

# Synchrotron radiography and x-ray topography studies of hexagonal habitus SiC bulk crystals

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Phase-sensitive synchrotron radiation (SR) radiography was combined with x-ray diffraction topography to study structural defects of SiC crystals. The particular bulk SiC crystals examined had a low micropipe density and a hexagonal habitus composed of prismatic, pyramidal, and basal faces well developed. X-ray diffraction topography images of the sliced (0001) wafers, which were formed due to the complex lattice distortions associated with defective boundaries, demonstrated the existence of two-dimensional defective boundaries in the radial direction, normal to the (0001) planes. In particular, those parallel to the  $\langle 11\bar{2}0 \rangle$  directions extended rather far from the seed. On the other hand, by phase-sensitive SR radiography the effect of micropipe collection was detected. Micropipes grouped mostly in the vicinities of the defective boundaries but rarely appeared between groups. Some general remarks about possible reasons for the development of such peculiar defect structures were made.

## I. INTRODUCTION

Silicon carbide (SiC) attracts persistent research interest due to its unique electronic, thermal, mechanical, and other properties. The problem of manufacture of large-size single crystals free from structural defects, however, remains unsolved. The solution requires detailed investigations of the SiC structural quality, which is sensitive to growth conditions. For nondestructive purposes, these are typically done by x-ray methods. X-ray imaging techniques allow the visualization of lattice defects within a crystal interior. Synchrotron radiation (SR) x-ray imaging has advantages caused by very high intensities and good collimations of synchrotron beams. SR diffraction topography has already allowed to show, in SiC wafers of relatively high structural quality, the generation of

hollow core dislocations, or micropipes (MPs), and the formation of other polytype inclusions.<sup>1,2</sup> Low divergence of the synchrotron beam at the Stony Brook Synchrotron Topography Facility made it possible to establish an accurate method to simulate the pure orientation contrast from super screw dislocations.<sup>3</sup> High flux and high energy of the x-ray beams available at the European Synchrotron Radiation Facility (ESRF), together with the large beam size, permitted studying the whole volume of the bulky SiC ingots.<sup>2</sup> Moreover, special properties of the third-generation synchrotron sources, like the ESRF, opened the way to novel imaging techniques like phase contrast radiography.<sup>4,5</sup> Being insensitive to strong deformations specific to sublimated grown SiC crystals of large area, this method allowed one to image MPs in samples of any crystalline perfection.<sup>6,7</sup> Besides the mapping and the determination of the size of the micropipes, phase-sensitive radiography provided information about their shape and spatial

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distribution. Principal differences in image formation mechanisms make x-ray topography and phase-sensitive radiography complementary techniques, and investigations of SiC structural quality will benefit by the use of both methods.

In this work we present such an approach to study SiC bulk crystals that were grown by the sublimation sandwich method described in Ref. 8. A peculiar feature in the studied crystals was a well-defined hexagonal habitus composed of prismatic, pyramidal, and basal faces. The other specific attribute was a very low micropipe density observed over a large area located at a boule periphery. We report the observation of SiC defects using both strain-sensitive (topography) and density gradient sensitive (radiography) x-ray imaging techniques and discuss their nature and peculiarities of distribution in the crystal interior.

## II. EXPERIMENTAL

### A. Instrumental

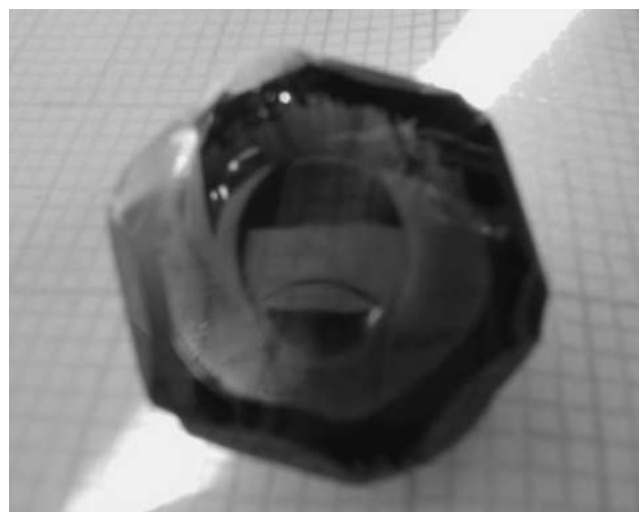
Phase-sensitive SR radiography experiments were performed at the Pohang Light Source (PLS Pohang, Korea). We used the 5C1 beamline, whose 1.32-T bending magnet gives an effective source size (measured at 5.5 keV operated at 2.5 GeV) of the order of 60  $\mu\text{m}$  in the vertical direction and 160  $\mu\text{m}$  horizontally, at a distance of 30 m from the studied sample. The experiments were performed in the edge detection regime with a sample-to-detector distance in the range from few centimeters to 1 m. A model suggested by Margaritondo and Tromba<sup>9</sup> and verified by Hwu *et al.*<sup>10</sup> showed that, in this regime, image sharpening could be achieved under less stringent conditions than those used in pioneering SR phase radiography experiments.<sup>4,5</sup> Namely, on one hand, the longitudinal coherence length is not a rigid requirement and, without too much loss of resolution, a polychromatic beam may be utilized. On the other hand, the requirements concerning the lateral coherence length are rather forgiving and can be fulfilled for SR beam parameters considerably less extreme than those achievable at “long” beamlines. We used unmonochromatized (“white”) light with no optical elements except beryllium windows.

The x-ray images were detected by a slow scan charge-coupled device (CCD) camera after converting the x-rays to visible light with a thin  $\text{CdWO}_4$  scintillator crystal. Different microscope objectives were used to change the field of view between  $4 \times 3 \text{ mm}^2$  and  $0.6 \times 0.4 \text{ mm}^2$ . The 12-bit CCD camera had a  $1392 \times 1040$  pixels matrix size. Resolution of a few micrometers was routinely obtained. Each sequence of measurements was accomplished by a background subtraction. All images displayed here are background-subtracted with no further processing.

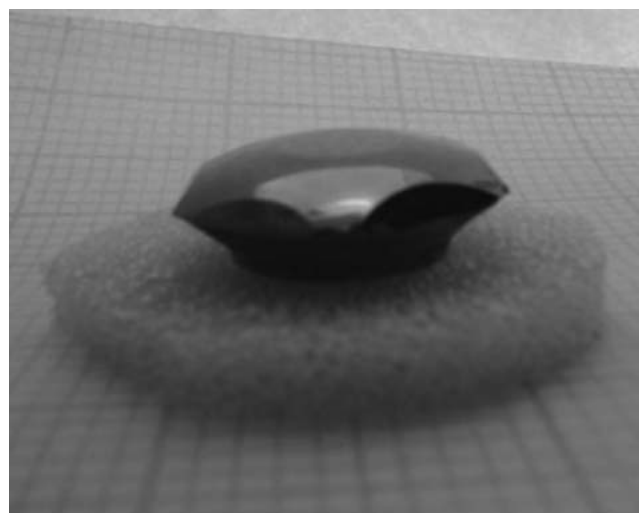
X-ray diffraction topography investigation was done on a commercial source in one crystal set up with the Lang method in Bragg and Laue geometries by using  $\text{Cu K}_\alpha$  and  $\text{Mo K}_\alpha$  radiations. In the plane of scattering, angular divergence of the beam at the specimen was of the order of angular separation between  $\alpha_1$  and  $\alpha_2$  x-ray lines. When the specimen to film distance was kept 5–10 mm, a few micrometers spatial resolution normal to the scattering plane could be achieved.

### B. Samples under study

Bulk SiC crystals were grown by a sublimation sandwich method in a tantalum container.<sup>8</sup> Features of the boules that attracted investigation are evident in the optical photographs presented in Figs. 1(a) and 1(b). In projection on basal plane the crystals were hexagonal in



(a)



(b)

FIG. 1. Photographs of the top (a) and side (b) views of SiC boule of hexagonal habitus.

shape with obtuse corners. The periphery of the boules exhibited remarkable prismatic faces crossed by sections of pyramidal faces, as illustrated in Fig. 2. A central part of the boule which was close to seed in the horizontal direction showed the lack of hexagonality (*l* in Fig. 2). A circular site (*O* in Fig. 2) parallel to the basal plane was placed just above the seed. Crystals grown by this method can be as large as 40 mm in diameter with the height up to approximately 20 mm. Note that the sample shown in Fig. 2 has an approximately 20-mm diameter. A boule near surface area (*l* and *O* in Fig. 2) was studied by x-ray diffraction topography in reflection geometry. For a further x-ray investigation the boules were sliced into wafers of a (0001) basal plane as well as (10 $\bar{1}$ 0) and (11 $\bar{2}$ 0) prism planes.

### III. RESULTS

#### A. X-ray diffraction topography

A typical example of an x-ray diffraction topograph of the (0001)-oriented wafer is shown in Fig. 3. Lines of bright contrast (high local intensity) extending out to the

wafer periphery in radial directions are visible, although their origin is not clear from the topographs alone. By correlating the positions of these lines in the wafers sliced at different distances from the beginning of growth, we suggested that they were basal plane traces

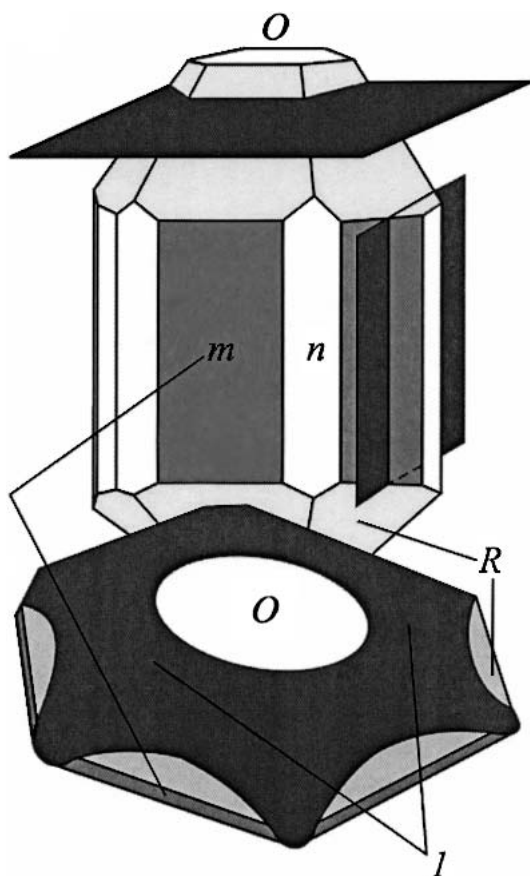


FIG. 2. Scheme of SiC crystal boule of a hexagonal habitus composed by the faces: (0001) (*O*); {10 $\bar{1}$ 0} (*m*); {10 $\bar{1}$ 1} (*R*). *l* is a part of the boule close to seed in the horizontal direction. Boules were sliced parallel to the planes (0001), (10 $\bar{1}$ 0), and (11 $\bar{2}$ 0) (*n*).

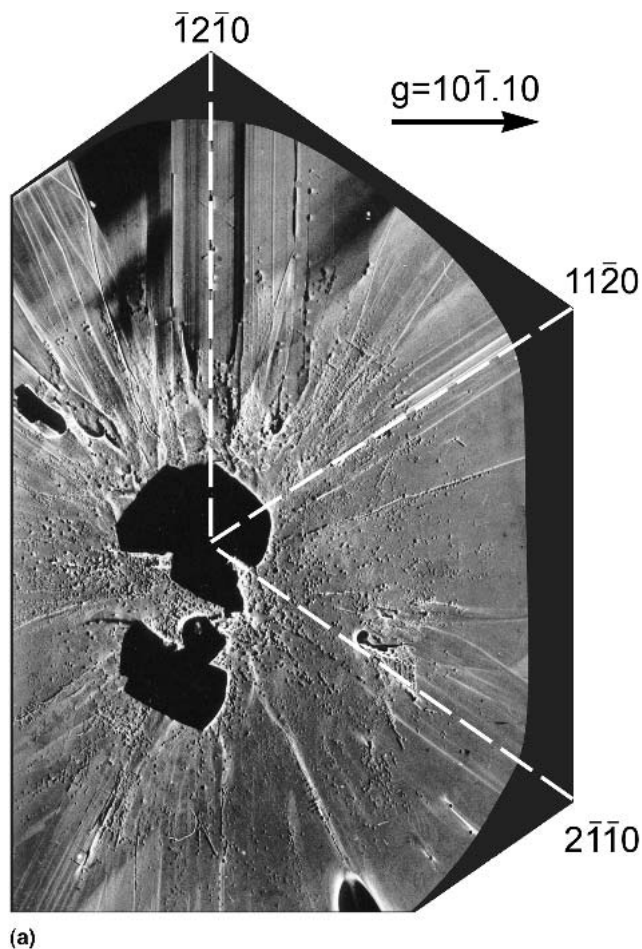


FIG. 3. (a) X-ray diffraction topograph of major area of (0001) SiC wafer: Lang method in Bragg geometry, 10 $\bar{1}$ .10, Cu K $\alpha$ , Bragg angle 35.8° (5 $\times$ ). (b) Higher magnification image: Lang method in Bragg geometry, 10 $\bar{1}$ .10, Cu K $\alpha$  (12 $\times$ ). The contrast is inverted to the original image.



of two-dimensional (2D) defects, or defective boundaries. These traces were produced by diffraction contrast in both basal plane and prism plane reflections due to lattice distortions associated with the defective boundaries.

In particular, the traces parallel to the  $\langle 11\bar{2}0 \rangle$  directions, as magnified in Fig. 3(b), were visible continuously through several basal plane slices, indicating that the corresponding defective boundaries extended rather far from the seed. On the other hand, those defects represented by curved lines in the topograph were mostly lost across the depth of one wafer. Comparing x-ray topographs obtained in the reflection and transmission geometries, such as shown in Fig. 4(a) and 4(b), from the part of the sample displayed in Fig. 3, reveals that these two types of defects indeed differ from each other. The marked traces [Fig. 4(b)] clearly visible in the transmission topograph while invisible in the reflection mode imply that the related boundaries are local.

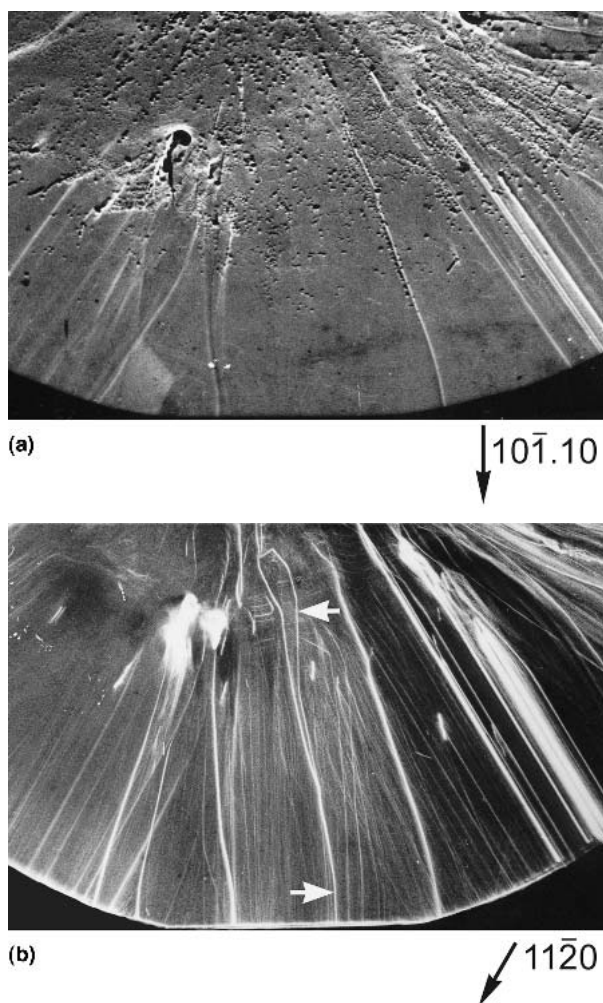


FIG. 4. (a) Reflection and (b) transmission topographs of the sample shown in Fig. 3: (a) Lang method in Bragg geometry,  $10\bar{1}.10$ ,  $\text{Cu K}\alpha$  ( $6\times$ ); (b) Lang technique in Laue geometry,  $11\bar{2}0$ ,  $\text{Mo K}\alpha$ , Bragg angle  $\theta = 13.3^\circ$  ( $6\times$ ).

The spotlike contrasts, as can be seen from the reflection topograph in Fig. 3, are another type of defect, which can be correlated with hollow-core screw dislocations or micropipe defects. Their distribution primarily concentrated in the central, close-to-seed area. A large area over the wafer periphery, on the other hand, was almost free from such diffraction spots [Fig. 3(b)]. With scanning electron microscopy etch pits associated with hollow-core screw dislocations in molten KOH treated wafers (data not shown) were observed in a very similar distribution. The dislocation density was found high in the central area but decreased with distance from the center. The density was also high close to the traces of the  $\langle 11\bar{2}0 \rangle$  defective boundaries with little in between. The resemblance therefore strongly suggested that spotlike contrasts we observed in the topograph were those micropipe defects.

The central part of the image [Fig. 3(a)] shows no contrast due to other polytype inclusions disoriented relative to the main lattice.

## B. Phase-sensitive SR radiography

Direct evidence for the distribution of tubular-shaped defects in crystal interiors was obtained by SR radiography. By this technique the wafers sliced at different distances from the seed parallel to  $(11\bar{2}0)$  and  $(10\bar{1}0)$  crystal planes were examined. Figure 5 shows the map of defects in the  $(11\bar{2}0)$ -oriented wafer cut, i.e., parallel to the micropipes and the direction normal to the basal plane traces of the far-reaching 2D defective boundaries, from

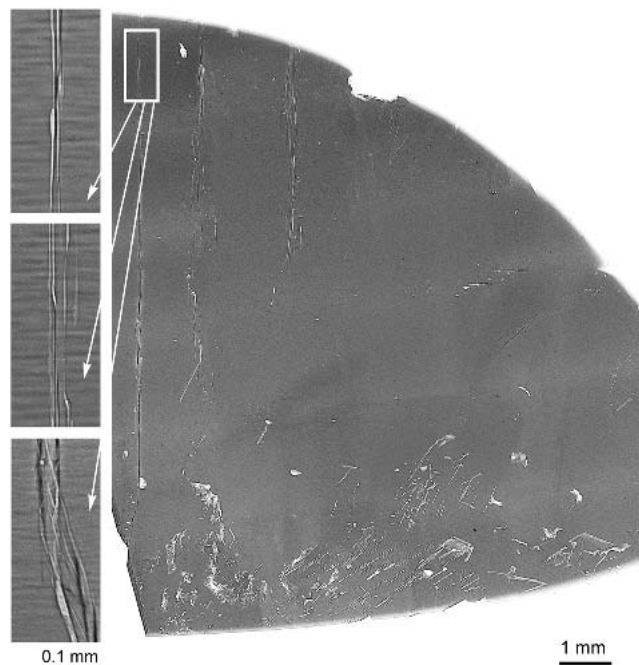


FIG. 5. SR microradiographs of  $(11\bar{2}0)$ -oriented SiC wafer. The inset shows a high-magnification image of bundled micropipes.

a central part of the boule. The most striking features are the defects vertically aligned in the central part of the sample (it should be noted that the radiograph presents only half of the wafer whose central region corresponds to the left-hand section of the image). As demonstrated in the left three figures that are the representative images magnified from the white square on top, this feature shown in the phase-sensitive radiograph resulted from closely spaced or interweaving tubular-shaped defects of approximately 10  $\mu\text{m}$  in diameter. We call such bundles of micropipes here “superpipes.” Increase or decrease in the image intensity of this feature was attributed to the depth variation of their sites under the wafer surface. The existence of the superpipes suggested that micropipes bundled into groups and interacted<sup>11</sup> near the center of the boule. Such micropipe bundling effect will be further discussed in Sec. IV.

The lower part of the image (Fig. 5) provides the view of macrodefects distribution over the close-to-seed area.

#### IV. DISCUSSION

Diffraction topographs and phase-sensitive radiographs of the defects in (0001)-, (10 $\bar{1}$ 0)-, and (11 $\bar{2}$ 0)-oriented slices provide complementary information on the structure and possible nature of these defects. If the 2D defective boundaries that were parallel to the [11 $\bar{2}$ 0] direction and extended rather far from the seed [Fig. 3(a)] were perpendicular to the basal plane, the projection of the defective boundaries onto the (11 $\bar{2}$ 0) plane would have a coincident edge at the hexagonal [0001] direction. That was the direction in which micropipes gathered in groups.

Due to the known distortion associated with the defective boundaries, one would expect that the driving force for the collection of micropipes is the accommodation of the lattice shear. Such transformation of micropipes to superpipes could therefore occur via coalescence, which, under these conditions, could be energetically favorable.

A micropipe can be considered as not only a super-screw dislocation having a Burgers vector but also as a tube having a free surface.<sup>11</sup> When two pipes interact, they always attract each other, on the one hand, due to micropipe free surfaces, regardless of Burgers vector signs. This force varies with distance as approximately  $1/r^3$ .<sup>11</sup> On the other hand, the dislocation components are repelling or attractive depending on their Burgers vector signs, repelling (attractive) for the same (opposite) signs. The force varies with distance as approximately  $1/r$ .<sup>12</sup> In dislocation theory, the first (approximately  $1/r^3$ ) type is called a short-range interaction, while the second (approximately  $1/r$ ) type is a long-range interaction.<sup>12</sup> Micropipes with Burgers vectors of opposite signs attract each other due to both types of interaction, resulting in

their coalescing. Meanwhile those with Burgers vectors of the same signs would interact in a more complicated manner. Nevertheless, one can still expect some equilibrium distance exists where the repelling force due to the long-range interaction is equal to the attractive force due to the short-range one. This distance is rather small and of the order of an average radius of interacting micropipes. For smaller distance, the short-range attracting force prevails and the micropipes coalesce. It is also possible that the long-range repelling force is suppressed by an “external” (with respect to this micropipe pair) shear stress due to residual thermal strains or other defects (micropipes, inclusions, internal boundaries, etc.). Gathering of a large amount of micropipes near the center of the boule and the 2D boundaries, which was attributed to higher defect density, resulted in the wafer regions between the group almost free from micropipes, as illustrated in Figs. 3 and 5. Meanwhile, in the wafers sliced far from the seed, micropipes were scarce. Figure 6 shows the example of the microradiograph demonstrating the defects in a (10 $\bar{1}$ 0)-oriented sample cut far from the central part of the boule. In this area micropipes were practically absent. The macrodefects seen on top were due to heavily damaged crystals containing fractions of material that were recrystallized by a spontaneous growth in the cell between the holder and the back side of the seed—the process scrupulously described in Ref. 8.

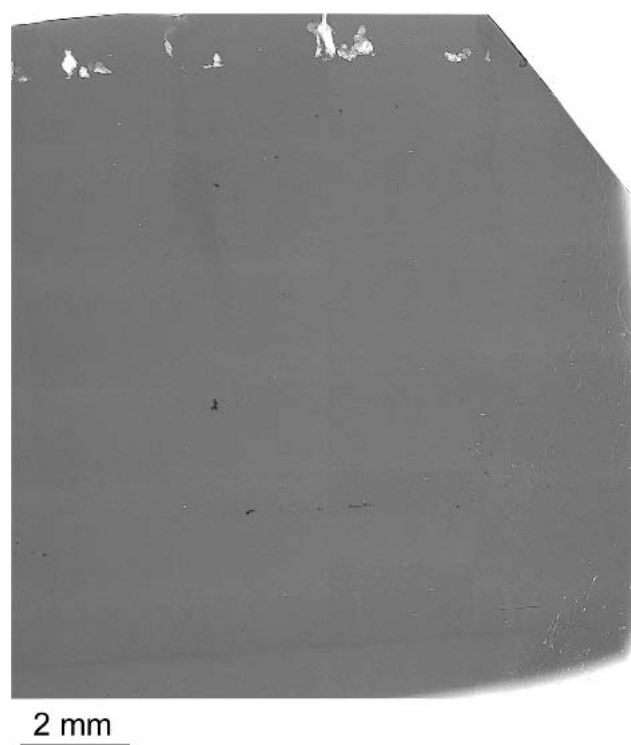


FIG. 6. SR microradiograph of (10 $\bar{1}$ 0)-oriented SiC wafer.

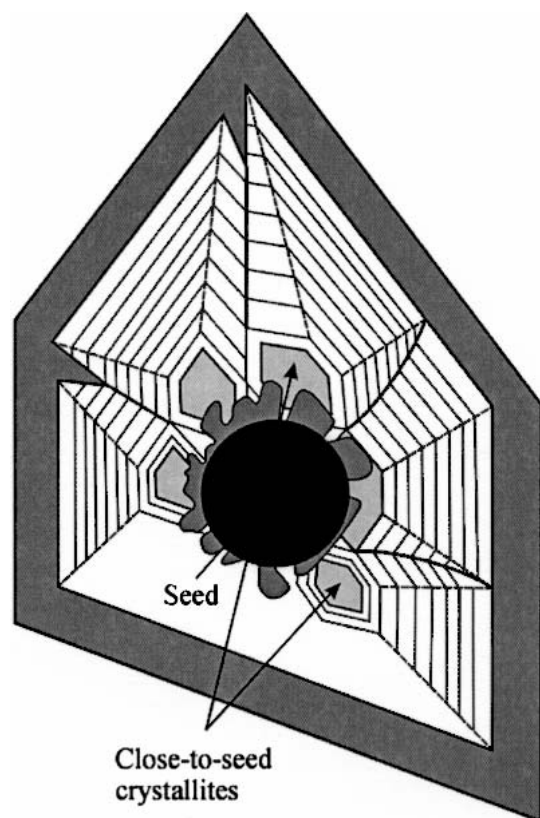


FIG. 7. Scheme illustrating a possible formation of internal boundaries (solid lines) between crystallites resulting from a multiple nucleation on the seed.

## V. CONCLUSIONS

Our x-ray imaging investigation by the combination of topography and phase-sensitive radiography techniques revealed that the distribution of structural defects in the SiC crystals studied here had the following special features: the 2D boundaries parallel to the directions  $\langle 11\bar{2}0 \rangle$  were associated with strong lattice distortions inside crystal interior and served as traps for micropipes. A fine structure of such regions still remains unclear. Not knowing defect arrangements, one can hardly hope to find mechanisms for their formation. Thus, in this section we present some general considerations calling for further investigation.

Specifically, we found the existence of the planar defective boundaries collecting the micropipes. We were also able to map their distribution and orientation. Armed with this information, it would be much easier to propose a suitable model for the growth and therefore understand the origin of the defects.

One could suppose that the growth of a hexagonal SiC crystal on a cylindrical seed started from multiple nucleation. With time, the nuclei appeared to be enclosed by a set of crystallites. The crystallites would finally come in contact and merge together leaving boundaries

inside the crystal body. A scheme illustrating the possible process of the formation of internal boundaries between close-to-seed crystallites is shown in Fig. 7. Thick solid lines trace the defective boundaries emerging at the joints between individuals. Interfacial regions between the merging crystallites were certainly imperfect, and crystalline defects were then generated to relax lattice strain. Specifically, the shear component could be accommodated by superscrew dislocations, or micropipes. The underlying mechanism of the bundling effect of the micropipes and their interaction with the planar defective boundaries however remains uncertain.

The investigation of the region close to the outer surface of the boule ( $l$  in Fig. 2) by diffraction topography in reflection geometry did not reveal the defective boundaries. It is likely that the joints were buried inside the crystal.

The importance of these problems is not confined by the scope of the present study. Sublimation-grown SiC crystals can contain internal boundaries of different types: between other polytype inclusions; between a recrystallized segment and a matrix; between a matrix and a void. The explanation of the evolution of micropipes in the lattice disturbed by these defects would be an important issue.

In summary, the combination of diffraction topography and phase-sensitive radiography has revealed that the micropipe defects are concentrated near internal boundaries buried inside the crystal interior, diminishing their average density in SiC wafers. This information, difficult to obtain with other techniques, is likely to provide a useful clue to the defect formation and the growth of SiC large crystals.

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