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Examining aerosol indirect effect under contrasting environments during the ACE-2 experiment

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Abstract

The Active Tracer High-resolution Atmospheric Model (ATHAM) has been adopted to examine the aerosol indirect effect in contrasting clean and polluted cloudy boundary layers during the Second Aerosol Characterization Experiment (ACE-2). Model results are in good agreement with available in-situ observations, which provides confidence in the results of ATHAM.

Sensitivity tests have been conducted to examine the response of the cloud fraction (CF), cloud liquid water path (LWP), and cloud optical depth (COD) to changes in aerosols in the clean and polluted cases. It is shown for two cases that CF and LWP would decrease or remain nearly constant with an increase in aerosols, a result which shows that the second aerosol indirect effect is positive or negligibly small in these cases. Further investigation indicates that the background meteorological conditions play a critical role in the response of CF and LWP to aerosols. When large-scale subsidence is weak as in the clean case, the dry overlying air above the cloud is more efficiently entrained into the cloud, and in so doing, removes cloud water more efficiently, and results in lower CF and LWP when aerosol burden increases. However, when the large-scale subsidence is strong as in the polluted case, the growth of the cloud top is suppressed and the entrainment drying makes no significant difference when aerosol burden increases. Therefore, the CF and LWP remain nearly constant.

In both the clean and polluted cases, the COD tends to increase with aerosols, and the total aerosol indirect effect (AIE) is negative even when the CF and LWP decrease with an increase in aerosols. Therefore, the first AIE dominates the response of the cloud to aerosols.

1 Introduction

Clouds and their related physical processes represent one of the most complicated and perplexing problems in the study of climate change and weather forecasting

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(Stephens, 2005). One important cloud-related process is the interaction between aerosols, clouds, and radiation which determines the so-called aerosol indirect effect (AIE). Recent estimates of the AIE from global climate models (GCMs) were summarized by Lohmann and Feichter (2005). The first AIE (i.e., the effect of aerosols on the initial cloud drop size distribution) varies from -0.5 to -1.9 W/m^2 , and the second AIE (i.e., the effect of aerosols on cloud precipitation efficiency, cloud lifetime, and cloud morphology) ranges from -0.3 to -1.4 W/m^2 at the top-of-the-atmosphere (Note: the negative sign indicates a cooling effect). If taken at face values, the largest of these estimates would imply that the magnitude of the AIE is larger than estimates of the total forcing from greenhouse gases plus natural causes. However, a total negative (natural plus greenhouse gases plus aerosols) forcing would not be consistent with the observed changes of temperature (Knutti et al., 2002; Anderson et al., 2003). So this might indicate these estimates are too large.

Growing evidence for the difficulty in simulating aerosol effects in GCMs also stems from the application of large eddy simulation (LES) models and/or cloud resolving models (CRMs) to particular cloud systems (Jiang et al., 2002; Ackerman et al., 2004; Lu and Seinfeld, 2005; Jiang and Feingold, 2006) and to comparisons of the modeled response of CRMs and the modeled response of single column models (SCMs) to changes in aerosols (Johnson, 2005; Ovtchinnikov and Ghan, 2005). Direct comparison of CRMs and SCMs also shows that the latter often lack the refined resolution and physical interactions required to properly simulate the aerosol-cloud-radiation interactions (Ghan et al., 2000; Xu et al., 2002; Xie et al., 2002, 2005). Therefore, in this study, we apply a CRM, ATHAM (Active Tracer High-resolution Atmospheric Model) (Oberhuber et al., 1998; Herzog et al., 1998, 2003; Textor et al., 2003; Guo et al., 2006¹), to explore the response of clouds to aerosols, i.e., the AIE.

It is well known that the attenuation of solar radiation by clouds is highly dependent

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on the cloud liquid water path (LWP) (Hansen and Travis, 1974; Stephens, 1978). Han et al. (1998) illustrated that changes in the cloud optical depth (COD) from changes in the cloud LWP could even be larger than aerosol-induced changes of cloud droplet number concentration (N_d) for optically thin clouds ($COD < 15$). The LWP in clouds is highly dependent on the meteorological conditions that lead to the formation of the cloud, because the cloud liquid water content (LWC) is affected by various processes including cloud dynamics (updraft velocity, turbulence) and thermodynamics (Rogers and Yau, 1989). Jiang et al. (2002) pointed out that a weaker subsidence could increase the cloud LWP. Ackerman et al. (2004) emphasized that the response of the cloud LWP to changes in aerosols was determined by the competition between moistening from the suppression of precipitation and drying from the entrainment of overlying air. Only if the air above cloud layer was moist enough or the surface precipitation was significant enough, did the cloud LWP increase with aerosol loading. Otherwise, the cloud LWP could decrease with aerosol loading. Lu and Seinfeld (2005) showed that the cloud LWP was highly sensitive to the large-scale subsidence and sea surface temperature (SST). Therefore, the relationship between cloud LWP and aerosols is complicated by the conditions that influence the cloud LWP. Under different meteorological conditions, cloud LWP may either increase or decrease with increases in aerosols, and thereby either magnify or reduce the total AIE.

In addition to cloud LWP, cloud fraction is another critical factor determining the cloud radiative forcing. Chen and Penner (2005) pointed out that cloud fraction (CF) is the second largest uncertainty after aerosol burden for estimating aerosol forcing. However, the representation of CF is still elusive in weather and climate models; and changes in CF caused by aerosols have not been carefully evaluated. Recently measurement-based assessments showed that CF increased with non-absorbing aerosols (Kaufman et al., 2005; Kaufman and Koren, 2006). However, Lohmann et al. (2006) performed ECHAM4 GCM simulations, used a similar statistical analysis as that in Kaufman et al. (2005), and showed that CF and aerosol concentrations were positively correlated even without any aerosol indirect effect. Thus, in these simula-

tions, the increase in CF was determined by the dynamics rather than the AIE. Penner et al. (2006) also studied the response in CF to changes in aerosols in 3 GCMs and showed that the computed response was not statistically significant.

Large-eddy simulation results have demonstrated that CF can increase with increasing aerosol burden (Ackerman et al., 2003). However, Xue and Feingold (2006) presented opposite results. They ascribed this decrease of CF with increasing aerosol burden to the ease of evaporation of the smaller cloud droplets that form in this case. The two aforementioned studies had different (or even contradictory) conclusions. Therefore, it is important to investigate these effects further. The meteorological conditions, in addition to aerosols, also determine the radiative properties of clouds, and even dominate them. In order to “quantify” the assessment of the AIE, we should and must link the AIE with the meteorological conditions.

To advance our knowledge of the AIE, the first step is to understand the aerosol effect on a regional scale. The second Aerosol Characterization Experiment (ACE-2), was promoted by the International Global Atmospheric Chemistry project (Raes et al., 2000); and was carried out from 16 June to 24 July, 1997 between Portugal, the Azores, and the Canary Islands (29.4 N, 16.7 W). The CLOUDY-COLUMN (CC) experiment, one of the six field projects during ACE-2, aimed to examine aerosol effects on the microphysical and radiative properties of marine boundary layer clouds (Brennguier et al., 2000a, 2003). Previous observations showed that the clean marine air over this area often alternates with the anthropogenic pollution originating from Europe and North Africa, and this alternation was also observed during ACE-2 (Verver et al., 2000). Conditions during this alternation allowed us to examine aerosols, cloud properties, and their related processes in contrasting clean and polluted environments. This CC experiment also provided extensive, concomitant, and independent measurements of aerosols and clouds, which provided a good opportunity to explore the AIE for both experimentalists and modelers (Pawlowska and Brenguier, 2000, 2003; Menon et al., 2003).

The purpose of this study is to explore the impact of aerosols and meteorologi-

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cal conditions on the response of the cloud LWP, CF, and cloud COD to changes in aerosols. The following questions will be addressed: (1) how do cloud LWP and CF respond to changes in aerosols? (2) how do large-scale meteorological conditions influence the response of cloud LWP and CF to aerosols, and how large is this influence?

The paper is organized as follows: Sect. 2 describes the clean and polluted base cases, and the simulation set-up. Section 3 discusses the model results from these base cases, and compares them with available observations. Section 4 presents sensitivity tests of the effects of changes in aerosols and the meteorological setting on the COD and radiative forcing. Section 5 summarizes our results.

2 Case descriptions and simulation set-up

2.1 Case descriptions

In this study, we focus on two contrasting marine stratocumulus cloud cases during the CC experiment. The first case is a pristine case on 26 June 1997, while the second is a polluted case on 9 July 1997. On 26 June the non sea-salt (nss) sulfate mass concentration measured at the surface site at Punta del Hidalgo was $0.3 \mu\text{g}/\text{m}^3$; whereas on 9 July, it was $2.8 \mu\text{g}/\text{m}^3$. The airborne condensation nuclei (CN) counter measurement of the total aerosol number concentration was 218 cm^{-3} and 636 cm^{-3} on 26 June and on 9 July, respectively (Snider et al., 2000). In this study, aerosol mass mixing ratio was assumed to be constant with altitude below cloud base because the boundary layer was well mixed (see Fig. 1) (Delene and Deshler, 2001; Penner et al., 2004). These measured sulfate concentrations and total aerosol number concentration were used in the parameterization of droplet nucleation which was based on Chuang and Penner (1995). For our implementation of this parameterization, we assumed that 75% of the mass fraction of sulfate (nss-sulfate) was produced by the aqueous phase pathway, and resided on pre-existing accumulation mode particles, which followed a specified three-mode size distribution (Chuang and Penner 1995; Chuang et al., 1997). Here-

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after, we denoted the clean base case (on 26 June) as “CACM” for “clean aerosol clean meteorology” and the polluted base case (on 9 July) as “PAPM” for “polluted aerosol polluted meteorology” (Table 1).

The alternation between clean and polluted cases (over the ACE-2 area) was governed by the location of the Azores High. Between 25 June and 3 July, Western Europe was under the influence of a cyclone, and this cyclone brought pristine marine polar air into the ACE-2 area. However, between 4 and 10 July, Western Europe was under the influence of the extension of the Azores High; and this high pressure brought about an outflow of anthropogenic pollution from Europe into the ACE-2 area (Verver et al., 2000). Between the “CACM” and the “PAPM”, both the aerosol burden and the background meteorological conditions changed as a result of these two events.

2.2 Simulation set-up

In this study, the three-dimensional version of ATHAM was applied (Oberhuber et al., 1998; Herzog et al., 1998, 2003; Textor et al., 2003; Guo et al., 2006¹). The horizontal domain used 64×64 grid points with a uniform spacing of 100 m. The vertical spacing was uniformly 30 m below 2 km and then stretched to 300 m near the model top. We used a time step of 2 s, and model results were archived every 5 min.

ATHAM was initialized with the reanalysis data from European Centre for Medium-Range Weather Forecasts (ECMWF) (Fig. 1). In the “CACM” case, initially the boundary layer was capped by a strong inversion layer. Within the boundary layer, both the potential temperature and the specific humidity were almost constant with height (about 292 K and 10 g/kg). Above the inversion layer, the air was relatively dry (relative humidity ~25%). In the “PAPM” case, initially the boundary layer exhibited similar features to that in the “CACM”, but it was shallower, warmer, and moister. This shallow boundary layer was due to the large-scale subsidence that is associated with the influence of the Azores High during the “PAPM” period. We assumed cloud free conditions at the beginning of all simulations.

The time- and height-dependent horizontal large-scale advection of temperature (T)

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and moisture (Q) were prescribed from the ECMWF reanalysis data. The corresponding vertical large-scale advection was calculated from the vertical gradients of the simulated T and Q profiles and the prescribed large-scale vertical velocity (W) as shown in Fig. 2. It is evident that there is a weak large-scale subsidence in the “CACM” case but a very strong large-scale subsidence (up to -1.2 cm/s) in the “PAPM” case between 1 and 2 km, which is consistent with the dominance of the Azores High over the ACE2 area in the “PAPM” case. The modeled horizontally averaged wind components (U , V) from ATHAM were nudged towards the ECMWF re-analysis data with a relaxation time scale of 1 h (Grabowski et al., 1996). The ECMWF reanalysis data were available every six hours, so they were linearly interpolated to the temporal (and vertical) resolution of the ATHAM grids between these times.

Our model simulations began at 18:00 LST (local standard time) and the entire simulation period was 30 h. The first 6-h simulation was regarded as a spinup and our analysis was only performed for the last 24-h.

3 Base case results

In the base case study, we compared the model results with available observations in order to examine the extent to which our model is able to reproduce observations. The most extensive observations were conducted close to local noon (Brenquier et al., 2000a, 2000b; Snider et al., 2000; Snider et al., 2003). The flight time, on 26 June was from 11:50 LST to 14:50 LST, while on 9 July, it was from 12:40 LST to 15:50 LST, and, our comparisons were conducted at the corresponding flight time.

3.1 Cloud vertical profiles

Figures 3 and 4 show the simulated and observed vertical profiles of cloud liquid water content (LWC), cloud droplet number concentration (N_d), and cloud droplet mean volume diameter (d_v) within cloud layers for the “CACM” and “PAPM” cases, respectively.

The LWC increases with the height above cloud base; N_d remains almost constant within cloud layers; and d_v increases with height above cloud base as expected if N_d is constant.

The comparison of LWC with the adiabatic LWC is illustrated in Figs. 3–4a and d. The difference between the modeled LWC and the adiabatic LWC becomes larger with the height above cloud base due to the mixing of dry air from above the cloud, especially in the case of “PAPM”.

The simulated N_d is slightly higher than the observed N_d near cloud base (Figs. 3–4b and e), which is partly due to the fact that the lower size limit of the Fast-FSSP (Forward Scattering Spectrometer Probe) measurement is $1.3\ \mu\text{m}$ in radius (Brenguier et al., 2003). N_d decreases at cloud top, due to the entrainment of dry air from above the cloud. We also note that the difference between d_v and the adiabatic d_v (solid lines in Figs. 3–4c and f) is small from cloud base to cloud top in both the ATHAM results and airborne measurements. This might suggest that the mixing between cloudy air and dry air above clouds tends to be heterogeneous because N_d becomes smaller but d_v remains close to the adiabatic d_v at cloud top. However, this convergence of d_v towards the adiabatic d_v is less evident in ATHAM than in the measurements. The standard deviation of the simulated d_v is about $2\sim 5\ \mu\text{m}$, which is larger than that of the observed d_v ($2\sim 3\ \mu\text{m}$) (Pawlowska et al., 2006). This might be due to the lack of representations of heterogeneous mixing in ATHAM, a problem that many CRMs or large eddy simulation models share (Herzog et al., 2003; Grabowski et al., 2006).

3.2 Vertical velocity and cloud droplet number concentration spectra

As noted above, the mechanistic aerosol activation scheme of Chuang and Penner (1995) was adopted, in which the initial N_d predicted at cloud base depends on the updraft velocity w . Figure 5 compares probability density functions (PDFs) of w near cloud base from ATHAM and from the airborne measurements for the “CACM” case and for the “PAPM” case. The solid lines represent modeled PDFs of w and the dashed lines represent the measurements around local noon. The modeled w captures the

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peaks and the spread of the frequency distribution of the observed w very well. The simulated mean values (\overline{w}) and the standard deviations (σ_w) agree reasonably well with the observed values (difference <30%) for both the “CACM” and the “PAPM” cases. \overline{w} is less than the measurement bias of 0.3 m/s (Brown, 1993; Guibert et al., 2003), but σ_w is relatively large, especially in the “PAPM” case. Therefore, the average of the vertical velocity is close to zero within the boundary layer; and local fluctuations due to boundary layer turbulence are a major contributor to the updraft at cloud base. This turbulence, therefore, leads to the activation of aerosols, as well as the vertical transport of moisture, mass, and energy.

N_d is calculated from a budget equation (Lohmann et al., 1999), which takes advective transport, turbulent transport, and microphysical sinks and sources into account as follows:

$$\frac{\partial N_d}{\partial t} = R(N_d) + Q_{\text{nucl}} - Q_{\text{auto}}(N_d) - Q_{\text{self}} - \frac{N_d}{\text{LWC}}(Q_{\text{accr}} + Q_{\text{evap}}), \quad (1)$$

where $R(N_d)$ refers to the advective and turbulent transport of N_d , and Q_{nucl} , $Q_{\text{auto}}(N_d)$, and Q_{self} refer to the rates for nucleation, autoconversion, and self-collection of cloud droplets, respectively. Q_{accr} and Q_{evap} refer to the rates of accretion and evaporation, respectively.

Figure 6 presents cloud droplet number concentration spectra at five evenly-divided cloud layers with $N_d > 20 \text{ cm}^{-3}$ from ATHAM and from the airborne measurements for the “CACM” case. The sample section criterion of $N_d > 20 \text{ cm}^{-3}$ is to achieve a statistically significant estimate of d_v and LWC derived from the Fast-FSSP measurements (Brenquier et al., 2003). At cloud base, the simulated droplet number concentration spectra agree well with the observed spectra. This is expected because the updraft velocity w compares favorably with the observed w and the aerosol chemical composition and number concentrations used in ATHAM are from observations. However, near cloud top, the agreement between the simulated and the observed spectra is not as good as that at cloud base, although the mean value ($\overline{N_d}$) and the associated standard deviation (σ_{N_d}) are still close to the observed values (difference <10%). These differ-

ences might be caused by the lack of representation of heterogeneous sub-grid scale mixing as mentioned in Sect. 3.1. For the polluted “PAPM” case, the results of the comparison between modeled and observed number concentration spectra are similar to those for the “CACM” case, and therefore, are not shown here.

5 3.3 Other cloud properties

Figures 7 and 8 show the time series of CF, in-cloud LWP, domain average LWP, N_d , d_v , and COD. It is clear that the model results are generally well within the uncertainties of the observations. For example, the averaged in-cloud LWPs from ATHAM are $\sim 19.6 \text{ g/m}^2$ and $\sim 13.5 \text{ g/m}^2$ for the “CACM” and for the “PAPM” cases, respectively, which are well within the observations ($18.5 \pm 17.8 \text{ g/m}^2$ and $11.0 \pm 10.8 \text{ g/m}^2$, respectively) (Menon et al., 2003). We also note that the in-cloud LWP for “PAPM” case is smaller than that for the “CACM” case, which would not be expected if caused by the second AIE. The reason is that the polluted cloud system in the case of “PAPM” tends to be shallower and drier (as shown in Figs. 3–4) largely due to its continental origin (Verver et al., 2000; Pawlowska and Brenguier, 2003).

CF is defined as the fraction of model cloudy columns, and a cloudy column is a column containing cloudy grid cells with liquid water mixing ratio $> 0.01 \text{ g/kg}$ and $N_d > 5 \text{ cm}^{-3}$ (Pawlowska and Brenguier, 2003). The observed CFs are from a Compact Airborne Spectrographic Imager (CASI) (indicated as “*” in Figs. 7–8a) and from the International Satellite Cloud Climatology Project (ISCCP) (indicated as “o”) (Schroder et al., 2002; Menon et al., 2003). In both the “CACM” and the “PACM” cases, the simulated CFs are overestimated, especially at night and in the early morning. This discrepancy might be due to three reasons: (1) the sample area: the ATHAM domain ($6.4 \times 6.4 \text{ km}^2$) is much smaller than that of ISCCP (250 km); (2) the cloud scene identification criterion in the retrieval of CF is too strict: for example, in both the “CACM” and the “PAPM” cases, CFs are $\sim 90\%$ from the airborne sampling, but they are $< 80\%$ from the CASI retrieval (Brenguier et al., 2003); (3) the existence of shallow and broken clouds, which are difficult to retrieve (Menon et al., 2003).

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4 Sensitivity tests

Here we consider sensitivity studies to examine how the CF and cloud LWP respond to changes in the aerosol burden, and how the meteorological conditions might amplify/reduce this response. Two sensitivity tests were performed to isolate the effect of aerosols and the effect of the meteorological conditions on cloud properties (Table 1). For these tests, the aerosol burdens between the “CACM” and the “PAPM” cases were swapped. We denoted the test which combined the aerosol loading in the “PAPM” case with the meteorological setting in the “CACM” case as the “PACM”. The case where the aerosol burden from the “CACM” case was used with the meteorological setting in the “PAPM” case was denoted “CAPM”. As in the base cases, these two sensitivity tests were run for 36 h, and we analyzed the results of the last 24 h.

4.1 Response of CF and LWP to changes in aerosol burden

4.1.1 Clean meteorology case

As discussed in the introduction, the CF and LWP are important for determining the cloud radiative forcing. In both the “CACM” and the “PACM” cases, CFs and LWPs have a significant diurnal variation: CFs reach their maxima ($\sim 100\%$ and $\sim 95\%$) during the local early morning and their minima ($\sim 75\%$ and $\sim 45\%$) in the local afternoon; the in-cloud LWPs reach 65.9 g/m^2 and 48.7 g/m^2 in the morning, and 14.8 g/m^2 and 10.5 g/m^2 in the afternoon. However, both the CF and LWP are generally larger in the “CACM” case than in the “PACM” case (Fig. 7). At local noon when the solar radiation is strong, the CF in the “CACM” case is about twice of that in the “PACM” case; the LWP in the “CACM” case is larger by 50%. Ackerman et al. (2004) pointed out that the cloud water content is affected not only by precipitation, but also by the entrainment of dry air from above clouds. Since surface precipitation is small ($< 0.05 \text{ mm/day}$) in both the “CACM” and in the “PACM” cases, the response of cloud water to aerosols is primarily determined by the entrainment of dry air at cloud top.

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Figure 9a presents the time series of the averaged cloud top height. It is clear that the cloud top growth rate is smaller in the “CACM” case (0.71 cm/s) than in the “PACM” case (0.75 cm/s). Since the entrainment rate can be approximated by the difference between the cloud top growth rate and the large-scale vertical velocity (Moeng et al., 1999; Jiang et al., 2002; Stevens et al., 2003a, 2003b), and since the large-scale vertical velocity (subsidence) is the same in the two cases, the entrainment rate is smaller in the “CACM” case than that in the “PACM” case. This is consistent with the large eddy simulation results (Jiang et al., 2002; Ackerman et al., 2004). For the clean case, the net latent heat release associated with precipitation formation (which allows increased condensation) offsets the radiative cooling near the cloud top, which lessens the negative buoyancy and vertical motions (particularly in the downdrafts) of air parcels (Fig. 9c), thereby decreasing the kinetic energy (Fig. 9b) and reducing the entrainment (Fig. 9a). Moreover, the air above clouds is relatively dry with a relative humidity (RH) of ~25% (Fig. 10). Therefore, the drying by entrainment in the “CACM” case is not as efficient as that in the “PACM” case, and the clouds are moister and CF and LWP are larger in the “CACM” case.

We also note that the largest difference in the CF between the “CACM” case and the “PACM” case occurs at local noon when the CF reaches a minima. This might be due to the dependence of the entrainment rate on the inverse of the cloud radius (Blyth et al., 1988). The decreased CF in the “PACM” case, or the smaller cloud radius, allows cloudy air to have a larger surface area to mix with the drier ambient air, and thus enhances the entrainment. This in turn reduces the cloud radius (and the cloud moisture) further.

4.1.2 Polluted meteorology case

CFs and LWPs are similar in the “CAPM” and in the “PAPM” cases (Fig. 8). The differences in their daily averages are well within their standard deviations (Table 2). There is negligible surface precipitation and the air above clouds is even drier than that in the “CACM” and “PACM” cases (RH <20%) (Fig. 10), but the cloud top growth is highly

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suppressed by the strong large-scale subsidence on 9 July (Figs. 2 and 11a). This strong subsidence warming dominates the net latent heating associated with precipitation formation in the cloud; and masks the difference between the vigor of the vertical motions, TKE, and entrainment between the “CAPM” and the “PAPM” cases (Figs. 11).

5 As a result, the difference in the cloud top growth rates in the “CAPM” case (0.47 cm/s) and in the “PAPM” case (0.47 cm/s) is insignificant, which means that the difference in the entrainment drying is insignificant, so that this does not give rise to a noticeable difference in cloud water. Consequently, cloud water contents are similar in the “CAPM” and in the “PAPM”, and so are CFs and LWPs.

10 4.2 Response of COD to changes in aerosol burdens

Although the response of the CF and LWP to aerosols is negative or neutral (as discussed in Sect. 4.1), the response of the COD to increases in aerosol burden tends to be positive (Figs 7–8). As expected from the 1st AIE (Twomey et al., 1977), N_d increases and d_v decreases when aerosol burdens increase (Figs. 7–8). Since the broad band COD can be approximated by $\frac{3LWP}{\beta d_v \rho_w}$ where β is ~ 1.08 for maritime clouds (Martin et al., 1994), a decrease of COD with increases in aerosol burden would only occur if the decrease of LWP with aerosols is faster than the decrease of d_v . Although a decrease of d_v with an increase in aerosols is observed (Penner et al., 2004), the LWP can either increase or decrease with aerosols (Lu and Seinfeld, 2005). In this study, the decrease of LWP is not as significant as the decrease of d_v (Table 2). On the whole, the COD increases with increasing aerosol burden.

4.3 Radiative forcing estimates

Here, we evaluate radiative effects of changes in aerosol burden in the clean and polluted cases. ΔF_t is the difference in the net flux (shortwave plus longwave) either at the top-of-the-atmosphere (TOA) or at the surface between simulations with clean and

polluted aerosol burden but with the same meteorological conditions, and is given by

$$\Delta F_t(\text{"CM"}) = F_{\text{net}}(\text{"PACM"}) - F_{\text{net}}(\text{"CACM"})$$

or

$$\Delta F_t(\text{"PM"}) = F_{\text{net}}(\text{"PAPM"}) - F_{\text{net}}(\text{"CAPM"}).$$

5 Where ΔF_t ("CM") is the total AIE with the meteorological setting of the clean base case on 26 June; and ΔF_t ("PM") is the total AIE with the meteorological setting of the polluted base case on 9 July.

The first and second aerosol indirect effects are often entangled in the real atmosphere, so that there is no simple way to isolate them. Because the first AIE is mainly induced by the increases in N_d due to anthropogenic pollution (Twomey et al., 1977), and N_d is more or less constant within cloud layers (Fig. 3), we replaced N_d in the base cases with the averaged N_d from the sensitivity tests, and estimated the first AIE offline (ΔF_1). Thus, we calculated the ΔF_1 under the clean meteorological conditions for the clean base case on 26 June (ΔF_1 ("CM")) and under the polluted meteorological conditions for the polluted base case on 9 July (ΔF_1 ("PM")).

15 The second indirect effect (ΔF_2) is estimated as the difference between ΔF_t and ΔF_1 ; so ΔF_2 is given by

$$\Delta F_2(\text{"CM"}) = \Delta F_t(\text{"CM"}) - \Delta F_1(\text{"CM"})$$

or

$$20 \Delta F_2(\text{"PM"}) = \Delta F_t(\text{"PM"}) - \Delta F_1(\text{"PM"}).$$

Table 3 summarizes the daily averaged ΔF_t , ΔF_1 , and ΔF_2 using above method. The magnitude of ΔF_t ("CM") is smaller than that of ΔF_1 ("CM"), and results in a positive ΔF_2 ("CM"). This is expected because both CF and LWP decreased with an increase in aerosols in the "CACM" and "PACM" cases. But the positive ΔF_2 ("CM") is dominated by the negative ΔF_1 ("CM"), so that the net effect is negative (cooling).

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The magnitudes of ΔF_t ("PM"), ΔF_1 ("PM"), and ΔF_2 ("PM") are smaller than those of ΔF_t ("CM"), ΔF_1 ("CM"), and ΔF_2 ("CM") mainly because the clouds on 9 July are relatively drier and shallower. The magnitude of ΔF_t ("PM") is similar to that of ΔF_1 ("PM"), and ΔF_2 ("PM") is close to 0, which is consistent with the neutral responses of CF and LWP to increases in aerosols under the strong large-scale subsidence on 9 July.

5 Conclusions and discussion

A cloud resolving model has been used to study aerosol indirect effect during the second Aerosol Characterization Experiment (ACE-2). The model is able to reproduce the cloud characteristics in both the clean and polluted base cases. The simulated cloud vertical profiles, updraft velocity near cloud base, and cloud droplet number concentration spectra are in good agreement with the observations.

The mean value of the updraft velocity near cloud base is generally small, but the standard deviation is large, which indicates that local fluctuations caused by turbulence, rather than the mean updraft, determine the aerosol activation process. Heterogeneous mixing between cloudy air and clear air near cloud top is not well simulated because the underlying assumption for the turbulence mixing scheme in ATHAM is homogeneous (Herzog et al., 2003).

Sensitivity tests were performed by swapping the aerosol burden and the meteorological conditions between the clean and polluted base cases in order to explore the response of the cloud fraction, cloud water path, and cloud optical depth to increases in aerosols. Our results show that neither the cloud fraction nor the cloud liquid water path necessarily increase with increases in aerosols. When there is a weak large-scale subsidence and dry overlying air above clouds, the stronger entrainment of dry air (due to a reduction of the net latent heating implied by the suppression of precipitation; and thus an enhancement of the cooling of air parcels and their negative buoyancy and kinetic energy) in the polluted aerosol scenario effectively dries the cloud, and gives rise to a lower cloud fraction and cloud liquid water path (i.e., a negative response).

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However, when the large-scale subsidence is strong and suppresses the rise of the cloud top, this entrainment drying does not result in a significant difference in cloud water for different aerosol burdens. The cloud fraction and cloud liquid water path are not sensitive to changes in the aerosol burden (i.e., neutral response).

5 The second aerosol indirect effect could either enhance or reduce the first aerosol indirect effect, depending on the large-scale subsidence. So different meteorological conditions could either magnify or diminish the first indirect effect. Because of these complications, evaluation of the second aerosol indirect effect is complex and challenging.

10 Nevertheless, the total aerosol indirect effect still tends to be negative, that is, it exerts a cooling effect. Even when the cloud fraction and liquid water path decrease with increases in aerosols, the cloud optical depth tends to be larger in the polluted aerosol scenario (due to the smaller cloud droplet size). For the clean meteorological case, we found a positive second aerosol indirect effect that decreased the total indirect effect, but did not change the sign of the total indirect effect.

15 *Acknowledgements.* This work was supported by the U.S. Department of Energy Atmospheric Radiation Measurement (ARM) Program under grant #DEFG02-97-ER62370.

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Table 1. Descriptions of two clean and polluted base cases (“CACM” and “PAPM”), and two sensitivity tests (“PACM” and “CAPM”).

Tests	Aerosol loading	Meteorological setting
“CACM”	Clean (26 June)	Clean (26 June)
“PACM”	Polluted (9 July)	Clean (26 June)
“CAPM”	Clean (26 June)	Polluted (9 July)
“PAPM”	Polluted (9 July)	Polluted (9 July)

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Table 2. Daily-mean and standard deviation of cloud fraction (CF), in-cloud liquid water path (LWP), cloud droplet number concentration (N_d), droplet volume mean diameter (d_v), and cloud optical depth (COD) for the four cases.

Cases	CF (%)		In-cloud LWP (g/m ²)		N_d (cm ⁻³)		d_v (μm)		COD	
	mean	Std.	mean	Std.	mean	Std.	mean	Std.	mean	Std.
“CACM”	91.93	6.66	33.89	15.22	41.51	3.95	16.90	2.40	4.95	1.57
“PACM”	79.19	15.07	26.72	13.98	236.18	14.16	9.27	1.00	7.24	3.01
“CAPM”	73.18	7.91	14.62	2.15	47.59	2.82	13.92	0.70	2.62	0.31
“PAPM”	72.80	6.66	14.94	2.76	248.33	24.26	8.85	0.65	4.44	0.59

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Table 3. The daily and spatially averaged total aerosol indirect effect (ΔF_t), first aerosol indirect effect (ΔF_1), and second aerosol indirect effect (ΔF_2) at the top-of-the-atmosphere and at the surface. These averages are derived from instantaneous calculations of every 5 min of model time. “CM” refers to the meteorological conditions for the clean base case on 26 June; and “PM” refers to the meteorological conditions for the polluted base case on 9 July.

	top-of-the-atmosphere (W/m ²)			surface (W/m ²)		
	ΔF_t	ΔF_1	ΔF_2	ΔF_t	ΔF_1	ΔF_2
“CM”	−8.8	−19.5	+10.7	−10.6	−20.9	+10.3
“PM”	−6.4	−6.0	−0.4	−6.8	−6.2	−0.6

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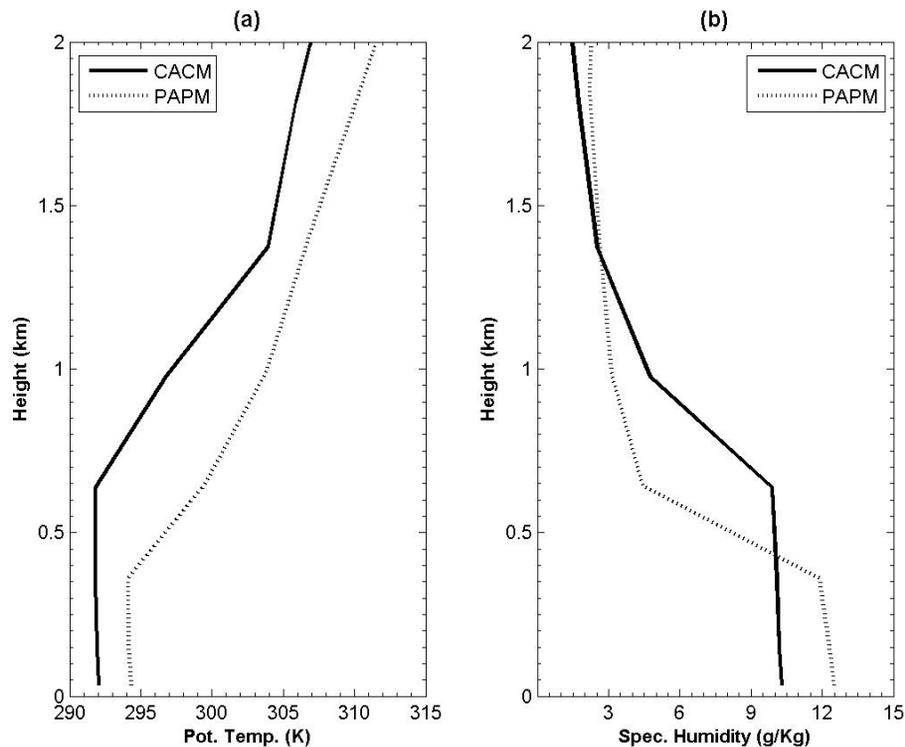


Fig. 1. Vertical profiles of the initial conditions: **(a)** potential temperature and **(b)** specific humidity for the clean base case on 26 June (“CACM”) and for the polluted base case on 9 July (“PAMP”).

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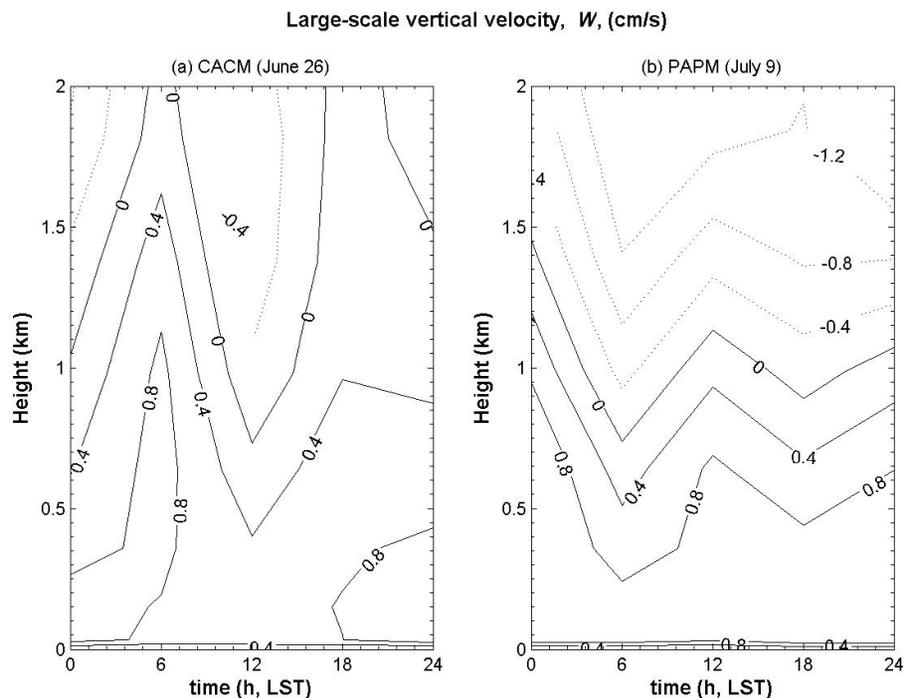


Fig. 2. Large-scale vertical velocity (W) from the ECMWF reanalysis data for the clean base case on 26 June (“CACM”) and for the polluted base case on 9 July (“PAMP”). Contour intervals are 0.4 cm/s. [Note: solid lines are non-negative values, and dotted lines are negative values].

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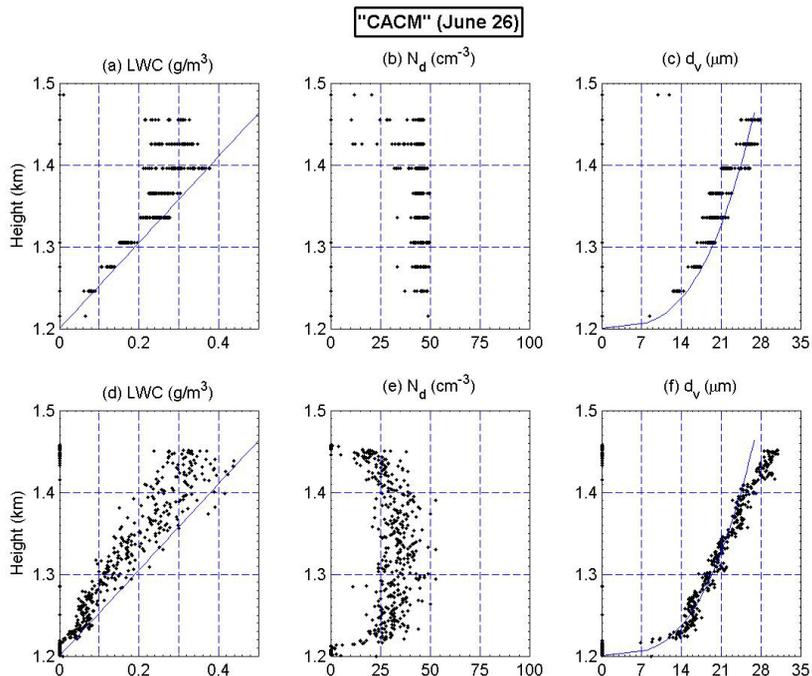


Fig. 3. Vertical profiles of the cloud microphysical properties for the “CACM” case on 26 June from the model (ATHAM) results **(a)–(c)** and from airborne measurements **(d)–(f)** sampled in the cloudy region around local noon: cloud liquid water content (LWC) is shown in (a) and (d), cloud droplet number concentration (N_d) in (b) and (e), and droplet volume mean diameter (d_v) in (c) and (f). The solid lines in (a) and (d) refer to the adiabatic LWC; and the solid lines in (c) and (f) refer to the adiabatic d_v , assuming $N_d=50\text{ cm}^{-3}$ for the “CACM” case.

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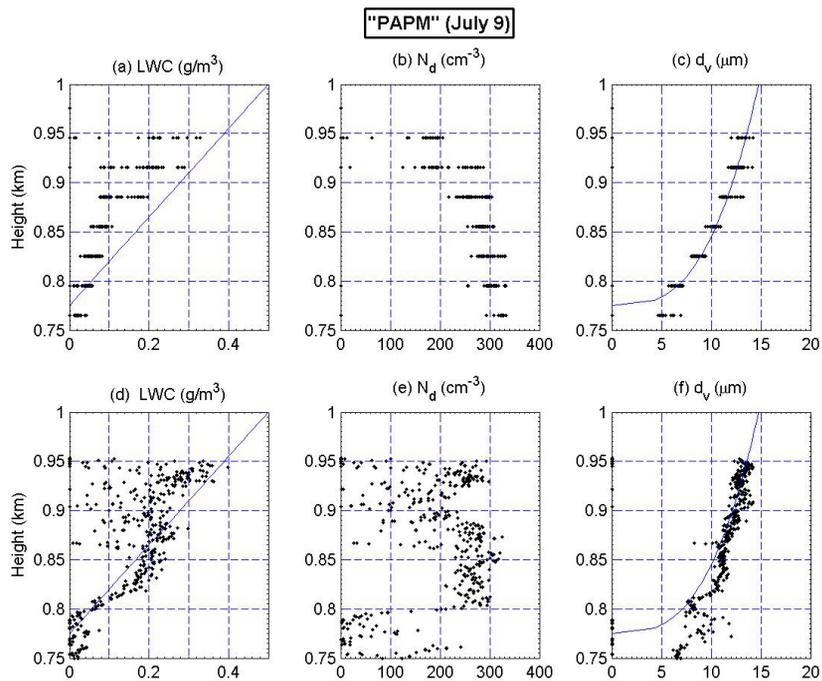


Fig. 4. Same as in Fig. 3 but for the “PAPM” case. The solid lines in (c) and (f) refer to the adiabatic d_v , assuming $N_d = 300 \text{ cm}^{-3}$ for the “PAPM” case.

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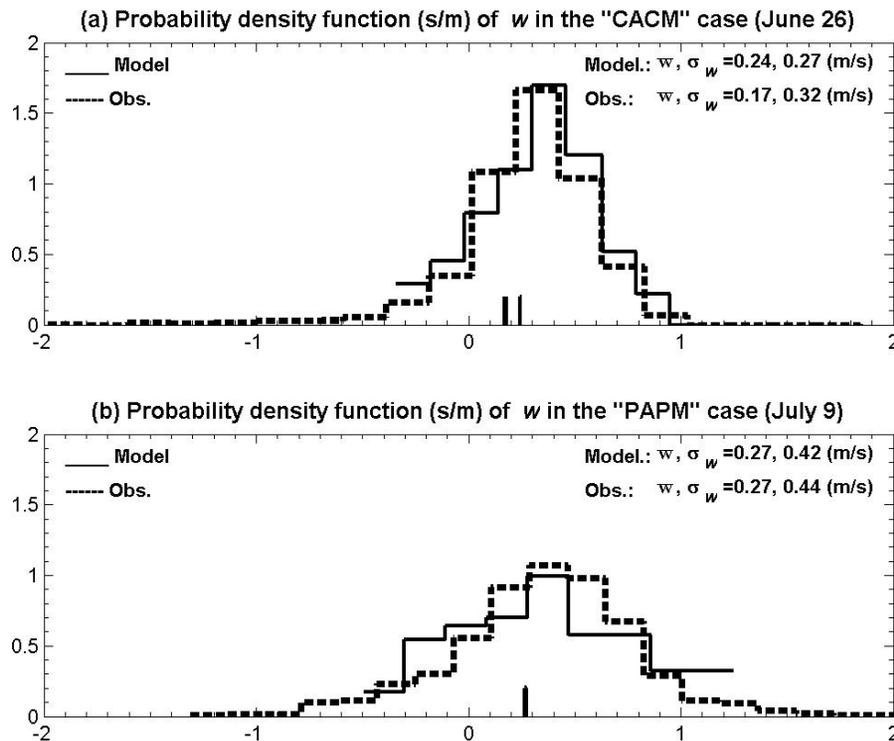


Fig. 5. Probability density function of the vertical velocity (w) near cloud base from the model (ATHAM) results (solid) and from airborne measurements (dashed) in the “CACM” case in **(a)** and in the “PAPM” case in **(b)**. The vertical bars on the x-axis represent the mean values (\bar{w}).

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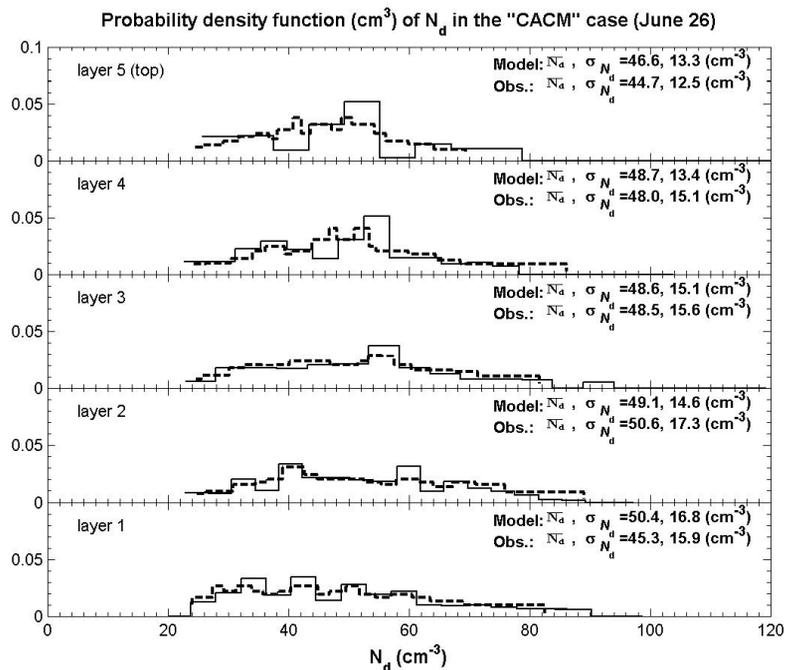


Fig. 6. Probability density function for the droplet number concentration (N_d) at five evenly-divided cloud layers with $N_d > 20 \text{ cm}^{-3}$ from the model (ATHAM) (solid) and from airborne measurements (dashed) in the "CACM" case.

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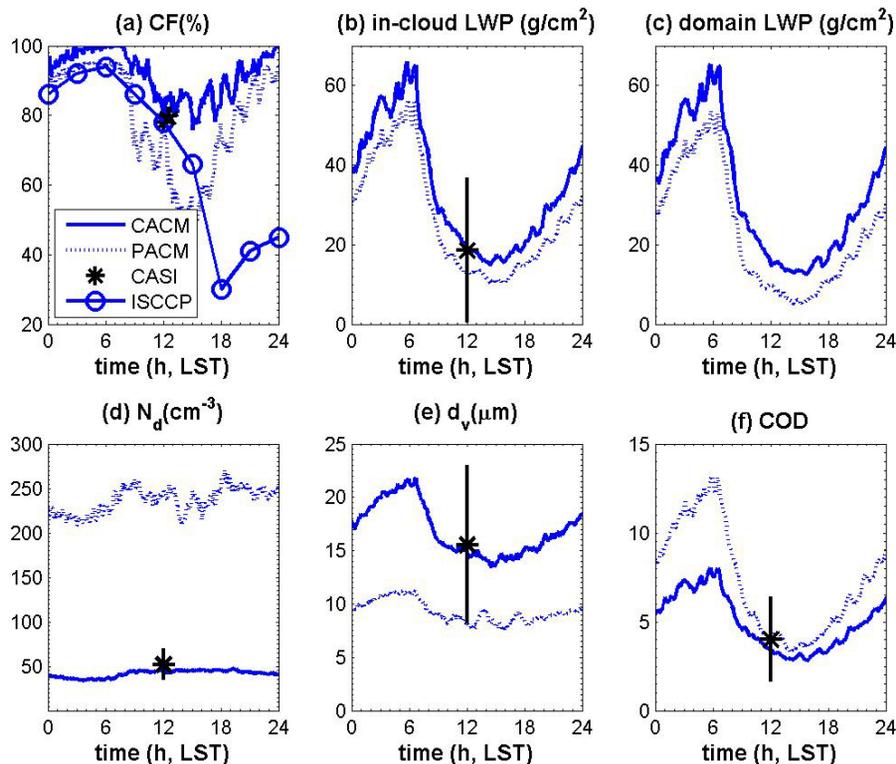


Fig. 7. Time series of cloud fraction **(a)**, in-cloud liquid water path (LWP) **(b)**, domain-average LWP **(c)**, droplet number concentration N_d **(d)**, droplet volume mean diameter d_v **(e)**, and cloud optical depth (COD) **(f)** for the “CACM” case (solid) and for the “PACM” case (dotted). Observed values and their uncertainties are also indicated with “*” and vertical bars.

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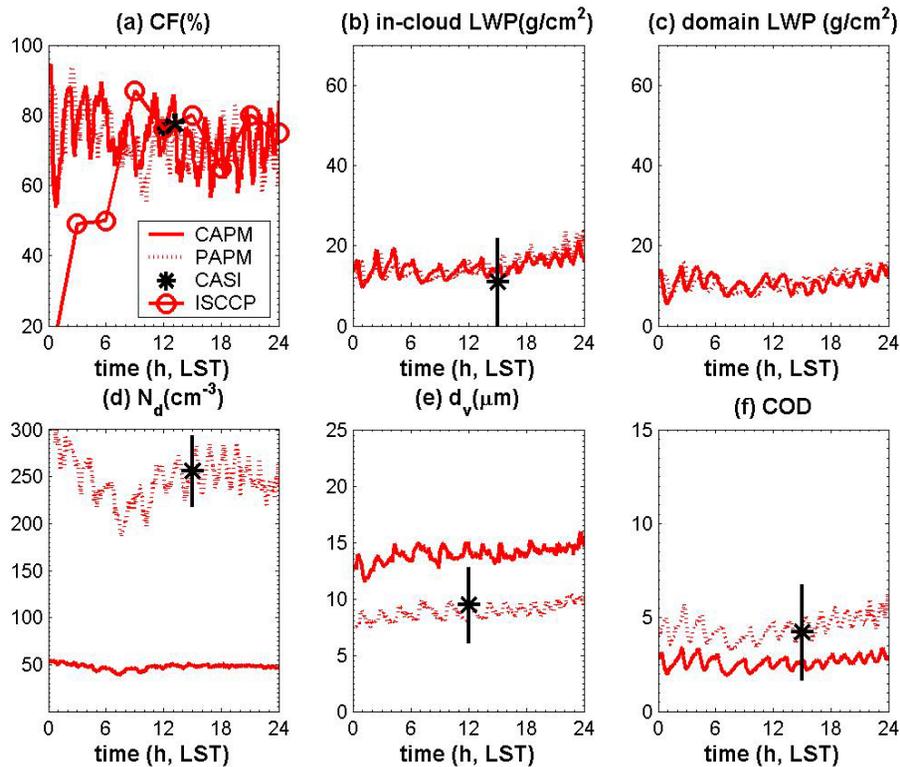


Fig. 8. Same as in Fig. 7 but for the “CAPM” case (solid) and for the “PAPM” case (dotted).

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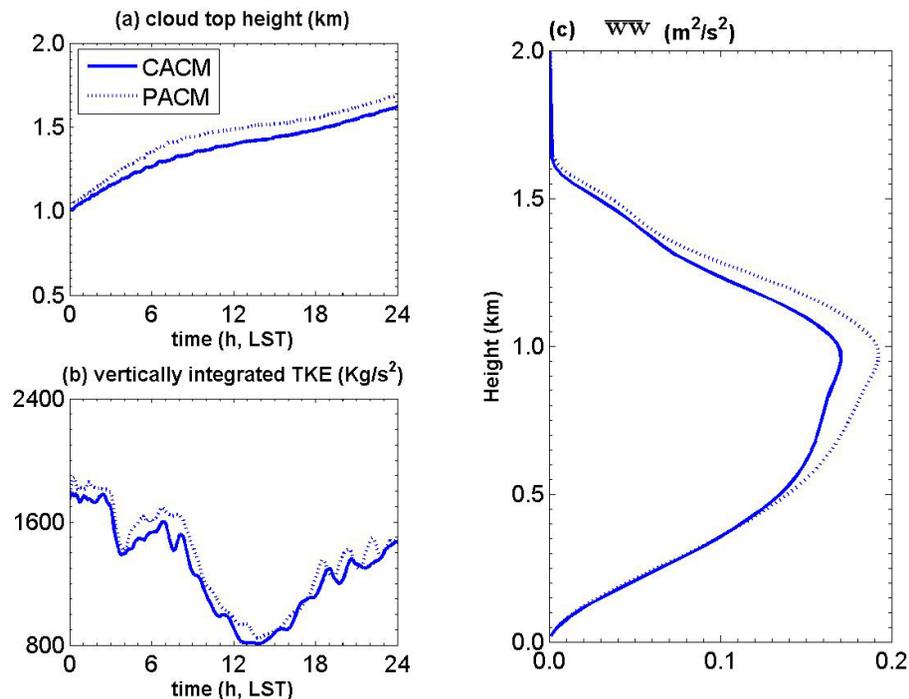


Fig. 9. Time series for the “CACM” case (solid) and for the “PACM” (dotted) case: of the **(a)** horizontally-averaged height of the cloud tops (km), and **(b)** vertically integrated turbulent kinetic energy (TKE) under 2 km (kg/s^2); **(c)** is the time average over the last 24 h of the profiles of w variance $\overline{w w}$ (m^2/s^2).

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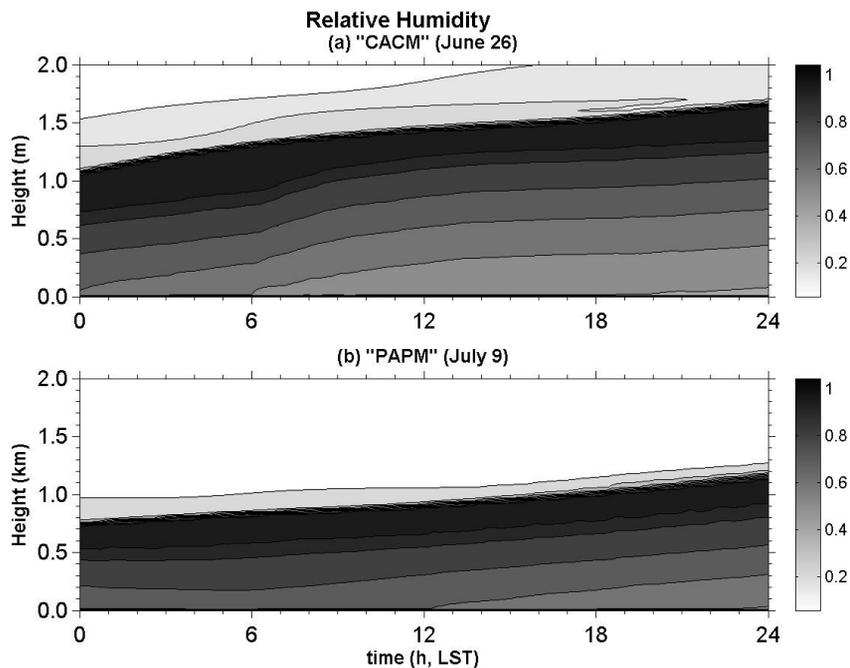


Fig. 10. Time-height cross section of the horizontally-averaged relative humidity (RH) for the “CACM” case **(a)** and for the “PAPM” case **(b)**.

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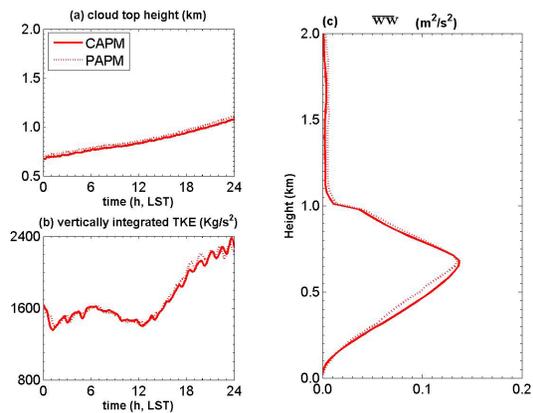


Fig. 11. Same as in Fig. 9 but for the “CAPM” case (solid) and for the “PAM” case (dotted).

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