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Impact of prescribed SSTs on climatologies and long-term trends in CCM simulations

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Abstract

Chemistry-Climate Model (CCM) simulations are commonly used to project the past and future development of the dynamics and chemistry of the stratosphere, and in particular the ozone layer. So far, CCMs are usually not interactively coupled to an ocean model, so that sea surface temperatures (SSTs) and sea ice coverage are prescribed in the simulations. While for future integrations SSTs have to be taken from precalculated climate model projections, for CCM experiments resembling the past either modelled or observed SSTs can be used. This study addresses the question to which extent atmospheric climatologies and long-term trends for the recent past simulated in the CCM E39C-A differ when choosing either observed or modelled SSTs. Furthermore, the processes of how the SST signal is communicated to the atmosphere, and in particular to the stratosphere are examined. Two simulations that differ only with respect to the prescribed SSTs and that span years 1960 to 1999 are used.

Significant differences in temperature and ozone climatologies between the model simulations are found, but long-term trends over 40 years in annual mean temperature and ozone differ only in the troposphere, where temperatures are directly influenced by the local SST trends. However, differences in trends are found on shorter time scales and the results suggest that these differences in trends are induced by associated SST trends. The SST trends lead to modifications in planetary wave activity and therefore a modulation of the Brewer-Dobson Circulation (BDC). This results in time series of tropical upwelling, as a measure of the strength of the BDC, differing strongly between the simulations. A reverse from negative to positive trends is found in the simulation using observed SSTs while trends are positive throughout the simulation when using modelled SSTs.

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1 Introduction

Chemistry-Climate Models (CCMs) were used extensively in the last decade to study the evolution of the global ozone layer and to identify processes leading to changes in ozone in the past and future. For the World Meteorological Organization's scientific assessments of ozone depletion (e.g. WMO, 2007) and as a part of the SPARC CCM Validation Activity CCMVal (Eyring et al., 2005; Eyring et al., 2008), common model strategies have been defined including transient simulations designed to closely resemble the past ("REF1") and to predict possible future developments of the ozone layer ("REF2" and "SCN2", without and with including natural variability, namely the QBO and the 11-year solar cycle). While in REF1, the boundary conditions including natural and anthropogenic variability are mainly deduced from observations, the REF2 or SCN2 scenarios are subjected to the use of future projections of the development of the boundary conditions. As lower boundary condition, sea surface temperatures (SSTs) and sea ice coverage are prescribed to the CCM simulations. For the REF2/SCN2 scenarios, these are commonly taken from climate model projections that include an interactively coupled ocean (i.e. atmosphere ocean general circulation models, AOGCM). As all models, AOGCMs are not capable of reproducing the reality perfectly, i.e. there exist model biases that are also reflected in the SSTs (see e.g. Meehl et al., 2007; Newman, 2007). This raises the question of how CCM simulations are affected by the SST differences introduced by using modelled instead of observed SSTs.

The modelled stratosphere is known to be highly sensitive to the prescribed SSTs. For example, Braesicke and Pyle (2004) studied the impact of inter-annual variability in SSTs by comparing a simulation with long-term mean SSTs to a simulation with transient SSTs. The results showed that the variability in SSTs can act to induce ozone trends in the stratosphere due to a modification of the wave activity flux entering the stratosphere. Other studies dealt mainly with the impact of the El Niño-Southern Oscillation (ENSO) on the stratosphere (Manzini et al., 2006; Sassi et al., 2004) and

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found significant anomalies in ENSO warm phases. Recently, a connection between SSTs and the lower stratosphere was suggested for the tropical regions (Deckert and Dameris, 2008a,b; Rosenlof and Reid, 2008). In Deckert and Dameris (2008a), SSTs were linked to tropical upwelling via deep convection, so that changes in SSTs might drive changes in upwelling. While this study is based on CCM simulations, similar evidence of a linkage between tropical SSTs and the lower stratosphere was found from radiosonde observational data by Rosenlof and Reid (2008). The issue of predicted changes in tropical upwelling (as the tropical branch of the Brewer-Dobson Circulation) in a changing climate was addressed frequently in the recent past (e.g. Butchart et al., 2006), but the reasons for changes are still widely unknown. If SSTs can act as a driver of upwelling, predicted changes of the Brewer-Dobson Circulation (BDC) might be strongly affected by the use of different SST data.

In this study, two CCM simulations starting in 1960 and corresponding to the REF1 and SCN2 scenarios, respectively, are compared. For the period 1960 to 1999, both simulations use the same boundary conditions except of the SSTs (including sea ice coverage; in the following, the term SSTs always includes sea ice coverage). Thus the comparison of these simulations can give an estimate of the bias introduced in the CCM simulation by using different SST data. Since for future predictions only modelled SSTs can be used, the issue of how the SST differences between observed and modelled data sets affect both the climatological mean state and trends and variability of the modelled atmosphere is highly relevant to assess the uncertainty of future projections. For example, when comparing a past period of the REF1 simulation with a future period from REF2/SCN2, one has to be aware of the fact that the changing pattern is a superimposition of changes due to changing boundary conditions with time (“climate change”), and differences due to the use of a different SST data basis. Since the real future development of SSTs is unknown, we can refer only to the past to obtain a rough idea of the quantity of uncertainties due to model deficits in capturing the real SST development. Apart from that, the comparison of the two model simulations can be used to learn about the processes of how SSTs act to influence the atmosphere,

and in particular the stratosphere, which we focus on in this study.

The paper is structured as follows. After a brief model description and a discussion on the SSTs in Sect. 2, some methods used in the analyses are given in Sect. 3. Section 4 compares the climatological mean state of the REF1 and SCN2 model simulation, while in Sect. 5 differences in the long-term evolution are investigated.

2 Model description and SSTs

In this study, the CCM ECHAM4.L39(DLR)/CHEM/ATTILA (E39C-A) is used. The model is an updated version of ECHAM4.L39(DLR)/CHEM (E39C) (Hein et al., 2001; Dameris et al., 2005) with the former semi-Lagrangian advection scheme replaced by the fully Lagrangian advection scheme ATTILA (Reithmeier and Sausen, 2002). The use of the fully Lagrangian advection scheme ATTILA for tracer-transport improved the model performance substantially (Stenke et al., 2008a; Stenke et al., 2008b). E39C is based on the spectral general circulation model ECHAM4.L39(DLR) (Land et al., 2002) and the chemistry-module CHEM (Steil et al., 1998). The spectral horizontal resolution of the model is T30, corresponding to approximately $3.75^\circ \times 3.75^\circ$ on the transformed latitude-longitude grid. In the vertical, the model consists of 39 layers, extending from the surface to the uppermost layer which is centred at 10 hPa. The chosen time step is 24 min. The chemistry module CHEM is based on a generalised family concept and includes homogeneous and stratospheric heterogeneous ozone chemistry and the most relevant chemical processes for describing the tropospheric background chemistry. For more details on E39C-A, see Stenke et al. (2008b).

The experimental set-up of the two transient simulations used in this study correspond to the REF-B1 and SCN-B2d scenarios as defined for the next WMO ozone assessment (Eyring et al., 2008). The REF1 simulation used here is the same as described in detail in Stenke et al. (2008b). Briefly, the external forcings defined as boundary conditions are the 11-year solar cycle, the quasi-biennial oscillation, sulphate aerosol loadings and radiative effects of major volcanic eruptions (namely Agung, El

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Chichón and Mount Pinatubo), the loading of ozone depleting substances, greenhouse gas concentrations and natural as well as anthropogenic NO_x emissions.

The SSTs in the REF1 simulation are prescribed as monthly means following the updated global sea ice coverage and sea surface temperature data set HadISST1 provided by the UK Met Office Hadley Center (Rayner et al., 2003). Following a 10-year spin-up period, the simulation spans the time period 1960 to 2004.

The SCN2 experiment, which spans years 1960 to 2049, is set up in exactly the same way as the REF1 simulation for years 1960–1999 except of the prescribed SSTs. After 1999, the boundary conditions for the SCN2 simulation follow future scenarios, so that years 2000 to 2004 of REF1 do not correspond to the same years in SCN2 (e.g. in terms of phase of the solar cycle). Therefore, in the paper it will be focused on years 1960 to 1999 only. Instead of the HadISST1 data set, SSTs and sea ice in SCN2 are taken from the HadGEM1 general circulation model (Martin et al., 2006; Johns et al., 2006). The SSTs for 1960 to 1999 are from the HadGEM1 transient simulation with anthropogenic forcing only (Stott et al., 2006, 'ANTHRO' in their Table 1), i.e. excluding volcanic eruptions and the solar cycle. The future SST projections used for SCN2 (years 2000 to 2049) are from the SRES A1B run of HadGEM1, which is initiated from ANTHRO so that a consistent SST data set is obtained. These HadGEM1 simulations are part of the World Climate Research Programme's (WCRP's) Coupled Chemistry Climate Intercomparison Project phase 3 (CMIP-3) multimodel dataset used for the 4th Intergovernmental Panel on Climate Change (IPCC) Assessment Report and were provided by the Program for Climate Model Diagnosis and Intercomparison (PCMDI, available at www.pcmdi.llnl.gov).

A comparison of the HadGEM1 and HadISST SST climatologies is shown in Fig. 1 (see also Johns et al., 2006). The difference pattern is dominated by lower SSTs in HadGEM1 in the tropics, the subtropics and northern mid-latitudes, where the differences are up to 3 K. Smaller regions of higher SSTs in HadGEM1 are found on the west coasts of South America and Africa, as well as on the east coasts of North America and Asia in northern winter (i.e. the storm track regions). There are also large

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differences in sea ice extent, with generally more sea ice in HadGEM1. Time series of anomalies of global and hemispheric mean SSTs from HadISST and HadGEM1 are compared in Fig. 2. Both the global mean and Southern Hemisphere mean SSTs show generally rather constant values in the 1960s, and rising SSTs since about the mid-70s. However, in the Northern Hemisphere the SST data set deduced from observations show generally higher SSTs in the 1960s, a strong decrease in the 1970s and rising temperatures only after the late 1970s. The modelled time series does not reproduce this behaviour, here SST anomalies are negative throughout the 1960s. In terms of variability, the interannual variance due to the El Niño-Southern Oscillation (ENSO) is underestimated and the ENSO signal is too regular (Johns et al., 2006) in the HadGEM1 model. A detailed study of annual and decadal variability of SST eigenmodes from both observations and several GCMs including HadGEM1 revealed large deficits of the models to capture the variability, and showed that the patterns are less persistent than in reality (Newman, 2007).

3 Methods of data analyses

3.1 Trend calculations and tests on significance

In the following, both differences between climatologies (i.e. mean values) and long-term changes are analysed. The statistical significance of differences is tested with the Student-T-test (see e.g. von Storch and Zwiers, 2002). To quantify long-term changes, trends are calculated by applying a linear least squares fit of the time series $y(t)$ to a linear function (i.e. $\hat{y}(t)=a+bt$ ($t=1\dots N$), where N is the number of timesteps). The uncertainty in the trend (i.e. the slope of the linear function b) is then given by $\sigma^2 = \frac{\chi^2}{N-2} \frac{1}{S_{tt}}$ where χ^2 is the sum of the squared residuals ($\chi^2 = \sum_{t=1}^N [y(t) - \hat{y}(t)]^2$) and $S_{tt} = \sum_{t=1}^N t^2 - N(\frac{N+1}{2})^2$. The trend is statistical significant if the null-hypothesis $b = 0$ can be rejected, which is tested by comparing the T-test value $t=b/\sigma$ against the critical

value of the T-distribution with $N-2$ degrees of freedom.

3.2 Eliassen-Palm fluxes

As a measure of wave activity, Eliassen-Palm (EP) fluxes were calculated using the full ageostrophic form (e.g. Andrews et al., 1987). The fluxes were scaled by the mass of an annular ring of a certain latitude band and pressure interval (see Edmon et al., 1980, their Eq. 3.12). Using pressure coordinates, the EP flux F at latitude ϕ and pressure level p is then given by:

$$F(\phi, p) = 2\pi r_e \cos(\phi) / g \times \left[r_e \cos(\phi) \left(-\overline{u'v'} + \frac{\overline{v'\Theta'}}{\Theta_p} u_p \right), \right. \\ \left. r_e^2 \cos(\phi) \left((f - u_\phi) \frac{\overline{v'\Theta'}}{\Theta_p} - \overline{u'\omega'} \right) \right] \quad (1)$$

where r_e is the Earth's radius, g the gravitational constant, Θ the potential temperature and $f = \Omega \sin(\phi)$ the Coriolis parameter (Ω is the Earth's angular speed of rotation). Subscripts denote the partial derivative of the variable. The EP divergence is then $\nabla \cdot \mathbf{F} = \partial F_\phi / \partial \phi + \partial F_p / \partial p$. Due to the mass scaling, the EP flux is in units [$\text{m}^3, \text{m}^3 \text{Pa}$] and the divergence in [m^3], which equals Nm/Pa , i.e. energy/pressure (Edmon et al., 1980).

3.3 Tropical upwelling and downward control

The meridional residual circulation ($\overline{v^*}, \overline{\omega^*}$) as introduced in the transformed Eulerian mean (TEM) equations (e.g. Holton, 2004) is used here for quantifying the large-scale circulation in the stratosphere. As a measure of the strength of the upwelling branch

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of the circulation, mean tropical upwelling is calculated by integrating over all latitudes within 60° N/S where the residual vertical velocity points upward (i.e. negative $\overline{\omega}^*$ since $\overline{\omega} = \frac{\partial p}{\partial t}$):

$$U_{\text{tropics}} = \int_{\phi_{\overline{\omega}^* < 0}} \frac{2\pi \cos(\phi) r_e^2}{g} \overline{\omega}^* d\phi \quad (2)$$

In addition to this direct calculation from the residual vertical velocity, tropical upwelling was calculated using the downward control (DC) principle (Haynes et al., 1991). The DC principle states that the residual circulation at a certain level is driven by the forcing on the mean flow above this level. By using the EP divergence as forcing, the contribution of planetary wave activity in driving the BDC can be quantified.

The DC calculations were performed as follows. From the continuity equation for $(\overline{v}^*, \overline{\omega}^*)$, the residual streamfunction $\overline{\chi}^*$ can be introduced:

$$(\overline{v}^*, \overline{\omega}^*) = \frac{1}{\cos(\phi)} \left(-\frac{\partial \overline{\chi}^*}{\partial p}, \frac{1}{r_e} \frac{\partial \overline{\chi}^*}{\partial \phi} \right) \quad (3)$$

Using the TEM momentum equation and assuming a non-zero meridional gradient of the zonal mean angular momentum $\overline{m} = r_e \cos(\phi)(\overline{u} + r_e \Omega \cos(\phi))$ and stationary conditions, the residual streamfunction $\overline{\chi}^*$ can be expressed as (for details see Haynes et al., 1991):

$$\overline{\chi}^*_{\overline{m}_0}(p) = \int_p^0 \left[\frac{r_e^2 \cos(\phi)^2 G}{\partial \overline{m} / \partial \phi} \right]_{\phi = \phi(\overline{m}_0)} dp' \quad (4)$$

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Here G is assumed to be the EP divergence (appropriately rescaled to units of m/s^2), and the subscript $\phi = \phi(\bar{m}_0)$ denotes integration along lines of constant angular momentum. The condition $d\bar{m}/d\phi \neq 0$ limits the latitude region of the calculation to outside the inner tropics. In practice, the integration along constant angular momentum can only be performed when \bar{m}_0 is a real function of pressure, i.e. the \bar{m}_0 contour is not closed in the domain of calculation.

For estimating tropical upwelling as given in Eq. 2, it is used that $\bar{\omega}^* = 1/(r_e \cos \phi) \bar{\chi}^*_\phi$ so that upwelling between latitudes ϕ_1 and ϕ_2 is given by $2\pi r_e/g[\bar{\chi}^*(\phi_2) - \bar{\chi}^*(\phi_1)]$ (Randel et al., 2002). In order to capture the same latitude band as for the direct tropical upwelling calculation (Eq. 2), the turnaround latitudes (where $\bar{\omega}^*$ changes sign) are deduced from the residual vertical velocity for each month and the streamfunction is then calculated at these latitudes. In a few cases, the turnaround latitudes were too deep in the tropics, so that the integration along \bar{m}_0 was not possible. In these cases, tropical upwelling was set to zero. The mean location of the turnaround latitudes did not differ between the REF1 and the SCN2 simulation in a statistical significant way. Also, trend analyses of the turnaround latitudes showed that there are no significant changes in the location and the width of the upwelling region in both simulations. Therefore changes in total upwelling can not be attributed to changes in the width of the upwelling region. However, due the relatively coarse horizontal resolution of the model (T30), changes of the width of the tropical belt in the order of one to two degrees per decade, as reported by Seidel and Randel (2007), might not be detectable here.

4 Climatological differences

4.1 Temperature and ozone

To address the question to which extent CCM simulations using observed and modelled SSTs, respectively, differ in their climatological mean state, the 40-year climatologies (1960–1999) of temperature and ozone mixing ratios are compared (Fig. 3).

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The tropospheric temperatures in the SCN2 simulation are significantly lower than in REF1, with differences maximising in the tropical middle troposphere. There the temperature is up to 1.5 K lower in SCN2. In the stratosphere, the temperature differences are generally weaker but temperatures are significantly higher in SCN2 in the tropical lower stratosphere and at high latitudes at around 200 hPa. The temperature differences in the troposphere between the two simulations are of the same order of magnitude as changes over the 40 years of each model simulation (see Fig. 6). The large differences in the climatologies are not unexpected since differences between the SST climatologies of the two runs are more than 1 K in the zonal mean for most latitudes (see Fig. 1), and even up to 3 K at about 50° N.

The differences in the ozone climatology show significantly higher ozone mixing ratios in the tropical lower stratosphere and lower ozone mixing ratios at southern middle and northern high latitudes in SCN2 compared to REF1. Additionally, there are significant differences in the troposphere with higher ozone mixing ratios in the tropics and at high latitudes. The difference pattern in ozone mixing ratio in the stratosphere is clearly indicative of weaker transport from the tropics into higher latitudes in SCN2. To examine the differences more closely, Fig. 4 (left panel) shows the annual cycle of zonal mean total column ozone for all latitudes (from REF1, 1960–1999) and respective differences between SCN2 and REF1. The negative ozone differences are strongest in northern winter at high latitudes (maximising in February) and in southern mid-latitudes in May to October, occurring in both hemispheres one to two month before total ozone maximises.

4.2 Residual circulation

The distribution of total column ozone is determined by both chemical and dynamical processes. While the photochemical production of ozone is highest in the tropics, the BDC is transporting ozone from the tropics towards higher latitudes. The classical view is that transport by the dominant winter cell of the BDC leads to the ozone maximum in middle to high latitudes of the winter/spring hemisphere. However, next to the direct

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transport of ozone by the BDC, also other mechanism might play a role, like eddy mixing of tropical air masses into mid-latitudes or the local ozone chemistry. Grewe (2006) showed that in the mid-latitudes the contribution of ozone originating in the tropics is about half of the total amount of ozone, with the other half being produced locally. Detailed analyses of the processes important to the redistribution of ozone, and in particular the relation of changes in these processes to changes in ozone concentrations are not scope of this study and are subject of ongoing work. However, in the following the classical view of poleward transport of ozone with the BDC will be adopted. Jiang et al. (2007) showed that changes of the BDC can be related to changes in the amount of ozone with reduced ozone concentrations in the tropics and enhanced ozone concentrations in mid-latitudes.

Figure 4 shows next to the ozone climatology of REF1 (left) the climatology of the streamfunction of the residual velocity $\overline{\chi^*}$ at 52 hPa, as measure of the BDC (right). The streamfunction clearly shows the dominant winter circulation cells, with a maximum of the Northern Hemisphere circulation in January and of the Southern Hemisphere circulation in June. Total column ozone maximises poleward and about 1–2 month after the circulation maxima.

The difference pattern in the residual circulation is rather patchy. While there is a strong negative anomaly in the Northern Hemisphere in January when the residual streamfunction is at its maximum, the anomaly is positive in March. However, the strong negative anomaly in January relates very well to the negative total column ozone anomaly seen in February, and the positive anomaly in transport in March leads to weak positive differences in total column ozone in April at high latitudes.

In the Southern Hemisphere, where negative values of the residual streamfunction indicate poleward transport, a positive anomaly is apparent at subtropical latitudes (about 0° – 30° S) from January through June. From March on, when the meridional circulation turns to the Southern Hemisphere, the positive anomaly corresponds to a weakening of poleward transport. This relates well to the negative total ozone anomaly maximising in July.

4.3 Planetary wave activity

Since the large-scale circulation in the stratosphere is largely driven by planetary waves, the differences in the meridional circulation as found in the last section are expected to result from differences in wave activity. The climatology of EP fluxes of the REF1 simulation was shown in Stenke et al. (2008b), and showed a generally realistic representation of planetary wave (PW) activity in the model. Differences between SCN2 and REF1 in EP divergences and fluxes are shown in Fig. 5 for January, where differences in the residual circulation are largest in the Northern Hemisphere. It can be seen that there is less wave activity entering the stratosphere in northern extratropics, resulting in less wave dissipation and therefore a weaker forcing of the BDC in the SCN2 simulation. Following the EP flux anomaly downward, the negative differences in the stratosphere seem to follow from less wave generation in the troposphere at latitudes poleward of 50° N. Even though there is a strong positive anomaly in wave generation equatorward of 50° N, stronger divergences at around 400 hPa lead to less wave activity reaching the stratosphere.

5 Long-term trends and variability

While in the last section, the climatological differences due to different SST data sets were studied, the following part of the paper aims to analyse differences in the temporal evolution.

5.1 Temperature and ozone

The linear trend in annual mean zonal mean temperature and ozone mixing ratios over the years 1960 to 1999 was calculated for each latitude and pressure level, and the resulting trend patterns are shown in Fig. 6. The results show that the trend patterns calculated for REF1 and SCN2, respectively, are similar both in terms of geographical

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locations and magnitudes of trend maxima and minima for temperature as well as for ozone.

In terms of temperature, both simulations show a warming of approximately 1 K over 40 years in the troposphere and a cooling pattern in the stratosphere with strongest cooling of up to 4.5 K near the south pole and a local cooling maximum in the tropics. The temperature trends are significantly different (on the 95% level) only in the Northern Hemisphere mid-latitudes in the troposphere, consistent with weaker SST trends in the Northern Hemisphere mean HadISST data as in HadGEM1 (see Fig. 2).

The ozone trend pattern shows a strong decrease of ozone mixing ratios in the stratosphere, especially in high southern latitudes and an increase in the troposphere. The differences between the two simulations in the 40-year zonal mean ozone trends are not statistically significant for all latitudes and levels.

To study the temporal evolution of temperature and ozone more closely, time series of total column ozone and 30 hPa temperatures (as representative for the lower stratosphere) for certain latitude bands are shown in Figs. 7 and 8. Only winter/spring months are shown since the stratospheric variability is highest in this time of the year. Also, SSTs are expected to influence the stratosphere mainly through differences in planetary wave activity, which is communicated to the stratosphere only in winter when wave propagation is not blocked by easterly winds.

Much of the interannual variability in ozone can be explained by external forcings like the QBO and the solar cycle (e.g. Steinbrecht et al., 2006). The co-variability of the time series of the two simulations, as measured by the correlation coefficient, gives a measure of how deterministic the system is, i.e. how strong the common external forcings (like the QBO or the solar cycle) act to determine the anomalies in individual years. The correlation of the detrended time series of total column ozone (Fig. 7) is significant for the tropics and for the Southern Hemisphere (with correlation coefficients of approximately 0.5), while in the Northern Hemisphere the high dynamic variability causes the correlation not to be significant. Since both model runs are driven by the same boundary conditions (apart from the SSTs), anomalies that are driven by e.g. the

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QBO or the solar cycle should be apparent in both simulations. For example, the very low total column ozone value over southern mid-latitudes in 1985 is reproduced by both model runs. The negative anomaly was a combined effect of the QBO, the solar cycle and their timing in the year (Bodeker et al., 2007).

5 The temperature time series of the SCN2 and REF1 runs as shown in Fig. 8 correlate significantly (after being detrended) in the tropics and in the mid-latitudes of both hemispheres. The correlation coefficients are 0.58, 0.47 and 0.35 in the northern mid-latitudes, the tropics and the southern mid-latitudes, respectively. Especially in the tropics, the variability seems to be largely dominated by the QBO, which is apparent particularly in the last decade of the time series.

10 The essential question when comparing the REF1 and SCN2 simulations is on the reproduction of the temporal evolution of the atmosphere. Above, it was shown that the 40-year trend pattern in ozone and temperature are not significantly different from each other. But is this still true when the length of the time series is shortened? To answer this question, the linear trend in total column ozone and temperature was calculated for periods of increasing length, starting with 1960 to 1974 and extending the period in steps of one year until 1999 (see right panels of Figs. 7 and 8). If the trend would be strictly linear, it would be constant for all periods. For example, if total ozone was increasing during the first 20 years and decreasing with the same rate during the second half of the total period, the trend would be positive during the first 20 years, then decrease and equal out to zero at the end.

15 The results for total column ozone (Fig. 7) show good agreement in trends calculated from the two simulations in southern high latitudes, where the strongest negative trend appears due to rapid ozone depletion. The trend increases for longer periods since the rapid ozone depletion only started in the 1980s, so that the trend is close to zero for a short period and increases afterwards. In the mid-latitudes and tropics, the absolute values of the ozone trends are very similar when including more than about 25 years of model simulation. In the tropics, the trend estimated for a period spanning from 1960 to around 1980 is positive in the REF1 run but close to zero or negative for the SCN2

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run. In turn, the trends at middle and high northern latitudes in the same period are strongly negative in REF1 in the late 70s, but close to zero in SCN2. This leads to the trend calculated over the whole period being weaker in REF1 as in SCN2.

The trend calculations in the same manner for the temperature time series give similar results (see Fig. 8 right panels). While the linear trend calculated over time periods including more than 25 years agree quite well between the model runs, the trends calculated for shorter periods differ both in magnitude and sign (for the tropics). Again, the trends in northern high latitudes are strongly negative for periods 1960 to around 1979 for REF1 and close to zero for SCN2. The behaviour is reversed in the tropics (with positive trends in REF1 and small (negative) trends in SCN2), even though here the differences in trends are barely significant.

This difference in trends during the period 1960 to the late 1970s in both total column ozone and temperature might either be caused by (i) chance, due to internal dynamic variability or (ii) a deterministic process, forced externally in the REF1 simulation but not in the SCN2 simulation. If it was a process forced by an external boundary condition, the only possible boundary condition to cause this modulation in trends are the prescribed SSTs, since this is the only difference in the external forcing of the REF1 and SCN2 simulations. As shown in Fig. 2, the time series of SSTs differ in particular in the Northern Hemisphere in the 1960/1970s. Therefore, we suspect that changes of SSTs alter the temporal evolution of stratospheric ozone and temperature in a way similar to the climatological differences found in Sect. 4. The mechanism of how SSTs are expected to act on stratospheric temperature and ozone distributions is via dynamical processes, i.e. the modulation of PW activity. Indeed, the fact that positive trends in the 1960/1970s (in both temperature and ozone) in the tropics seem to be balanced by strong negative trends in middle and high latitudes in the winter hemisphere suggest that the underlying mechanism is related to changes in the large-scale circulation (which is in turn driven by PW activity). In order to investigate the hypotheses of SST-induced dynamical changes, the temporal evolution of the BDC and PW activity for the two simulations with different underlying SSTs will be investigated in the next section.

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5.2 Brewer-Dobson circulation and planetary wave activity

To quantify changes in the large-scale circulation of the stratosphere, time series of mean tropical upwelling for December to February (DJF) are shown in Fig. 9 for REF1 and SCN2. Significant changes in tropical upwelling with time were found only for DJF, so that in the following only northern winter months are considered. As Fig. 9 shows, the temporal evolution of tropical upwelling is fairly different in the two simulations. While in SCN2, tropical upwelling increases throughout the period 1960 to 1999, in REF1 a clear reverse in the trend is seen in the late 1970s. As mentioned above, the temperature and ozone time series are indicative of a weakened BDC in the REF1 simulation in the 1960s and 70s, and the tropical upwelling time series strongly supports this suggestion. As discussed above, the fact that the reverse in trends only occurs in the simulation with observed SSTs strongly suggests that differences in the temporal evolution of the SSTs drive these changes in the large-scale circulation. However, the most likely mechanism for SSTs affecting the BDC is via changes in planetary wave generation or propagation as the main drivers of the BDC. Therefore, in the following the contribution of planetary wave activity in forcing upwelling and changes in the wave activity are analysed.

The contribution of planetary wave activity to forcing changes in upwelling can be quantified by calculating tropical upwelling via the downward control principle (see Sect. 3). Using this method, the trends in upwelling can be reproduced, especially the strong negative trend in DJF in the first 20 years in REF1. This shows that the changes in the BDC can be linked to changes in the dissipation of planetary wave activity in the stratosphere. PWs are largely generated in the troposphere, propagate upward and dissipate due to respective background conditions (see e.g. Charney and Drazin, 1961). Therefore, changes in wave activity divergence can be caused by either changes in the generation, the propagation and/or the region of dissipation of waves.

Figure 10 shows the linear trend in EP fluxes and divergence and in the streamfunction of the residual circulation for the REF1 simulation over the period 1960 to 1978 in

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January, where tropical upwelling decreases. For SCN2, trends over the period 1960 to 1999, where tropical upwelling increases, are shown. For REF1, the period 1978 to 1999 also shows an increase in tropical upwelling and the trend pattern in the residual circulation and EP fluxes are qualitatively similar to the SCN2 1960–1999 period, even though the shorter time period makes it more difficult to obtain significant trends.

The REF1 1960–1978 trend pattern of the residual velocity shows a strong significant hemisphere-wide decrease in the circulation strength of the Northern Hemisphere (Fig. 10 left panel). This weakening in the residual circulation is clearly linked to a negative trend in EP divergence in the extratropical stratosphere above 60 hPa. The trend in EP fluxes shows that the origin of the EP divergence change lies in the tropospheric production of wave activity. The decrease in tropospheric wave generation is found in northern mid-latitudes, in particular between 35° N and 60° N.

For the period 1960 to 1999 in SCN2 a clear increase in the residual streamfunction and therefore in the strength of the BDC is seen in both hemispheres (Fig. 10 right panel). The trends in the residual streamfunction are significant in the (sub-)tropics below 80 hPa, with rising motion around the equator and downwelling equatorward of about 20° N/S. Also in the Southern Hemisphere at middle and high latitudes, the trend towards a stronger BDC is significant. The (sub-)tropical circulation changes can be associated with increases in EP convergence in the same region. In the mid-latitudes, both hemispheres show enhanced EP flux entering the stratosphere and therefore enhanced EP convergence at higher levels, but the changes are statistically significant only in the Southern Hemisphere.

The changes in EP flux propagating into the southern stratosphere in mid-latitudes can not be attributed to stronger wave generation but the trend pattern indicates reduced dissipation in the upper troposphere. The changes in wave activity in southern summer are most likely a result of changes in the dynamics due to ozone loss. For example, Li et al. (2008) suggested a mechanism of increased descent in the Antarctic vortex leading to a stronger meridional temperature gradient and therefore a stronger jet which in turn leads to enhanced upward planetary wave propagation. The fact that

trends in the Southern Hemisphere are absent the period 1960 to 1978 fortifies this hypotheses.

6 Summary and conclusions

The comparison of two transient E39C-A simulations using observed and modelled SSTs, respectively, showed that there exist large differences in temperature and ozone climatologies. These differences are linked to small, but statistically significant differences in the residual meridional circulation, which are in turn driven by differences in PW activity fluxes. Detailed trend analyses showed generally good agreement in long-term trend pattern of zonal mean temperature and ozone mixing ratios. However, differences in trends were found when considering shorter time series.

In the comparison of the two simulations that are identical except of the SST (including sea ice coverage) data sets, it is straightforward to attribute the climatological differences to SST differences. When it comes to time series (and trend) comparison, differences in the temporal evolution in the two simulations could simply be due to internal variability of the system. It would be desirable to have more than one realisation of each simulation to test whether the behaviour is deterministic. However, the following results suggest that the differences in trends between the simulations in the 1960/1970s are induced by differences in the SSTs: (i) there are quite large deviations between the behaviour of the time series of the two simulations, especially in tropical upwelling in northern winter where the reversed trend in REF1 over 20 years is statistically significant, (ii) there occur deviations in the Northern Hemisphere mean SST anomalies between the two data sets in the same time period and (iii) the reverse in trends in the 1960/1970s in REF1 also occurred in earlier simulations with the predecessor model E39C (Dameris et al., 2005, 2006).

Both from inter-simulation and intra-simulation (i.e. time series) comparison, we suggest enhancement or reduction in planetary wave generation, which in turn drives changes in the large-scale circulation in the stratosphere, as the mechanism of how

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SST signals are communicated to the stratosphere. This can be seen in differences of ozone climatologies between the two runs, which show clear signals of changes in transport of ozone from the tropics to higher latitudes. Also the temperature climatology comparison shows significantly higher temperatures in the tropical lower stratosphere, an indication of weaker rising of air masses in this region. The difference pattern in the residual circulation confirms that the BDC is weakened, as a result of less PW flux propagating into the stratosphere. Also the trend pattern of the residual circulation and EP fluxes for the period 1960 to 1979 in REF1 shows a weakened BDC and less PW propagation into the stratosphere, opposing the trend towards a stronger BDC as found in the later period in the same run and in the SCN2 simulation. As discussed above, the trend anomalies in REF1 in the 1960/1970s are attributed to SST trends in this period. Both the climatological EP differences and the 1960-79 EP trends are indicative of less PW generation in the extratropical troposphere. While it is well understood how changes in EP influx drive changes in the BDC in the stratosphere (e.g. Haynes et al., 1991), the question of how SSTs can affect the excitation of planetary waves is poorly understood (Kushnir et al., 2002). Peng et al. (1995) and Peng and Whitaker (1999) showed that the response to SST anomalies is highly sensitive to the background state, i.e. the response to the same anomaly might be of different nature for different month, but also appears different when using different models. Therefore, more studies will be necessary to show how PW generation changes are linked to SST anomalies, and whether there are regions that are of particular importance.

As it was mentioned in the introduction, most climate and chemistry-climate models show an increase of the BDC in a changing climate with rising greenhouse gas (GHG) concentrations. The findings of the study presented here agree in that an increase of tropical upwelling with time in the SCN2 simulation and in REF1 after 1980 is diagnosed.

The trend pattern of the residual streamfunction in the simulations suggest that the changes consist of a strongly enhanced circulation constrained to the (sub-)tropics and of significant changes in the Southern Hemisphere. The former is linked to changes

in EP fluxes within the tropics, and can probably be attributed to the mechanism suggested in Deckert and Dameris (2008a). The latter is forced by stronger extratropical planetary wave flux into the stratosphere, which in turn seems to be caused by changes in EP propagation.

5 Even though the decrease of the BDC in REF1 in the 1960/1970s seems to be driven by changes in PW generation driven by SST changes, from our current findings we can not conclude on SSTs being the drivers of the general BDC enhancement. It could as well be other processes, linked for example to changes in the mean zonal winds due to changed GHG concentrations, that lead to increased extratropical PW activity, and
10 which are superimposed by SST-related modulation. However, it is shown that the temporal evolution of the BDC is dependent on the SST data set used, which suggests that SST trends can alter the BDC.

The fact that the reverse in trends is seen in the simulation using observed SSTs suggests that all the features discussed above should be detectable in data of the real atmosphere. The problem in verifying this directly from measurements is the insufficient data basis for analysing global stratospheric dynamics before 1979. One consistent data set of stratospheric temperatures (only for the Northern Hemisphere) that starts in 1957 and covers the time period to 2001 are the so-called Berlin analyses (Labitzke and collaborators, 2002, available at http://badc.nerc.ac.uk/data/berlin_strat). Labitzke
15 and Kunze (2005) presented trend calculations for 50 hPa, 80° N temperatures from the Berlin data set and two reanalysis data sets, and found trends with opposite sign for the period 1964–1979 than for 1979–2000, but the trends are mostly not statistically significant. Another study, based among other on ECMWF and NCEP reanalysis data, points out a clear reverse of trends around 1980 in the duration of easterlies in the stratosphere during summer (the so-called summer length) and attributed it to changes
20 in wave activity (Offermann et al., 2005). However, direct comparisons of the same diagnostics from these observational studies did not show the same behaviour in the model simulation, which would indicate that the model system is responding to the observed SST anomalies in a different way than the real atmosphere. Since in the

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model the SSTs are prescribed, the interaction between ocean and atmosphere is an one-way only process, so that feedbacks of the atmosphere acting to drive the SST anomalies are not included. This might be a possible explanation why we don't find the same atmospheric behaviour in model data and observations.

Based on the results presented here, we claim that the SST signal is communicated to the stratosphere by modulation of PW activity and therefore the BDC, implying that trends in these quantities are particularly sensitive to changes in SSTs. Trends in stratospheric temperature and ozone are largely driven by changes in concentrations of GHGs and ozone depleting substances, and are altered by SSTs only secondarily. Thus, the offset in the climatologies due to the SST biases can be assumed to not affect future predictions in temperature and ozone to a large degree.

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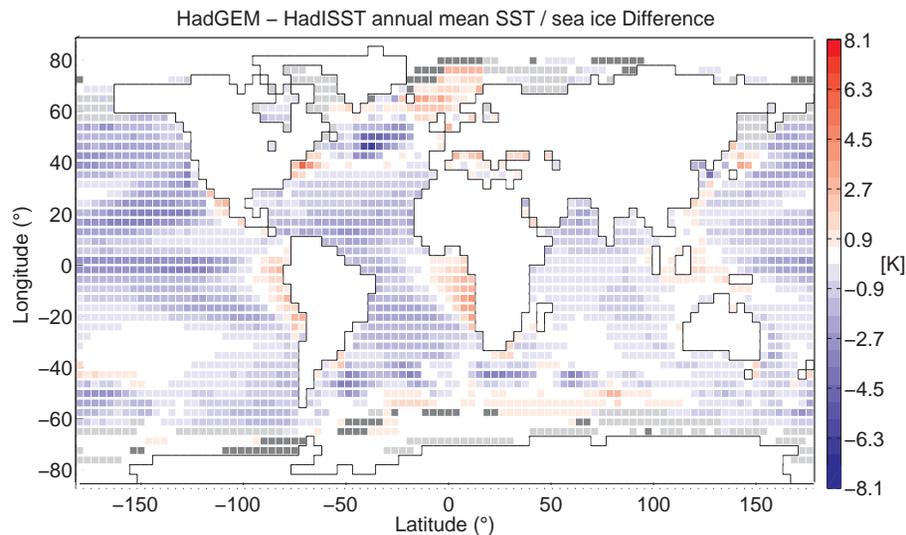


Fig. 1. Difference of 1960 to 1999 annual mean sea surface temperatures [in K] from HadGEM minus HadISST. Differences in sea ice extend are shown in grey, where light grey areas correspond to more sea ice in HadGEM compared to HadISST, and dark grey regions to less sea ice in HadGEM.

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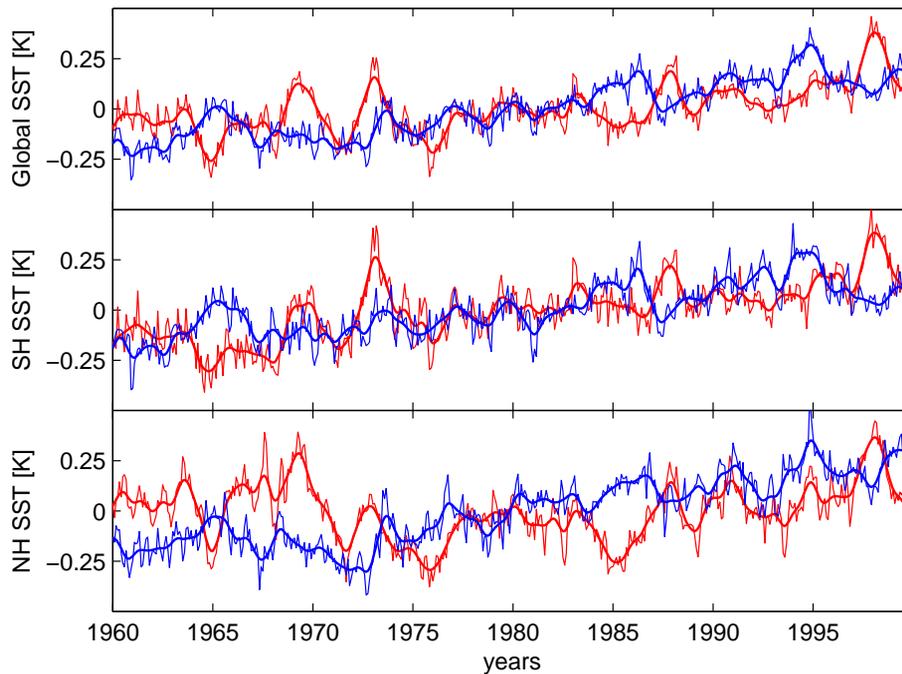


Fig. 2. SST time series with the 1960–1999 mean annual cycle subtracted for REF1 (red) and SCN2 (blue) for global (top), Southern (middle) and Northern (bottom) Hemisphere means. The thick solid lines are smoothed time series (by applying a 1-2-1 filter 20 times).

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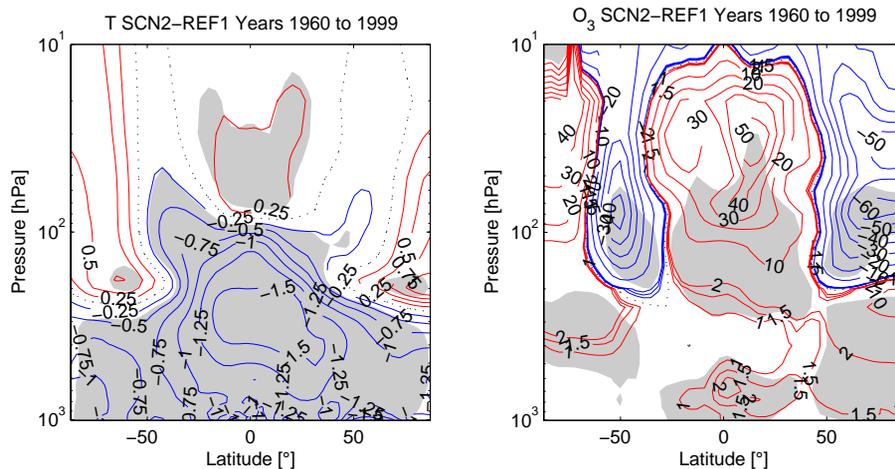


Fig. 3. Difference of 1960 to 1999 temperature [in K] (left) and ozone mixing ratio [in ppbv] (right) climatologies in REF1 and SCN2. Shown are differences SCN2 minus REF1, grey shading indicates statistic significance of the differences (on the 95% level).

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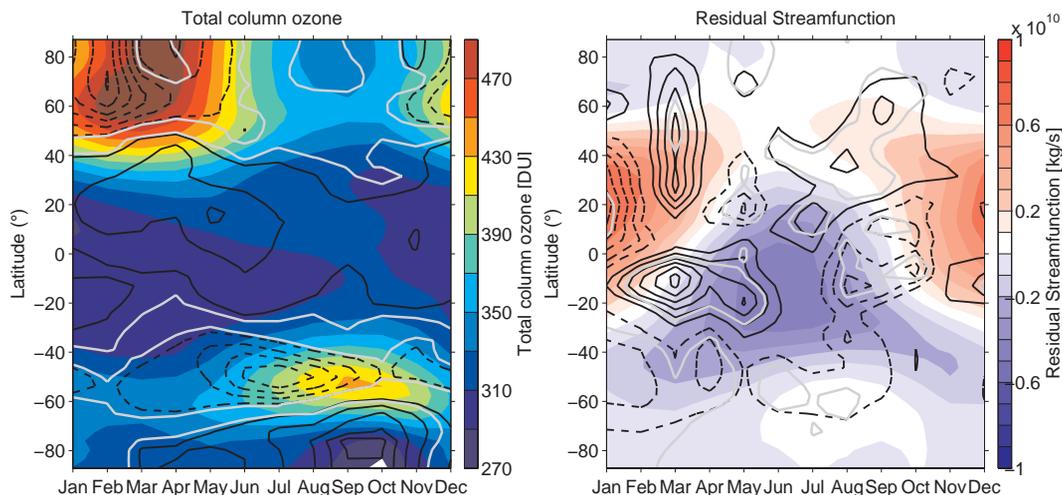


Fig. 4. Total column ozone (left) and residual streamfunction ($\overline{\chi^*}$, see Eq. 3) at 52 hPa (right) 40-year climatologies (coloured) of REF1. Differences SCN2-REF1 are shown as black solid contours (for positive values) and dashed contours (for negative values). The contour interval is 2 DU for total column ozone and 10^8 kg/s for the residual streamfunction. Within light grey contours, differences are statistical significant on the 95% level.

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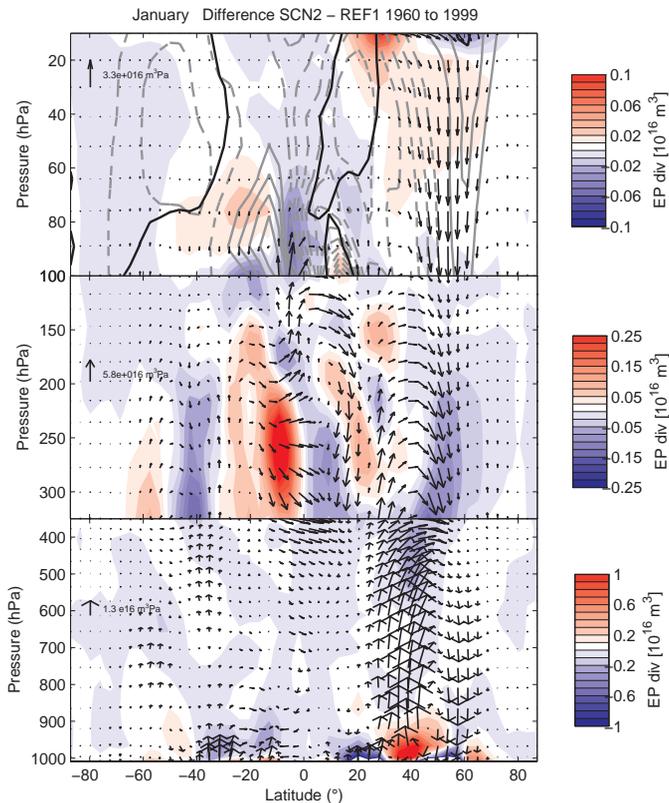


Fig. 5. Differences (SCN2-REF1) of 40-year climatologies of EP fluxes (black arrows) and EP divergences (coloured) for three different height regions. The scaling of the EP fluxes and divergences is indicated by the scale arrow and the colorbar in each panel. Light grey contours in the upper panel show differences (SCN2-REF1) in the residual meridional circulation, with solid lines denoting positive values and dashed lines denoting negative values (contour interval is 10^8 kg/s). Differences in the residual circulation values are statistically significant within black contours.

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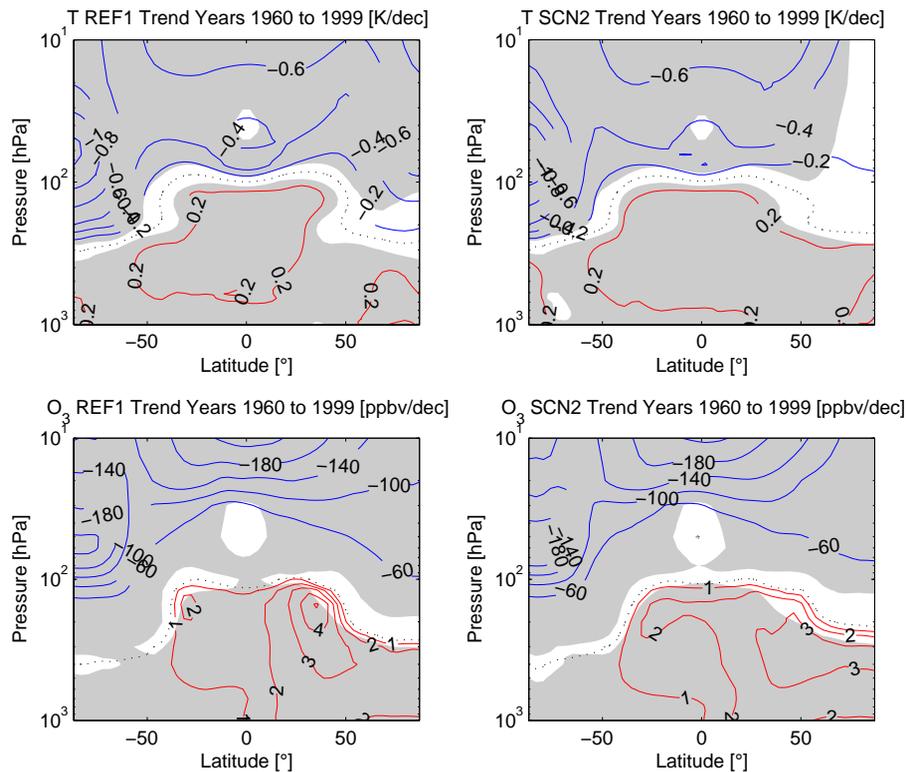


Fig. 6. Linear trends over years 1960 to 1999 of temperature (upper panels) and ozone mixing ratio (lower panels) for REF1 (left) and SCN2 (right). Units are K/decade and ppbv/decade, respectively. Grey shading indicates statistic significance of the trends (on the 95% level).

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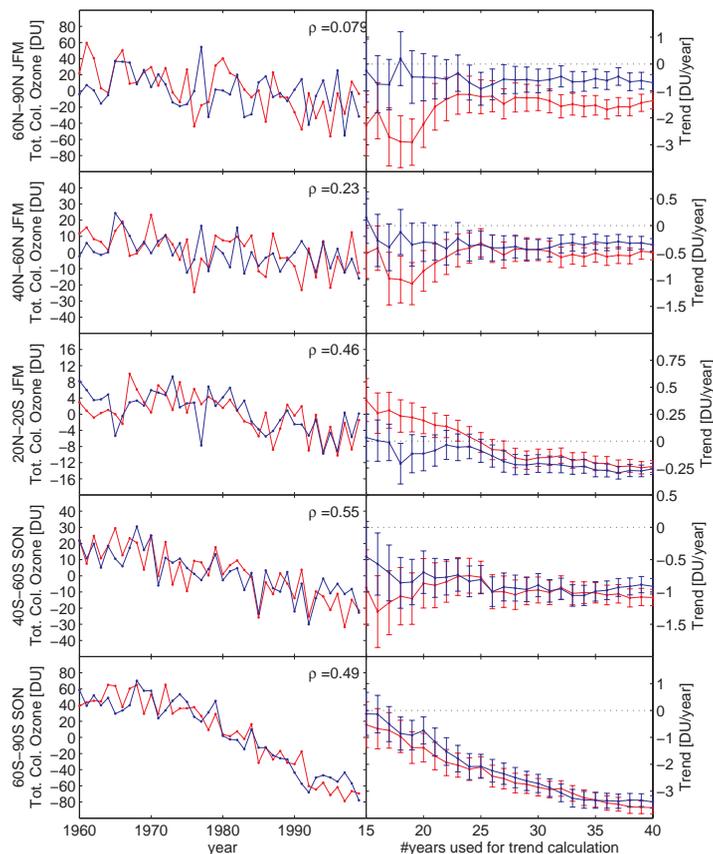


Fig. 7. Mean total column ozone anomalies in [DU] for various latitude bands for REF1 (red) and SCN2 (blue). The number in the upper right corner is the correlation coefficient between the two time series. Right panels show linear trends as a function of the number of years after 1960 used for the trend calculation.

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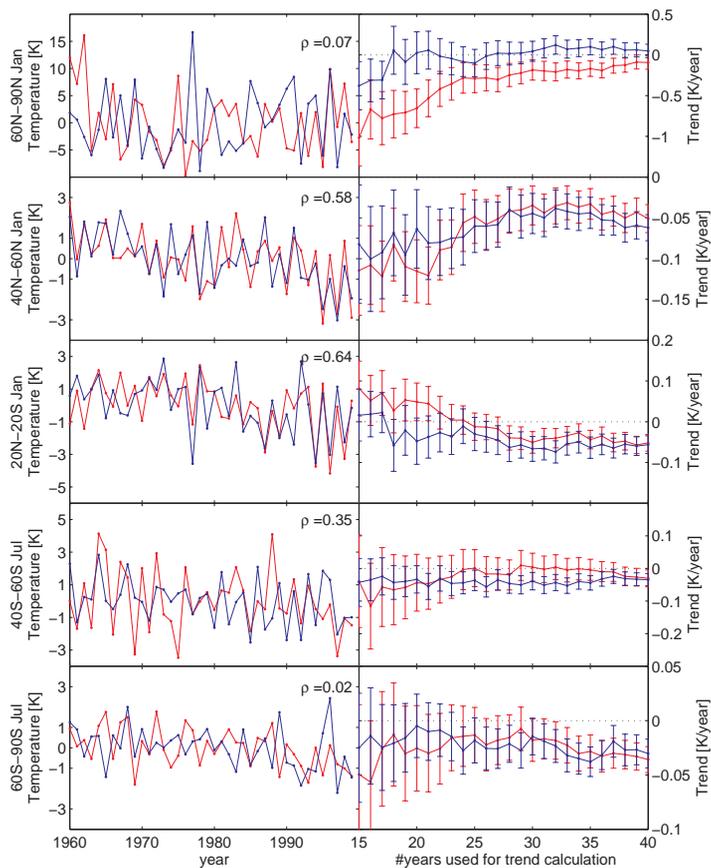


Fig. 8. As in Fig. 7 but for temperature in [K] at 30 hPa.

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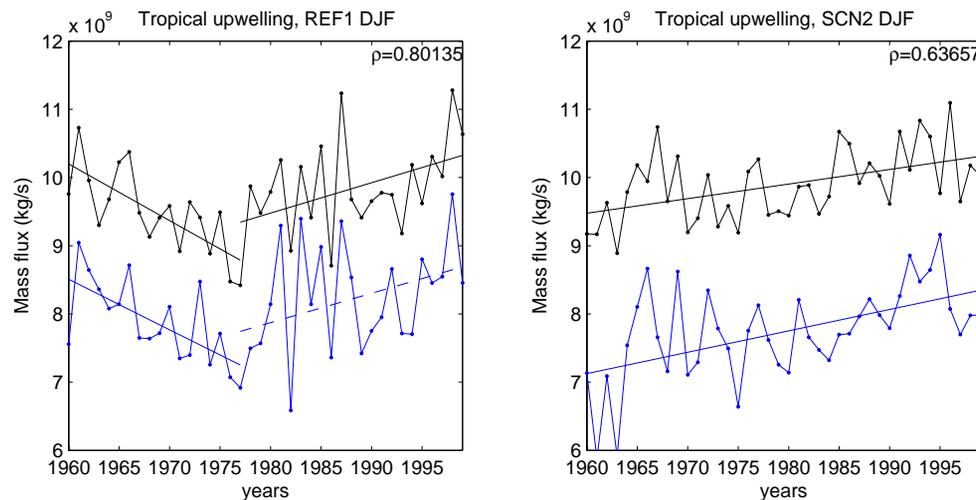


Fig. 9. Mean tropical upwelling (as measured by mass fluxes) at 76 hPa for December to February for REF1 (left) and SCN2 (right). Shown is tropical upwelling derived directly from the residual velocity (black lines) and upwelling calculated using downward control (blue lines). Together with the time series, linear trends are shown (as solid line for significant trends, otherwise dashed). In the upper right corner, the correlation coefficient between the time series of tropical upwelling calculated directly and via DC is given.

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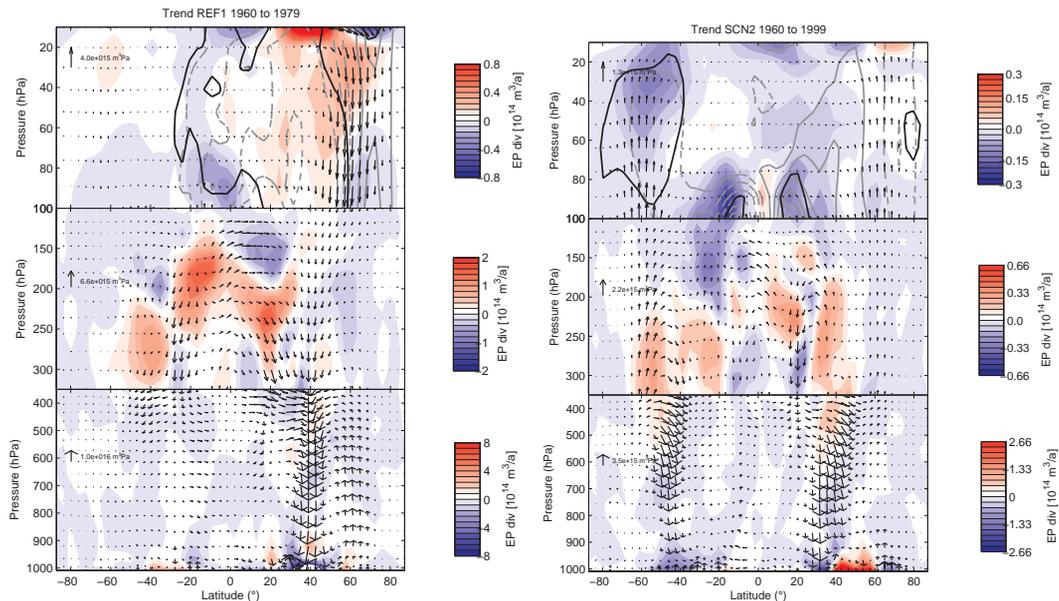


Fig. 10. As in Fig. 5, but showing the linear trend in EP fluxes and divergence and in the residual circulation for REF1, years 1960 to 1979 (left panels) and SCN2, years 1960 to 1999 (right panels). Contour intervals for the residual streamfunction is 2.5×10^7 kg/s/a in the left panel and 0.7×10^7 kg/s/a in the right panel.

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