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# Doping dependence of magnetic and transport properties in single crystalline Co-doped $BaFe_2As_2$

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#### ABSTRACT

We report the doping dependence of transport and magnetic properties in Co-doped BaFe<sub>2</sub>As<sub>2</sub>. With increasing Co concentration *x*, structural and magnetic transitions are suppressed and superconductivity emerges in the range of 0.3 < x < 0.15. *T*-linear resistivity is observed at the optimally doped composition x = 0.075 and the temperature exponent of the resistivity increases with *x*. Critical current density  $J_c$  at low temperatures and low fields obtained from bulk magnetization is reasonably large and the doping dependence shows a maximum at  $x \sim 0.075$  similar to  $T_c$ . The values of  $J_c$  at low temperatures reach about  $1 \times 10^6$  A/cm<sup>2</sup> around the optimally doped region, which is potentially attractive for technological applications.

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#### 1. Introduction

Since the discovery of the high- $T_c$  iron-based oxypnictide superconductor LaFeAsO<sub>1-x</sub> $F_x$  with  $T_c \sim 26$  K [1], other iron-based superconductors have attracted great interest. Following this discovery, oxygen-free iron-arsenide  $AFe_2As_2(A = Ba, Sr, Ca)$  was discovered. These materials show superconductivity with the substitution of alkali metals, such as Na, K, and Cs, for A resulting in the introduction of holes in the (FeAs)<sup>-</sup> layers [2], or Co for Fe resulting in the introduction of electrons [3]. Among them, Codoped BaFe<sub>2</sub>As<sub>2</sub> has attracted much attention because the superconductivity induced by the substitution in the conducting layers strongly compares well with a drastic suppression of  $T_c$  in cuprates [4] and may suggest a possibility of novel superconducting state. We report here the doping dependence of transport and magnetic properties in Co-doped BaFe<sub>2</sub>As<sub>2</sub> in order to investigate normal and superconducting properties. With increasing Co concentration x, structural and magnetic transitions are suppressed and superconductivity occurs in the range of 0.03 < x < 0.15. *T*-linear resistivity is observed at the optimally doped composition near x = 0.075 and temperature exponent of resistivity crosses over to  $T^2$  in the overdoped region. Critical current density  $I_c$  at low temperatures and low field obtained from bulk magnetization is relatively large and shows a doping dependence similar to  $T_c$ . The values of  $I_c$  reach about  $1 \times 10^6$  A/cm<sup>2</sup> at 2 K around the optimally doped region.

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#### 2. Experiments

Single-crystalline samples of Ba $(Fe_{1-x}Co_x)_2As_2$  were grown by the FeAs/CoAs self-flux method. FeAs and CoAs were prepared by placing mixtures of As pieces and Fe/Co powder in a silica tube and reacting them at 1065 °C for 10 h after heating at 700 °C for 6 h. A mixture with a ratio of Ba: FeAs/CoAs = 1:5 was placed in an alumina crucible. The whole assembly was sealed in a large silica tube, and heated up to 1150 °C for 10 h followed by slow cooling down to 800 °C at a rate of 5 °C/h, which is slightly different from the synthesis reported before [5]. After cleaving, we can obtain shiny samples. The typical dimensions of the resulting crystals is  $4 \times 4 \times 0.1$  mm<sup>3</sup>. Co concentrations were determined by energy dispersive X-ray spectroscopy measurements. Magnetization was measured by a commercial SQUID magnetometer (Quantum Design MPMS–XL5) and resistivity measurement was performed by a four-contact method.

### 3. Results and discussion

Fig. 1 shows the temperature dependence of in-plane resistivity  $\rho_{ab}$  for Ba(Fe<sub>1-x</sub>Co<sub>x</sub>)<sub>2</sub>As<sub>2</sub>. For BaFe<sub>2</sub>As<sub>2</sub>, x = 0, the sharp decrease of  $\rho_{ab}$  associated with the structural and magnetic transition [2,6] is observed at 134 K. With increasing x, temperature of the resistive anomaly is suppressed monotonically. For  $x \ge 0.038$ , superconductivity is observed and disappear at  $x \ge 0.15$ . At x = 0.075,  $T_c$  shows a maximum of ~24 K and very sharp transition within  $\Delta T_c < 1$  K is observed in the resistivity and magnetization as shown in the right inset of Fig. 1. We find *T*-linear resistivity at the optimally doped composition x = 0.075 and the temperature





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**Fig. 1.** Temperature dependence of the in-plane resistivity  $\rho_{ab}$  for Ba(Fe<sub>1-x</sub>Co<sub>x</sub>)<sub>2</sub>As<sub>2</sub>, x = 0, 0.023, 0.038, 0.06, 0.075, 0.113, and 0.15. The right inset shows the temperature dependence of  $\rho_{ab}$  and zero-field-cooled (ZFC) and field-cooled (FC) magnetization for x = 0.075. The left inset shows the phase diagram for Ba(Fe<sub>1-x</sub>Co<sub>x</sub>)<sub>2</sub>As<sub>2</sub>.  $T_s$  and  $T_{SDW}$  were by the temperature derivative of resistive anomaly and  $T_c$  is determined by the onset of diamagnetism. Dashed lines are guides for the eye.

exponent of the resistivity increases with *x*. At x = 0.15,  $T^2$ -behavior is observed below 40 K, which indicates the non-Fermi-liquid to Fermi-liquid crossover with increasing *x*. The doping dependence of resistivity for Ba(Fe<sub>1-x</sub>Co<sub>x</sub>)<sub>2</sub>As<sub>2</sub> are very similar to those reported

by Refs. [7,9]. The Left inset of Fig. 1 shows the phase diagram of  $Ba(Fe_{1-x}Co_x)_2As_2$  obtained from the transport and magnetization measurements, where structural transition temperature  $T_s$  and spin-density wave transition temperature  $T_{SDW}$  were determined by the temperature derivative of resistive anomaly and  $T_c$  is determined by the onset of diamagnetism, which is consistent with the previous study reported in Ref. [8].

The inset of Fig. 2 shows the field dependence of magnetization for x = 0.045, 0.06, 0.075, and 0.113 at several temperatures. For x = 0.045 and 0.113, corresponding to under- and over-doped region, respectively, a maximum at zero field is very sharp. By contrast, for x = 0.06 and 0.075 near the optimally doped region, a peak at zero field is broadened and the irreversible magnetization is much larger than those for x = 0.045 and 0.113. For x = 0.075, prominent fish-tail magnetization is observed at 15 K [5] while for x = 0.045, 0.06, and 0.113, the fish-tail effect is not observed in the present temperature and field range.

Fig. 2 shows the field dependence of critical current density  $J_c$  for x = 0.045, 0.06, 0.075, and 0.113.  $J_c$  is obtained from hysteresis loop in the magnetization shown in each inset using the Bean model,

$$J_c = 20 \frac{\Delta M}{a(1-a/3b)},\tag{1}$$

where  $\Delta M$  is  $M_{down} - M_{up}, M_{up}$  and  $M_{down}$  are the magnetization when sweeping fields up and down, respectively, and *a* and *b* are the sample widths (*a* < *b*). For *x* = 0.075, non-monotonious field dependence of  $J_c$  at high temperatures above 10 K reflecting the fish-tail magnetization is observed. On the other hand, for



**Fig. 2.** Field dependence of critical current density for  $Ba(Fe_{1-x}Co_x)_2As_2$ , (a)  $x = 0.045(330 \times 415 \times 30 \ \mu m^3)$ , (b)  $0.06 (410 \times 520 \times 35 \ \mu m^3)$ , (c)  $0.075 (600 \times 1200 \times 20 \ \mu m^3)$  and (d)  $0.113 (410 \times 540 \times 10 \ \mu m^3)$  obtained from magnetization using the Bean model. Inset shows the field dependence of magnetization at several temperatures.

x = 0.045, 0.06, and 0.113,  $J_c$  decreases monotonically with increasing field. The values of  $J_c$  at 2 K reach about  $1 \times 10^6$  A/cm<sup>2</sup> for x = 0.06 and 0.075 around optimally-doped region, which is relatively high and potentially attractive for technological applications. It should be noted that further enhancement of  $J_c$  by heavy-iron irradiation has been demonstrated [10].

In summary, we study the doping dependence of transport and magnetic properties in Co-doped BaFe<sub>2</sub>As<sub>2</sub>. With increasing Co concentration, structural and magnetic transitions are suppressed and superconductivity emerges in the range 0.03 < x < 0.15. *T*-linear resistivity is observed only near the optimally doped composition x = 0.075 and the temperature exponent of the resistivity increases with *x*.  $J_c$  at low temperatures and fields reaches about  $1 \times 10^6$  A/cm<sup>2</sup> around optimally doped region, which is potentially attractive for technological applications.

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