interact with other
293 chemical species without appreciable electronic modifications. The calculation of ADME
294 properties indicate that (**1**) has drug like properties and from similarity with other com-
295 mercial drugs, it is a promise candidate to act as a nonsteroidal anti-inflammatory and to
296 treat dementia, sleep disorders, alcohol dependence and psychosis.

297 For the solid phase, the crystal was grown and its structure was experimentally de-
298 termined using single crystal X-ray diffraction. Starting with the experimental crystal
299 structure, two structure optimizations were done considering van der Waals dispersion
300 forces, very important in organic crystals. The type of interactions between atoms and
301 their contribution to the crystal cohesion were graphically analyzed using the Hirshfeld
302 surface and fingerprints. The calculated density of states show that (**1**) is a low gap semi-
303 conductor and can act as an absorber for electromagnetic radiations with energy greater
304 that ~ 0.9 eV.

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Paper: "Synthesis, characterization and theoretical study in gaseous and solid phases of the imine 4-Acetyl-N-(4-methoxybenzylidene)aniline"

Authors: J.F.N. Batista, J.W. Cruz Jr., A.C. Doriguetto, C. Torres, E.T. de Almeida and I. Camps

Journal: Journal Molecular Structure

- The calculated chemical descriptors indicates an electronically stable molecule.
- Title compound has pharmacological drug like properties.
- Solid phase of title compound is a low gap semiconductor.
- Carbon and oxygen atoms contribute must to HOMO/LUMO and electronic bands.