- Huggins, M. L. J. Chem. Phys. 1941, 9, 440. (c) Huggins, M. L. Ann. N. Y. Acad. Sci. 1942, 43, 1.
- 4. Abrams, D.; Prausnitz, J. M. AIChE. J. 1975, 21, 116.
- 5. (a) Sanchez, I. C.; Lacombe, R. H. Nature 1974, 252, 381. (b) Sanchez, I. C.; Lacombe, R. H. J. Phys. Chem. 1976, 80, 2352. (c) Sanchez, I. C.; Lacombe, R. H. J. Phys. Chem. 1976, 80, 2568.
- 6. Kehiaian, H. V.; Grolier, P. E.; Bebson, G. C. J. Chem. Phys. 1978, 75, 1031.
- 7. (a) Okada, M.; Nose, T. Polymer J. 1981, 13, 399. ibid. 1981, 13, 591.
- 8. Panayiotou, C.; Vera, J. H. Polymer J. 1982, 14, 681.
- 9. Kumar, S. K.; Suter, U. W.; Reid, R. C. Ind. Eng. Chem. Res. 1987, 26, 2532.
- 10. Smirnova, N. A.; Victorov, A. I. Fluid Phase Equilibria 1987, 34, 235.
- 11. (a) You, S. S.; Yoo, K.-P.; Lee, C. S. Fluid Phase Equilibria 1994, 93, 193. ibid. 1994, 93, 215. (c) You, S. S.; Lee, C. S.; Yoo, K.-P. J. Supercritical Fluids 1993, 6, 69. ibid. 1994, 7, 251. (d) Shin, M. S.; Yoo, K.-P.; You, S. S.; Lee, C. S. Int. J. Thermophysics 1995, 16, 723. (e) Yoo, K.-P.; Shin, M. S.; Yoo, S. J.; You, S. S.; Lee, C. S. Fluid Phase Equilibria 1995, 111, 175. (f) Yoo, K.-P.; Kim, H. Y.; Lee, C. S. Korean J. Chem. Eng. 1995, 12, 277. ibid. 1995, 12, 289. (g) Yoo, K.-P.; Kim, J. S.; Kim, H. Y.; You, S. S.; Lee, C. S. J. Chem. Eng. Japan 1996, 29, 439. (h) Yoo, K.-P.; Lee, C.S. Fluid Phase Equilibria 1996, 117, 48. Yoo, S. J.; Yoo, K.-P.; Kim, H. Y.; Lee, C. S. Fluid Phase Equilibria

- 1996, 125, 21. (I) Yoo, S. J.; Yoo, K.-P.; Lee, C. S. J. Phys. Chem. B. 1997, 101, 1072.
- 12. Wilson, G. M. J. Am. Chem. Soc. 1964, 86, 127.
- 13. Hill, T. L. An Introduction to Statistical Mechanics; Addison Wesley: Massachusetts, 1960; p 296.
- 14. Timmermans, J. Physicochemical Constants of Pure Organic Compounds; Elsevier Scientific Publishing Company: New York, 1950; Vol. 1.
- 15. Reid, R. C.; Prausnitz, J. M.; Poling, B. E. The Properties of Gases and Liquids, 4th eds.; McGraw-Hill Book Co: New York, 1987.
- 16. Braker, W.; Mossman, A. L. Matheson Gas Data Book, 6th ed.; Matheson Gas Products, 1980.
- 17. van Krevelen, D. W. Properties of Polymers; Elsevier Scientific Publishing Company: New York, 1990; Ch. 4-
- 18. Bondi, A. Physical Properties of Molecular Crystals, Liquids and Glasses; Wiley: New York, 1986; Ch. 14.
- 19. Fredenslund, A.; Jones, R. L.; Prausnitz, J. M. AIChE. J. 1975, 21, 1086.
- 20. Kac, M.; Uhlenbeck, G. E.; Hemmer, P. C. J. Math. Phys. 1963, 4, 216.
- 21. Uhlenbeck, G. E.; Hemmer, P. C.; Kac. M. J. Math. Phys. 1964, 5, 60.
- 22. Guggenheim, E. A. Mol. Phys. 1965, 9, 43. ibid. 1965, 9, 199.
- 23. Longuet-Higgins, H. C.; Widom, B. Mol. Phys. 1964, 8,
- 24. Beret, S.; Prausnitz, J. M. Macromolecules 1975, 8, 536.

Oxygen Evolution Reaction at Electrodes of Single Phase Ruthenium Oxides with Perovskite and Pyrochlore Structures**

Eun-Ok Chi, Young-Uk Kwon, and Sun-il Mho*

Department of Chemistry, Sung Kyun Kwan University, Suwon 440-746, Korea *Department of Chemistry, Ajou University, Suwon 442-749, Korea Received May 12, 1997

Single phase ruthenium oxides with perovskite (ATi_{1-x}Ru_xO₃ (A=Ca, Sr)) and pyrochlore structure (Bi₂Ru₂O₇, Pb₂Ru₂O_{6.5}) have been prepared reproducibly by solid state reaction methods and their electrocatalytic activities for oxygen evolution have been examined by Tafel plots. Tafel slopes vary from a low value of 42 mV/decade up to 222 mV/decade at room temperature. The high exchange current densities and high Tafel slopes compared with those obtained from the RuO₂ DSA electrode at the crystalline single phase metal oxide electrodes suggest that they are better electrocatalysts at low overpotentials. A favorable change in the Tafel slope for the oxygen evolution reaction occurs as the ruthenium content increases. Substitution of Ti for Ru in the perovskite solid solutions enhanced their chemical stability by losing marginal electrochemical activity.

Introduction

The oxygen evolution and reduction reactions are of special importance in water electrolyzers, fuel cells, and batteries using air cathodes since slow kinetics of these reactions are the chief cause of efficiency losses.1 Transition metal oxides possess many of the desirable characteristics for practical electrodes for oxygen evolution.2 The poor stability of some metal oxides over anodic potential ranges where oxygen evolution takes place is a major problem. Dimensionally stable anode (DSA) materials have been attractive elec-

^{*}To whom correspondence should be addressed

^{**}This article is dedicated to Prof. Woon-kie Paik (Sogang University) in commemoration of his 60th birthday.

trocatalysts for both the chlorine and oxygen evolution reactions.³ DSA electrodes such as RuO_x are commonly prepared by thermal decomposition of RuCl₃ and Ru(NO)(NO₃)₃, which are applied as thin coating on metallic substrate, usually titanium.⁴⁻⁶ Generally, the oxide layers are nonstoichiometric with structural defects. The properties of the oxide layer, especially its catalytic effects depend on the thermal decomposition conditions such as concentration of the salt solution, temperature, atmosphere and duration of heating, and the substrate.⁷

Perovskite-type CaRuO₃ and SrRuO₃ are metallic conductors. Ru ions in these compounds can be replaced by Ti to form solid solutions ATi_{1-x}Ru_xO₃ (A=Ca, Sr) in the whole range of x. The conductivity mechanism changes from metallic to semiconducting as x decreases below 0.5.89 However, we expected that the chemical stability of these Ru compounds can be increased by the Ti substitution. These materials are expected to have a varying electrocatalytic activity and chemical stability depending on their chemical composition. Also, metal oxides with the pyrochlore structure such as Bi₂Ru₂O₇ and Pb₂Ru₂O_{6.5}, which are all characterized by an infinite, 3-dimensional cubic array of oxide ion as in the perovskite structure, are known to show encouraging performance for air electrodes. 10,11 In the present study, single phase ruthenium oxides with perovskite and pyrochlore structures have been prepared reproducibly by solid state reaction methods and their electrocatalytic activities for oxygen evolution have been examined by Tafel plots.

Experimental

 $SrTi_{1,r}Ru_rO_3$ (x=1.0, 0.67, 0.33), $CaTi_{1,r}Ru_rO_3$ (x=1.0, 0.8, 0.6, 0.2, 0) having a perovskite bulk structure, where ruthenium and titanium have an octahedral coordination of 6 oxygens, have been prepared by solid state reactions for electrode materials. The stoichiometric amounts of TiO2 and SrCO₃, and RuO₂ and SrCO₃ were mixed and heated at 1100 °C for 20-30 hours in air to prepare SrTiO₃ and SrRuO₃, respectively. The resulting powders were ground separately, and the mixture of (1-x) SrTiO₃ and x SrRuO₃ was heated at 1100 °C for 30 hours to synthesize $SrTi_{1-x}Ru_xO_3$ (x=1.0, 0.67, 0.33). $CaTi_{1,x}Ru_xO_3$ (x=1.0, 0.8, 0.6, 0.2, 0) were prepared by a similar procedure at 1300 °C. Bi₂Ru₂O₇ and Pb₂Ru₂O_{6.5} have also been prepared by solid state reactions. The stoichiometric amounts of Bi₂O₃ and RuO₂, and PbO and RuO2 were mixed. The mixture was placed in an alumina boat and sealed in a quartz tube in a vacuum line, which was heated at 700 °C for 15 hrs. and 900 °C for 24 hrs. consecutively for Bi₂Ru₂O₇, and at 850 °C for 20 hrs. and at 900 °C for 10 hrs. for Pb₂Ru₂O_{6.5}. The structures of the compounds were determined by a Rigaku X-ray powder diffractometer with Cu Ka₁ radiation.

To make electrodes for the electrochemical measurements, the products were pressed into pellets after mixing the powder with a few drops of binder (4% polyvinyl alcohol solution). The pellets were dried and heated at 1400 °C. One face of the disk was ground with SiC based emery papers (grit #1200, 1500) and polished with alumina (stepwise 1.0, 0.3, 0.05 μm), sonicated and washed thoroughly before use. For comparison, the RuO₂ DSA electrode was prepared

Table 1. Lattice parameters of the single phase ruthenium oxides with perovskite and pyrochlore structure

	Lattice type	Cell parameter	(Å)
SrRuO ₃	cubic	a=3.932(2)	
$SrTi_{1/3}Ru_{2/3}O_3$	cubic	a=3.929(1)	
$SrTi_{2/3}Ru_{1/3}O_3$	cubic	a=3.916(2)	
CaRuO ₃	orthorhombic	a=5.533(3), b=5.380(3),	c=7.644(5)
CaTi _{0.2} Ru _{0.8} O ₃	orthorhombic	a=5.531(4), b=5.364(3),	c=7.673(5)
CaTi _{0.4} Ru _{0.6} O ₃	orthorhombic	a=5.530(5), b=5.356(5),	c=7.672(4)
Bi ₂ Ru ₂ O ₇	cubic	a=10.315(1))
$Pb_2Ru_2O_{6.5}$	cubic	a=10.267(4))

by a common technique, *i.e.*, thermal decomposition of a metal salt, RuCl₃, on Ti wire.⁵

All electrochemical measurements were made by a Princeton Applied Research Model 263 potentiostat/galvanostat. A three electrode cell system was used: the metal oxide electrode as a working electrode, a platinum wire spiral counter electrode, and a saturated calomel reference electrode (SCE). All potentials reported here are referred to the SCE. The electrolyte was 1.0 M KOH in distilled deionized water. For the Tafel plot, the equilibrium potentials were measured with applying constant currents by chronopotentiometry.

Results and Discussion

The lattice parameters of the compounds of present study are listed in Table 1. While the compounds in the SrTi_{1-x}Ru_xO₃ solid solution crystallize in an ideal cubic perovskite structure, the Ca-analogues show orthorhombic distortions. The unit cell parameters decrease gradually with the increase in Ti content in both cases. The pyrochlore compounds show the same lattice parameters as those in the literature indicating that their compositions are close to the ideal ones, *i.e.*, Bi₂Ru₂O₇ and Pb₂Ru₂O_{6.5}.

A plot of log *i* (current) vs. η (overpotential), known as a Tafel plot, is a useful device for evaluating kinetic parameters.¹² For an anodic branch, the anodic current shows a relationship to the overpotential:

$$i=i_o \exp[(1-\alpha)nF\eta/RT]$$

or $\eta=[2.3RT/(1-\alpha)nF]$ (log $i-\log i_o$)

where n, F, R, and T are number of electrons, Faraday con-

Table 2. Values of the slopes and exchange current densities from Tafel plots for oxygen evolution in 1.0 N KOH solution at various metal oxides

Metal oxide electrodes		Slope (V/decade)	Exchange current density (A/cm ²)
perovskite	$SrTi_{1/3}Ru_{2/3}O_3$	0.096	1.49×10^{-7}
structure	SrTi _{2/3} Ru _{1/3} O ₃	0.161	1.30×10^{-5}
	CaRuO ₃	0.087	3.40×10^{-5}
	$CaTi_{0.2}Ru_{0.8}O_3$	0.197	6.85×10^{-5}
	$CaTi_{0.4}Ru_{0.6}O_3$	0.202	1.53×10^{-5}
	$CaTi_{0.8}Ru_{0.2}O_3$	0.254	2.82×10^{-5}
pyrochlore	Bi ₂ Ru ₂ O ₇	0.075	6.30×10^{-8}
structure	$Pb_2Ru_2O_{6.5}$	0.152	1.58×10^{-5}
DSA type	RuO ₂ on Ti	0.056	1.60×10^{-8}



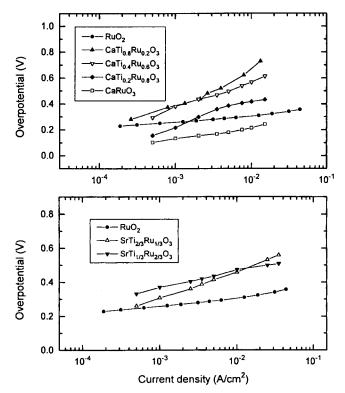


Figure 1. Tafel plots for oxygen evolution in 1.0 N KOH solution at single phase ruthenium oxides with perovskite structure. (a) $CaTi_{1-x}Ru_xO_3$ (x=0, 0.2, 0.6, 0.8, 1.0). (b) $SrTi_{1-x}Ru_xO_3$ (x=2/3, 1/3).

stant, gas constant, and temperature, respectively. The slope of the linear segment is $(1-\alpha)nF/2.3RT$. It extrapolates to an intercept of log i_o . The transfer coefficient α and the exchange current i_o are readily accessible from this kind of presentation. Table 2 summarizes the values of the slopes and exchange current densities from the Tafel plots for oxygen evolution reaction at single phase ruthenium oxide electrodes with perovskite and pyrochlore structures including $SrTi_{1-x}Ru_rO_3$ (x=1.0, 0.67, 0.33) and $CaTi_{1-x}Ru_rO_3$ (x =1.0, 0.8, 0.6, 0.2, 0), and $Bi_2Ru_2O_7$ and $Pb_2Ru_2O_{6.5}$. Tafel plots for oxygen evolution in 1.0 N KOH at various metal oxide electrodes with perovskite structure are also presented in Figure 1. Tafel slopes vary from a low value of 42 mV/ decade up to 222 mV/decade at room temperature. A favorable change in the Tafel slope for the oxygen evolution reaction occurs as the ruthenium content is increased, which corresponds to an increase in the transfer coefficient $(1-\alpha)$. $SrTi_{1/3}Ru_{2/3}O_3$ and $CaRuO_3$ and $Bi_2Ru_2O_7$ show lower Tafel slopes (better catalytic effects) for oxygen evolution reaction than the other mixed metal oxide electrodes in this study. The Tafel slopes from our metal oxide electrodes are in the same range as other published works, 30-235 mV/decade.1 Therefore, the mechanism of oxygen evolution at the metal oxides could be considered to be similar to the previous ones generally accepted13-15 and that only one electron is involved in the rate determining step at these electrodes. The reaction pathway that is often suggested for oxygen evolution in alkaline solution is as follows¹⁴:

S+OH → S-OH + e (S is transition metal ion site)

SOH + OH⁻
$$\rightarrow$$
 SO⁻ + H₂O
SO⁻ \rightarrow SO + e⁻
2 SO \rightarrow 2 S + O₂

The rate determining step (rds) for oxygen evolution of metal oxides varied depending on the electrode material. In addition, the valence states of the cations were found to play an important role in oxygen evolution on metal oxide electrocatalysts. For most cases, one electron process is suggested to be the rate determining step for oxygen evolution processes.¹⁵ From the present study, it was difficult to identify which process is the rds for oxygen evolution at these electrodes. More extensive analysis is required to identify the rds for oxygen evolution.

Table 2 shows that the exchange current densities are in the range from 5.1×10^{-5} to 2.1×10^{-2} A/cm². The exchange current density is based on the geometric surface area and the roughness factor is not taken into account. Because of the variability in Tafel slopes obtained from different electrodes, a direct comparison of the exchange current densities as a measure of electrocatalytic activity is not considered to be very useful.1 For practical purposes, a more meaningful comparison of the electrocatalytic activity of anodes is the overpotential for oxygen evolution that is obtained at a specific current density. CaRuO₃ shows the best performance and shows even lower overpotentials than RuO₂ DSA electrode at all current densities (Figure 1a). CaTi_{0.2}Ru_{0.8}O₃ shows lower overpotentials at lower current densities and higher overpotentials at higher current densities than the RuO₂ DSA electrode. Higher overpotentials with a large Tafel slope for the oxygen evolution reaction than at the RuO₂ DSA electrode are obtained at oxide electrodes with the lower ruthenium content, CaTi_{1-x}Ru_xO₃ (x= 0.2, 0.6). It is well known that in DSA type electrodes (containing RuO₂/TiO₂), the electrocatalytic activity for oxygen evolution decreases dramatically when the ruthenium content is less than 30 mol%. 15 Higher overpotentials for the oxygen evolution reaction than at the RuO2 DSA electrode are obtained at SrTi_{1-x}Ru_xO₃ (x=0.33, 0.67) (Figure 1b). Although SrTi_{1/3}Ru_{2/3}O₃ shows lower Tafel slopes for oxygen evolution compared with other oxides in this study, overpotentials are larger at all current densities than the RuO₂ DSA electrode. The SrRuO₃ electrode is too unstable in alkaline solutions to make electrochemical measurements. Electrodes of perovskite type $SrTi_{1-x}Ru_xO_3$ (x=1.0, 0.33, 0.67) and $CaTi_{1x}Ru_xO_3$ (x=1.0, 0.8, 0.6, 0.2) disintegrate during the oxygen evolution reaction in alkaline solutions.

The presence of the mixed valence states of ruthenium in the nonstoichiometric oxides with structural defects, which are prepared by thermal decomposition of metal salts at low temperature (ca. 400 °C), is known to correlate with the high electrocatalytic activity. Single phase ruthenium dioxide electrodes of perovskite structures prepared by solid state reaction methods at high temperatures (ca. 1200 °C) would have less defects resulting in the valence state of four and a smaller surface area with a crystalline structure. This trend is speculated as the main reason for not showing the expected catalytic effects on the oxygen evolution reaction. In general, the oxygen evolution reaction at the single phase ruthenium oxide electrodes studied in this work results in high exchange current densities and high Tafel

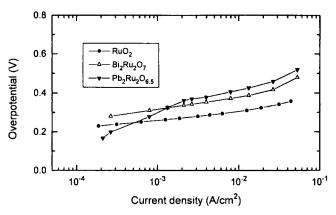


Figure 2. Tafel plots for oxygen evolution in 1.0 N KOH solution at Bi₂Ru₂O₇ and Pb₂Ru₂O_{6.5} electrodes of pyrochlore structure.

slopes, compared with the RuO₂ DSA electrode. Therefore, the crystalline single phase metal oxide electrodes appear better electrocatalysts at low overpotentials.

Some oxides with the pyrochlore structure, A₂B₂O_{7-y} (A= Pb, Bi; B=Ru, Ir; $0 \le y \le 0.5$), have been found to be active for both oxygen evolution and reduction reactions and are therefore candidate materials for a bifunctional oxygen electrode. 10,11 Oxides with the pyrochlore structure have the ideal chemical formula A₂B₂O₆O'; they contain a cubic B₂O₆ framework of corner-shared octahedra that is interpenetrated by an A₂O' subarray. The interpenetrating A₂O' array is flexible, and it may contain O' vacancies. Whereas Bi₂Ru₂O₇ has the ideal pyrochlore structure, the Pb analogues contain O'-site vacancies. 16-18 Tafel plots for oxygen evolution in 1.0 N KOH at the electrodes of pyrochlores, Bi₂Ru₂O₇ and Pb₂Ru₂O_{6.5}, are shown in Figure 2. In the case of Bi₂Ru₂O₇, although the overpotentials are larger at all current densities, the Tafel slope for the oxygen evolution reaction is about 60 mV/decade and the linear region in the Tafel plot is comparable to that observed at RuO2 DSA type electrodes. With Pb₂Ru₂O_{6.5}, the slopes for the oxygen evolution reaction are higher than RuO₂ DSA type electrodes. In addition, Tafel slope changes and the linear region in the plot is very limited. The change in the Tafel slope is known to be associated with a change in the valence state or the crystal structure of the metal oxide, e.g., α - or β -Ni(OH)₂.¹⁹

The compounds in the present study and RuO₂ have RuO₆ octahedral units as the building blocks in common but in different modes of connectivity. The octahedra are edgeshared in RuO₂ while they are corner-shared in the others. Since edge-sharing will result in broader d-band through direct d-orbital interactions than in corner-sharing, the mobility of conduction d-electrons will be the highest in RuO₂. Among the structures with corner-shared RuO₂ octahedra, the band width will decrease with the degree of bending in Ru-O-Ru bond angles, which increases along the order, SrRuO₃>CaRuO₃>Bi₂Ru₂O₇>Pb₂Ru₂O_{6.5}. However, in the last two pyrochlores 6s²-electrons of Bi and Pb would contribute to the conduction electron concentrations to raise their conductivities higher than those of the perovskites. Tiincorporation into the ARuO₃ (A=Ca, Sr) perovskites will localize d-electrons of Ru by creating random potentials and, at the same time, by reducing the d-electron concentration.

Therefore, the bulk conductivity of the Ru(IV) compounds will decrease according to the sequence, RuO₂>Bi₂Ru₂O₇ \geq Pb₂Ru₂O_{6.5}>ARuO₃>ATi_{1.x}Ru_xO₃ (A=Ca, Sr). Substitution of Ti for Ru of the perovskite solid solutions causes the loss of electrochemical activity. However, in the case of SrRuO₃ electrode, its chemical stability is enhanced by Ti substitution by losing marginal electrochemical activity. Although we have studied only few compositions in this study, more detailed experiments with narrow composition intervals may reveal the optimal composition of acceptable chemical stability and electrochemical activity in these compounds.

Conclusion

Single phase ruthenium oxides with perovskite and pyrochlore structure including SrTi_{1-x}Ru_xO₃ (x=1.0, 0.67, 0.33) and $CaTi_{1-x}Ru_xO_3$ (x=1.0, 0.8, 0.6, 0.2, 0), and $Bi_2Ru_2O_7$ and Pb₂Ru₂O_{6.5} have been prepared reproducibly by solid state reaction methods. Their electrocatalytic activities for oxygen evolution have been examined by Tafel plots. The slopes of Tafel plots for oxygen evolution in 1.0 N KOH at various single phase metal oxide electrodes vary from a low value of 42 mV/decade up to 222 mV/decade at room temperature. A favorable change in the slope occurs as the ruthenium content increases, which corresponds to an increase in the transfer coefficient $(1-\alpha)$ for the oxygen evolution. SrTi_{1/3}Ru_{2/3}O₃ and CaRuO₃ and Bi₂Ru₂O₇ show lower Tafel slopes (better catalytic effects) for the oxygen evolution reaction than the other mixed metal oxide electrodes in this study. The bulk conductivity of the ruthenium compounds decreases along the sequence, RuO₂>Bi₂Ru₂O₇≥ Pb₂Ru₂O₆₅>ARuO₃>ATi_{1-r}Ru_rO₃ (A=Ca, Sr). Substitution of Ti for Ru of the perovskite solid solutions causes the loss of electrochemical activity with the enhancement of its chemical stability. Single phase ruthenium dioxide electrodes of perovskite structure prepared by solid state reactions at high temperatures (ca. 1200 °C) have less defects, which results in the valence state of four and a smaller surface area with a crystalline structure. This is speculated as the main cause for not showing the expected favorable catalytic effect on oxygen evolution. In general, oxygen evolution reaction at the single phase ruthenium oxide electrodes studied in this work results in high exchange current densities and high Tafel slopes, compared with the RuO2 DSA electrode. Therefore, the crystalline single phase metal oxide electrodes are suggested to be better electrocatalysts at low overpotentials.

Acknowledgment. This work has been supported in part by the Ministry of Education (BSRI-96-3421) and in part by the Korea Science and Engineering Foundation (96-0501-05-01-3).

References

- Kinoshita, K. Electrochemical Oxygen Technology; Wiley: New York, 1992; and references cited therein.
- 2. Trasatti, S. Electrochim. Acta 1984, 29, 1503.
- 3. Trasatti S.; Lodi, G. in *Electrodes of Conductive Metallic Oxides*; Trasatti, S. Ed.; Elsevier: New York, 1980; Chap. 7 and references cited therein.

- 976
- Yeo, R. S.; Orehotsky, J.; Visscher, W.; Srinivasan, S. J. Electrochem. Soc. 1981, 128, 1900.
- Park, S.; Mho, S.; Chi, E.; Kwon, Y.; Yeo, I. Bull. Korean Chem. Soc. 1995, 16, 82.
- Ardizzone, S.; Falciola, M.; Trasatti, S. J. Electrochem. Soc. 1985, 136, 1545.
- 7. Melsheimer J.; Ziegler, K. Thin Solid Films 1988, 163, 301
- Kobayashi, H.; Nagata, M.; Kanno, R.; Kawamoto, Y. Mater. Res. Bull. 1994, 29, 1271.
- 9. Dabrowski, B.; Jorgensen, J.; Minks, D.; Pei, S.; Richards, K.; Vanfleet, H.; Kecker, K. *Physica C* **1989**, 162, 99.
- 10. Egdell, R. G.; Goodenough, J. B.; Hamnett, A.; Naish, C. C. J. Chem. Soc., Faraday Trans. 1983, 1, 79.
- Gokagac G.; Kennedy, B. J. J. Electroanal. Chem. 1993, 353, 71.

- 12. Bard, A. J.; Faulkner, L. R. Electrochemical Methods; Wiley: New York, 1980; p 106.
- 13. Bockris, J. O'M.; Otagawa, T. J. Electrochem. Soc. 1984, 131, 290.
- 14. Matsumoto, Y.; Manabe, H.; Sato, E. J. Electrochem. Soc. 1980, 127, 811.
- Rasiyah, P.; Tseung, A. C. C. J. Electrochem. Soc. 1983, 130, 2384.
- Beyerlein, R. A.; Horowitz, H. S.; Longo, J. M.; Leonowicz, E.; Jorgensen, J. D.; Rotella, F. J. J. Solid State Chem. 1984, 51, 253.
- 17. England, W. A.; Cross, M. G.; Mamnett, M.; Wiseman, P. J.; Goodenough, J. B. Solid State Ionics 1980, 1, 231.
- 18. Fujishima, A.; Honda, K. Nature (London) 1972, 238, 37.
- Gennero de Chialvo, M. R.; Chialvo, A. C. *Electrochim. Acta* 1988, 33, 825.

Ligand Field Approach to 4d¹ Magnetism Based on Intermediate Field Coupling Scheme

Jin-Ho Choy* and Jong-Young Kim

Department of Chemistry and Center for Molecular Catalysis, College of Natural Sciences, Seoul National University, Seoul 151-742, Korea Received May 16, 1997

The magnetic susceptibilities of molybdenum ions with $4d^1$ electronic configuration in the octahedral crystal field were calculated on the basis of ligand field theory. The experimental magnetic susceptibilities for molybdenum ions, which are stabilized at the octahedral site in the perovskite lattice of $Ba_2ScMo^VO_6$ and $Sr_2YMo^VO_6$, were compared with the theoretical ones. We have tried to fit their temperature dependences of magnetic susceptibility with ligand field parameters, spin-orbit coupling constant ζ_{so} , and orbital reduction parameter κ according to intermediate field coupling and strong field theory. Strong field coupling theory could not explain experimental curves without unrealistically large axial ligand field, since it ignores the mixing up between different state via spin-orbit interaction and ligand field. On the other hand, the intermediate field coupling theory could successfully reproduce experimental data in octahedral and trigonal ligand field. The fitting result demonstrates not only the fact that spin-orbit interaction is primarily responsible for the variation of magnetic behavior but also the fact that effective orbital overlap, enhanced by cubic crystal structure, reduces significantly orbital angular momentum as indicated by κ parameter.

Introduction

In order to interpret the magnetic property of second and third row element, more rigorous calculation is required than the simplified one such as weak field coupling, which has been known to be quite successful in interpreting the magnetic property of first row element. There is a marked difference in the magnetic and optical properties between the $4d^n$ and $5d^n$ transition metal ions and $3d^n$ ones, since spin-orbit interaction as well as ligand field are much stronger than interelectronic repulsion in the former, compared to the latter. $^{1-7}$ In fact, in going from 3d toward 4d and 5d, we notice that the interelectronic repulsion among the outer shell

electrons gets weaker while the spin-orbit interaction becomes stronger. Therefore, the influences of spin-orbit interaction and ligand field need to be relevantly considered. As it is well-known, in the weak field coupling scheme, the starting orbital is d and the zeroth order wavefunctions are

$$|d, m_l=2>, |d, 1>, |d, 0>, |d, -1>, \text{ and } |d, -2>$$

and the zeroth order Hamiltonian consists of the free ion Hamiltonian, given by

$$H_W = H_{core} - \sum_i \frac{Ze^2}{r_i} + \sum_{i \neq j} \frac{e^2}{r_{ij}}$$

where the summation is over all the d electrons. Here H_{core} involves the electrons in the closed shells. The symmetry group of H_w is the full rotation group R_3 in the weak field

^{*}To whom all correspondence should be addressed